

Balanced Flow on the Spinning Globe

Peter Lynch
School of Mathematical Sciences
University College Dublin

EMS Silver Medal Lecture
Met Soc & UCD
Science Hub, Belfield
Thursday 20th November 2014



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



Galileo Galilei (1564–1642)

Formulated **law of falling bodies**
... verified by measurements.

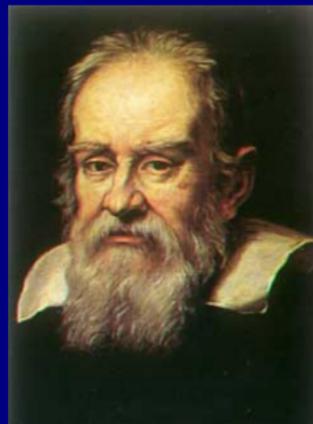
Constructed a **telescope**, and found

- ▶ lunar craters
- ▶ four moons of Jupiter

Galileo invented the **thermometer**

Evangelista Torricelli invented the **barometer**

Thus began quantitative meteorology.



Galileo Galilei and Leaning Tower of Pisa

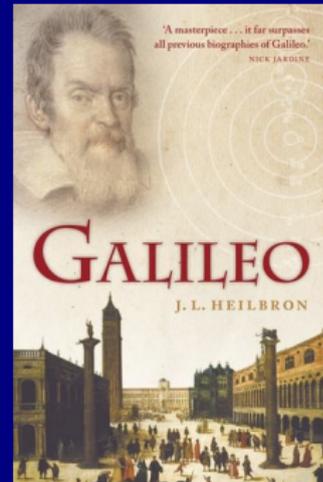


Leaning Tower

We have **Viviani's** word that Galileo dropped various weights from the Leaning Tower ...

“... to the dismay of the philosophers, different weights fell at the same speed ...”

Heilbron, John. *Galileo*. Oxford University Press, 2010



Galileo on the Universe

The Assayer (IL SAGGIATORE)
was published in Rome in 1623.

[The universe] ... is written in the
language of mathematics ...
without which it is ... impossible
to understand a single word of it.



As easy as A, B, C

Three-term equation:

$$A + B + C = 0$$



As easy as A, B, C

Three-term equation:

$$A + B + C = 0$$

Suppose one term is small relative to the others:

$$A \text{ SMALL} \implies B + C \approx 0$$



As easy as A, B, C

Three-term equation:

$$A + B + C = 0$$

Suppose one term is small relative to the others:

$$A \text{ SMALL} \implies B + C \approx 0$$

There are three possibilities:

$$A \text{ SMALL} \implies B + C \approx 0$$

$$B \text{ SMALL} \implies A + C \approx 0$$

$$C \text{ SMALL} \implies A + B \approx 0$$



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



A Most Surprising Property of Atmospheric & Oceanic Motion

The motion of the atmosphere and ocean systems is **remarkably persistent**.

Why doesn't air rush in to fill low pressure areas?



A Most Surprising Property of Atmospheric & Oceanic Motion

The motion of the atmosphere and ocean systems is **remarkably persistent**.

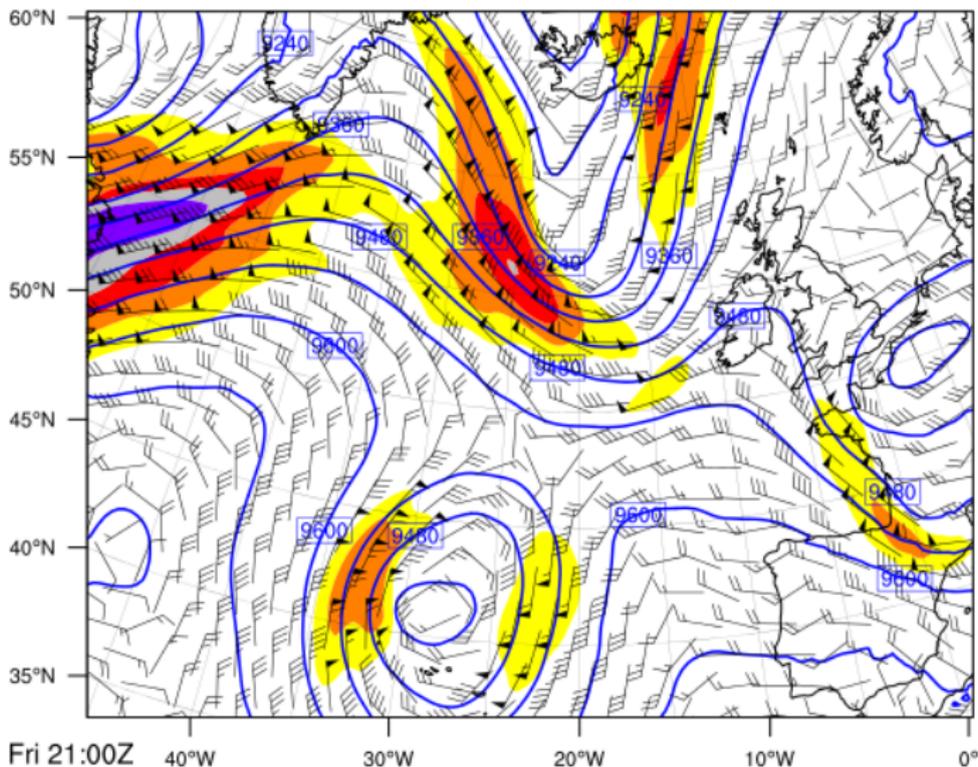
Why doesn't air rush in to fill low pressure areas?

The crucial factor is the **rotation of the Earth**.



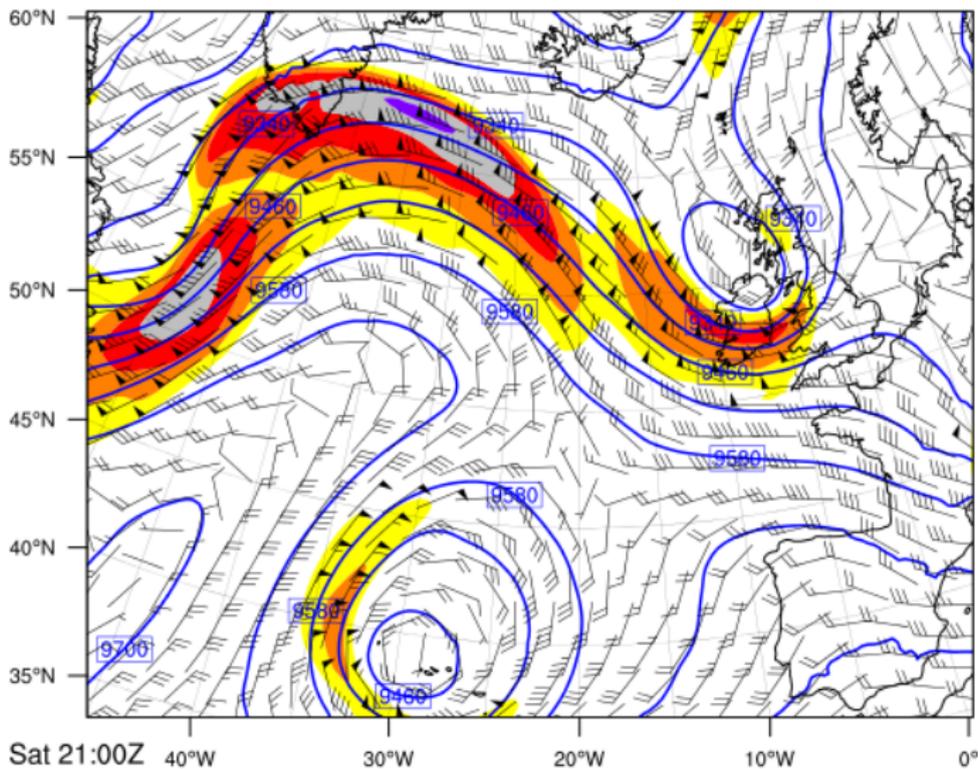
IT: 2014-07-25_06:00
UCD MCC FORECAST VT: 2014-07-25_21:00

Wind Speed (kts)
Height (m) at 300 hPa
Wind (kts) at 300 hPa



IT: 2014-07-25_06:00
UCD MCC FORECAST VT: 2014-07-26_21:00

Wind Speed (kts)
Height (m) at 300 hPa
Wind (kts) at 300 hPa



Jule Charney



*“If a stone is thrown into an **infinite resting ocean**, the gravitational oscillations engendered will radiate their energy to infinity, leaving the ocean . . . undisturbed;*



Jule Charney



*“If a stone is thrown into an **infinite resting ocean**, the gravitational oscillations engendered will radiate their energy to infinity, leaving the ocean ... undisturbed;*

*“If a stone is thrown into an **infinite rotating ocean**, some of the energy ... will be converted into rotational motions ... and these will persist”*

[Planetary Fluid Dynamics: *Dynamic Meteorology*, Ed. P. Morel, 1973]





Intro

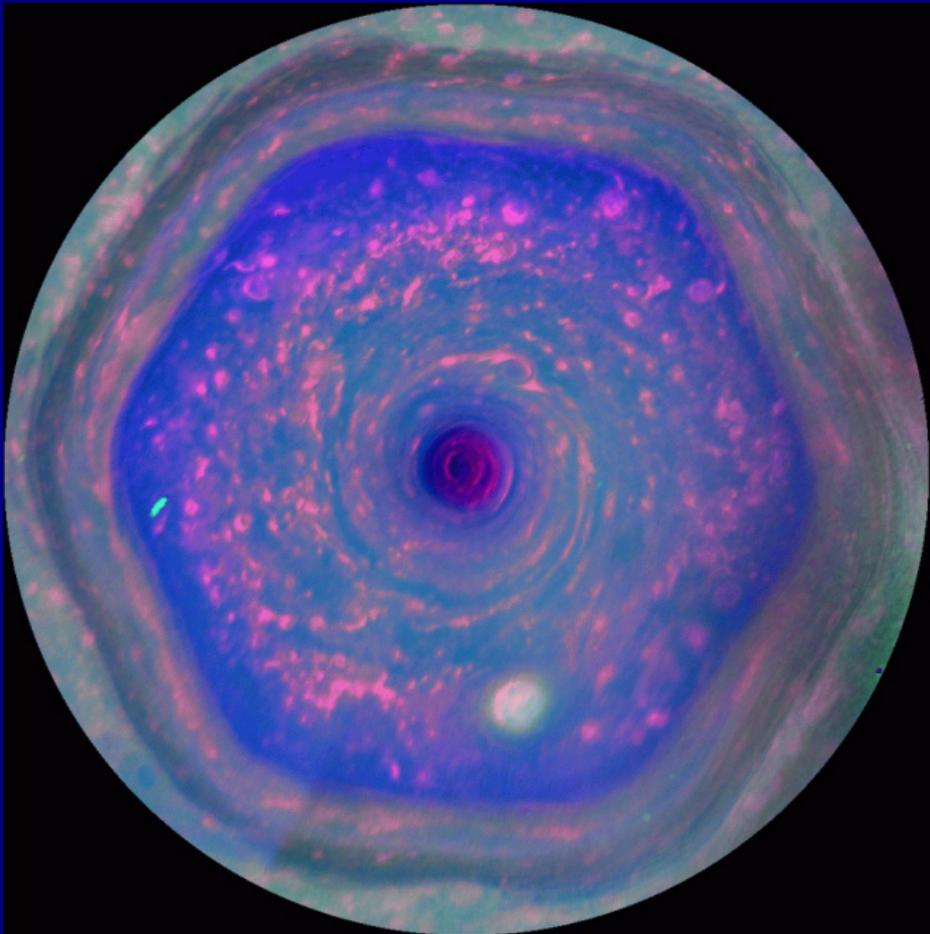
Balance

Foucault

Coriolis

LFR

ENIAC



Intro

Balance

Foucault

Coriolis

LFR

ENIAC

20th Century Reanalysis Project

A global reanalysis dataset spanning the entire twentieth century ...

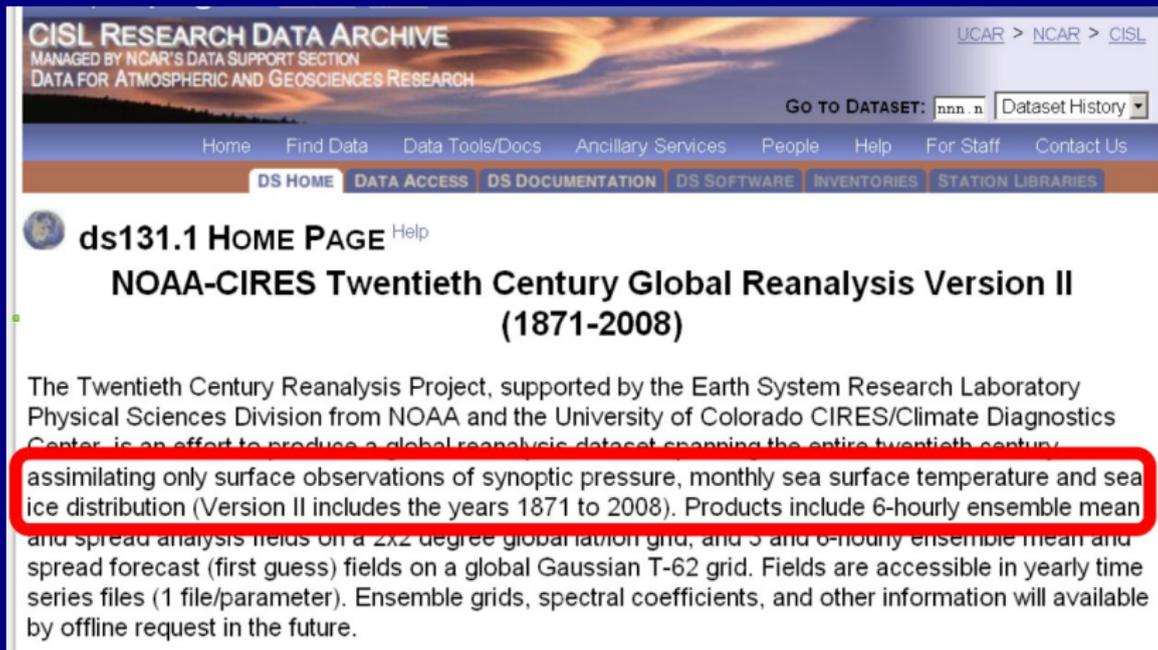
Assimilating only **surface pressure observations ...
... the analysis covers the entire troposphere.**

Resolution:

T62 (300km), 28 Levels. 56-Member Ensemble.



20th Century Reanalysis (20CRv2)



CISL RESEARCH DATA ARCHIVE
MANAGED BY NCAR'S DATA SUPPORT SECTION
DATA FOR ATMOSPHERIC AND GEOSCIENCES RESEARCH

UCAR > NCAR > CISL

Go TO DATASET: Dataset History ▾

Home Find Data Data Tools/Docs Ancillary Services People Help For Staff Contact Us

DS HOME DATA ACCESS DS DOCUMENTATION DS SOFTWARE INVENTORIES STATION LIBRARIES

 **ds131.1 HOME PAGE** [Help](#)

NOAA-CIRES Twentieth Century Global Reanalysis Version II (1871-2008)

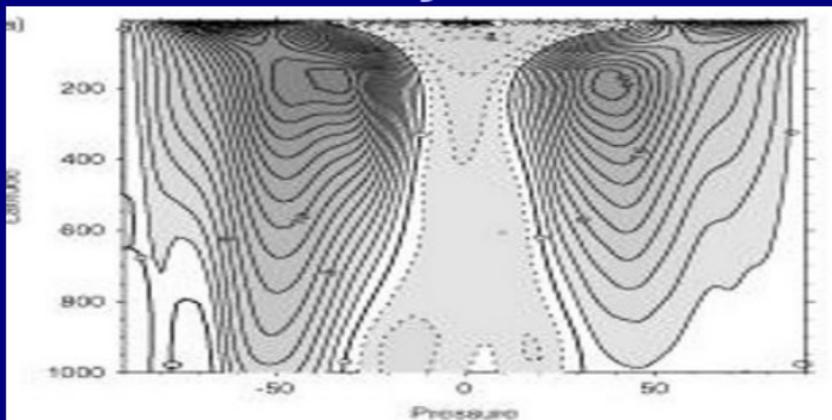
The Twentieth Century Reanalysis Project, supported by the Earth System Research Laboratory Physical Sciences Division from NOAA and the University of Colorado CIRES/Climate Diagnostics Center, is an effort to produce a global reanalysis dataset spanning the entire twentieth century.

assimilating only surface observations of synoptic pressure, monthly sea surface temperature and sea ice distribution (Version II includes the years 1871 to 2008). Products include 6-hourly ensemble mean and spread analysis fields on a 2x2 degree globalation grid, and 3 and 6-hourly ensemble mean and spread forecast (first guess) fields on a global Gaussian T-62 grid. Fields are accessible in yearly time series files (1 file/parameter). Ensemble grids, spectral coefficients, and other information will available by offline request in the future.

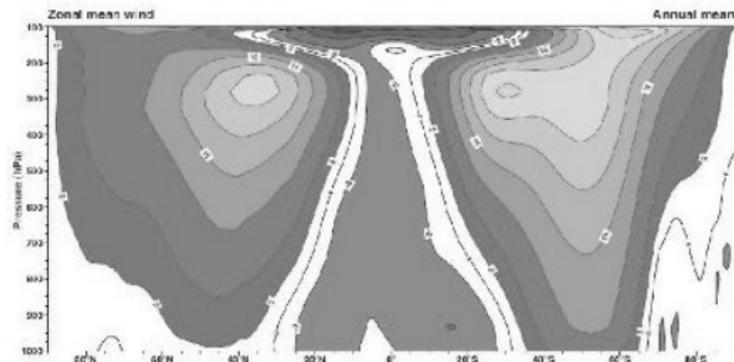


Mean Zonal Wind Analysis

20CR



ERA40



20th Century Reanalysis: Conclusion

How do they do that?

How do they reconstruct the troposphere from surface observations?



20th Century Reanalysis: Conclusion

How do they do that?

How do they reconstruct the troposphere from surface observations?

Reconstruction of the complete three-dimensional structure of the troposphere is possible ...

... because the atmosphere is in a state of balance.



20th Century Reanalysis: Conclusion

How do they do that?

How do they reconstruct the troposphere from surface observations?

Reconstruction of the complete three-dimensional structure of the troposphere is possible ...

... because the atmosphere is in a state of balance.

ERA-CLIM2: Ongoing ECMWF Project.
20th Century Reanalysis coming soon.

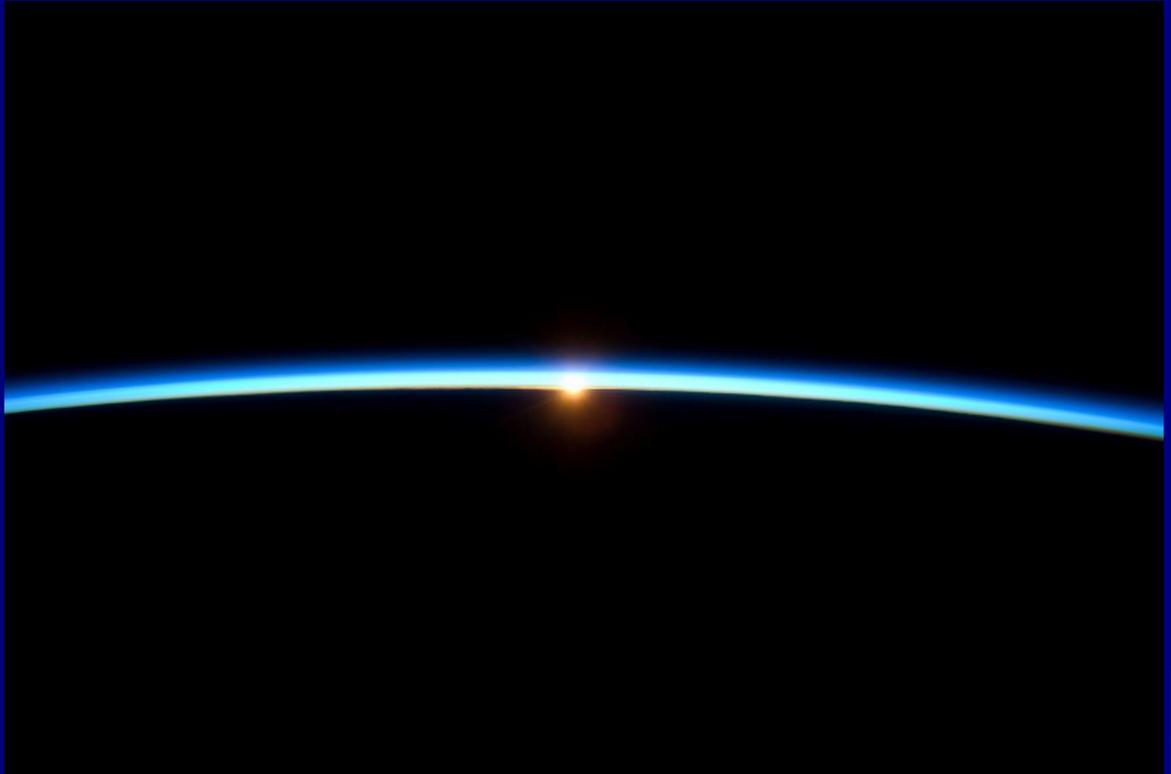


Examples of Balance in the Atmosphere

- ▶ **Hydrostatic balance**
- ▶ **Geostrophic balance**
- ▶ **Quasi-nondivergence**
- ▶ **Quasi-incompressibility**
- ▶ **Ocean atmosphere balance**
- ▶ **Energy balance**
- ▶ **Ice sheet balance**
- ▶ **Etc., etc., etc.**



The Thin Atmosphere



Hydrostatic Balance

What keeps the air aloft?

Something must be balancing gravity.

What is it?



Hydrostatic Balance

What keeps the air aloft?

Something must be balancing gravity.

What is it?

For a **parcel of air**:

- ▶ The air below is pushing it upwards.
- ▶ The air above is pushing it down.
- ▶ The push upwards is greater.
- ▶ The difference balances the pull of gravity.



Vertical Equation of Motion

Examine the terms in the vertical equation

$$\frac{dw}{dt} = \frac{u^2 + v^2}{r} + 2\Omega u \cos \phi - \frac{1}{\rho} \frac{\partial p}{\partial z} - g + F_z.$$

Vertical pressure gradient force and gravity dominate.



Vertical Equation of Motion

Examine the terms in the vertical equation

$$\frac{dw}{dt} = \frac{u^2 + v^2}{r} + 2\Omega u \cos \phi - \frac{1}{\rho} \frac{\partial p}{\partial z} - g + F_z.$$

Vertical pressure gradient force and gravity dominate.

Remember:

$$A + B + C = 0 \text{ with } A \text{ small means } B + C = 0.$$



Vertical Equation of Motion

Examine the terms in the vertical equation

$$\frac{dw}{dt} = \frac{u^2 + v^2}{r} + 2\Omega u \cos \phi - \frac{1}{\rho} \frac{\partial \rho}{\partial z} - g + F_z.$$

Vertical pressure gradient force and gravity dominate.

Remember:

$$A + B + C = 0 \text{ with } A \text{ small means } B + C = 0.$$

Keeping just the two large terms, we have:

$$\frac{\partial \rho}{\partial z} = -g\rho$$



Hydrostatic Balance

- ▶ The vertical wind is generally very small.
- ▶ There is balance between the vertical pressure gradient force and gravity.
- ▶ This balance is called *hydrostatic balance*.

$$\frac{\partial p}{\partial z} + g\rho = 0$$



Examples of Balance in the Atmosphere

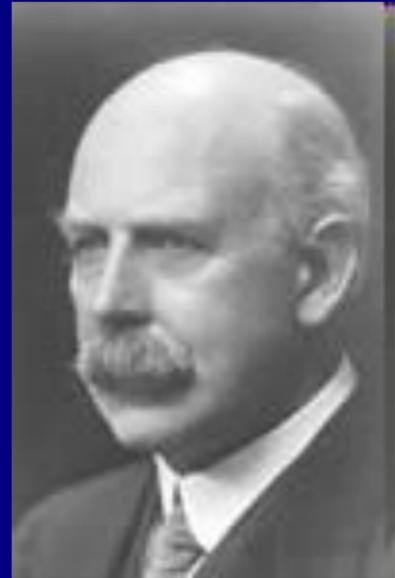
- ▶ **Hydrostatic balance**
- ▶ **Geostrophic balance**
- ▶ **Quasi-nondivergence**
- ▶ **Quasi-incompressibility**
- ▶ **Ocean atmosphere balance**
- ▶ **Energy balance**
- ▶ **Ice sheet balance**
- ▶ **Etc., etc., etc.**



Geostrophic Balance

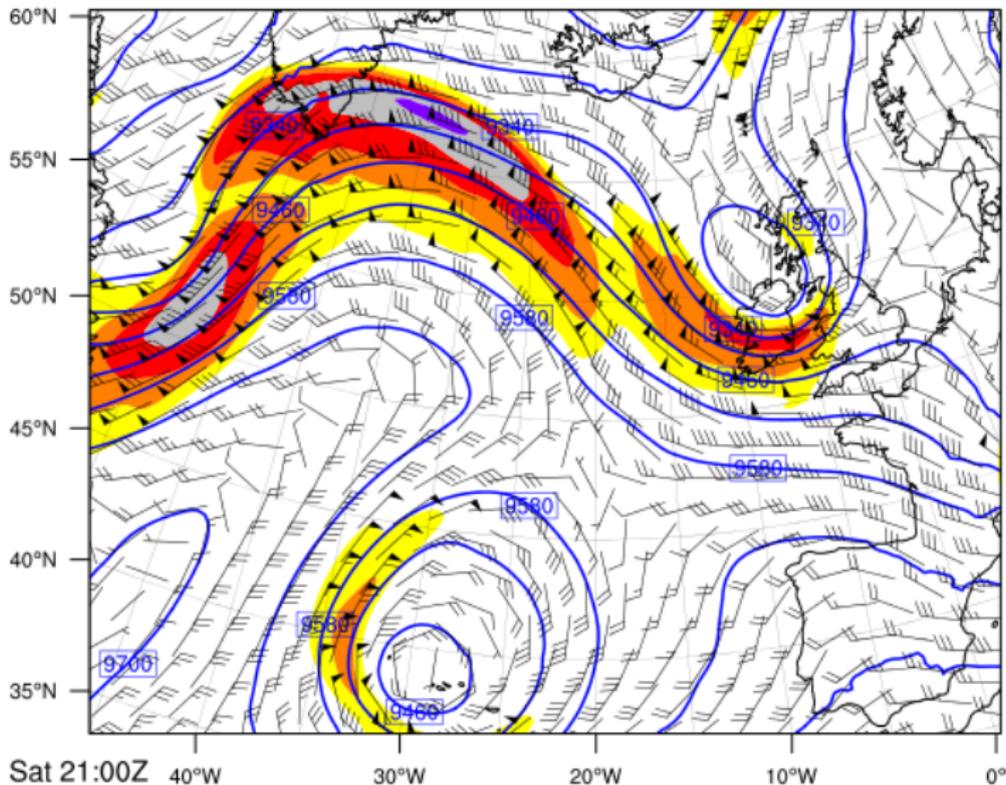
$\gamma\epsilon\omega$ στροφη
= *geo strophe*
= Earth Turning

The term was coined by
Sir Napier Shaw,
Director of the Met Office.



IT: 2014-07-25_06:00
UCD MCC FORECAST VT: 2014-07-26_21:00

Wind Speed (kts)
Height (m) at 300 hPa
Wind (kts) at 300 hPa



Buys Ballot

**Christophorus Henricus
Diedericus Buys Ballot
(1817–1890)**

**Dutch meteorologist
and chemist
and mineralogist
and geologist
and mathematician.**



Buys Ballot's Law

In the Northern Hemisphere, if you stand with your back to the wind, the Low Pressure is to your left.



Buys Ballot's Law

In the Northern Hemisphere, if you stand with your back to the wind, the Low Pressure is to your left.

The GPS Version:

If you stand with your back to the wind, and the low pressure is to your left, then you must be in the Northern Hemisphere.



Buys Ballot's Law

In the Northern Hemisphere, if you stand with your back to the wind, **the Low Pressure is to your left.**

The GPS Version:

If you stand with your back to the wind, and the low pressure is to your left, **then you must be in the Northern Hemisphere.**



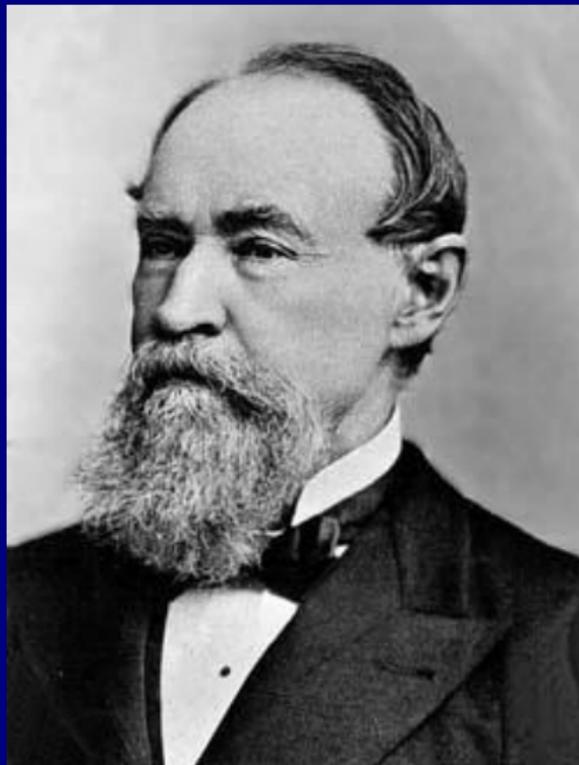
Raymond Hide



William Ferrel

**William Ferrel
(1817–1891)**

**American
meteorologist.**



Ferrel's 1856 Paper

An essay on the winds and the currents of the oceans.

[Nashville Journal of Medicine and Surgery, 1856.]



Ferrel's 1856 Paper

An essay on the winds and the currents of the oceans.

[Nashville Journal of Medicine and Surgery, 1856.]

“In consequence of the atmosphere’s revolving ... each particle is impressed with a **centrifugal force**.

“But if the rotatory motion of the atmosphere is greater than that of the Earth, this force is **increased**.

“and if ... [less] ... it is **diminished**.

“This difference gives rise to a **disturbing force** ... which **materially influences the motion**.”



Force Balance for Low and High Pressure

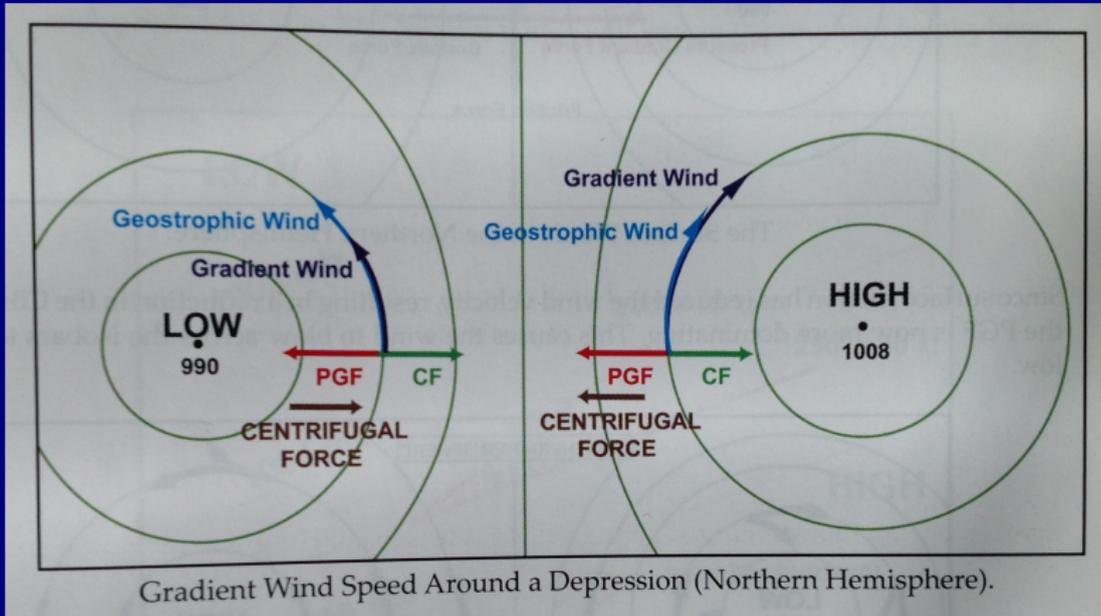


Image from ATPM Manual, Oxford Aviation Training

Gradient balance around low and high pressure.



Horizontal Equations of Motion

$$\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} + 2\boldsymbol{\Omega} \times \mathbf{V} + \frac{1}{\rho} \nabla p = \mathbf{0}$$



Horizontal Equations of Motion

$$\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} + 2\boldsymbol{\Omega} \times \mathbf{V} + \frac{1}{\rho} \nabla p = \mathbf{0}$$

For steady motion we get a three-way balance:

$$\underbrace{\mathbf{V} \cdot \nabla \mathbf{V}}_{\text{CFF}} + \underbrace{2\boldsymbol{\Omega} \times \mathbf{V}}_{\text{COR}} + \underbrace{(1/\rho) \nabla p}_{\text{PGF}} = \mathbf{0}$$



Horizontal Equations of Motion

$$\frac{\partial \mathbf{V}}{\partial t} + \mathbf{V} \cdot \nabla \mathbf{V} + 2\boldsymbol{\Omega} \times \mathbf{V} + \frac{1}{\rho} \nabla p = \mathbf{0}$$

For steady motion we get a three-way balance:

$$\underbrace{\mathbf{V} \cdot \nabla \mathbf{V}}_{\text{CFF}} + \underbrace{2\boldsymbol{\Omega} \times \mathbf{V}}_{\text{COR}} + \underbrace{(1/\rho) \nabla p}_{\text{PGF}} = \mathbf{0}$$

This is as easy as ABC:

$$A + B + C = 0$$



Three-way Balance

Also known as **Gradient Balance**:

$$\underbrace{\mathbf{V} \cdot \nabla \mathbf{V}}_{\text{CFF}} + \underbrace{2\boldsymbol{\Omega} \times \mathbf{V}}_{\text{COR}} + \underbrace{(1/\rho)\nabla p}_{\text{PGF}} = \mathbf{0}$$



Three-way Balance

Also known as **Gradient Balance**:

$$\underbrace{\mathbf{V} \cdot \nabla \mathbf{V}}_{CFF} + \underbrace{2\boldsymbol{\Omega} \times \mathbf{V}}_{COR} + \underbrace{(1/\rho)\nabla p}_{PGF} = \mathbf{0}$$

- ▶ *CFF* small \implies **Geostrophic Balance**
- ▶ *COR* small \implies **Cyclostrophic Balance**
- ▶ *PGF* small \implies **Inertial Balance**

Three for the price of one!



Geostrophic Balance

The **Coriolis Parameter** is $f = 2\Omega \sin \phi$.



Geostrophic Balance

The **Coriolis Parameter** is $f = 2\Omega \sin \phi$.

Then:

$$\underbrace{f \mathbf{k} \times \mathbf{V}}_{COR} + \underbrace{(1/\rho) \nabla p}_{PGF} = \mathbf{0}$$



Geostrophic Balance

The **Coriolis Parameter** is $f = 2\Omega \sin \phi$.

Then:

$$\underbrace{f\mathbf{k} \times \mathbf{V}}_{COR} + \underbrace{(1/\rho)\nabla p}_{PGF} = \mathbf{0}$$

**Balance between the Coriolis force
and the pressure gradient force:**

$$\mathbf{V}_{GEO} = \frac{1}{f\rho} \mathbf{k} \times \nabla p$$



Geostrophic Balance

The **Coriolis Parameter** is $f = 2\Omega \sin \phi$.

Then:

$$\underbrace{f\mathbf{k} \times \mathbf{V}}_{COR} + \underbrace{(1/\rho)\nabla p}_{PGF} = \mathbf{0}$$

**Balance between the Coriolis force
and the pressure gradient force:**

$$\mathbf{V}_{GEO} = \frac{1}{f\rho} \mathbf{k} \times \nabla p$$

We can determine the wind from the pressure!



Time-scale for Atmospheric Motions

Non-rotating Earth:

$$\underbrace{\frac{\partial \mathbf{V}}{\partial t}}_{V/T} + \underbrace{\frac{1}{\rho} \nabla p}_{\Delta p / \rho L} = \mathbf{0} \quad \text{or} \quad T = \frac{\rho L V}{\Delta p}$$

For typical synoptic values this gives $T \approx 3$ hours.



Time-scale for Atmospheric Motions

Non-rotating Earth:

$$\underbrace{\frac{\partial \mathbf{V}}{\partial t}}_{V/T} + \frac{1}{\underbrace{\rho}_{\Delta \rho / \rho L}} \nabla p = \mathbf{0} \quad \text{or} \quad T = \frac{\rho L V}{\Delta p}$$

For typical synoptic values this gives $T \approx 3$ hours.

Rotating Earth:

$$\underbrace{\frac{\partial \mathbf{V}}{\partial t}}_{V/T} + \underbrace{\mathbf{V} \cdot \nabla \mathbf{V}}_{V^2/L} = \mathbf{0} \quad \text{or} \quad T = \frac{L}{V}$$

For typical synoptic values this gives $T \approx 30$ hours.



Simple Geostrophic Adjustment

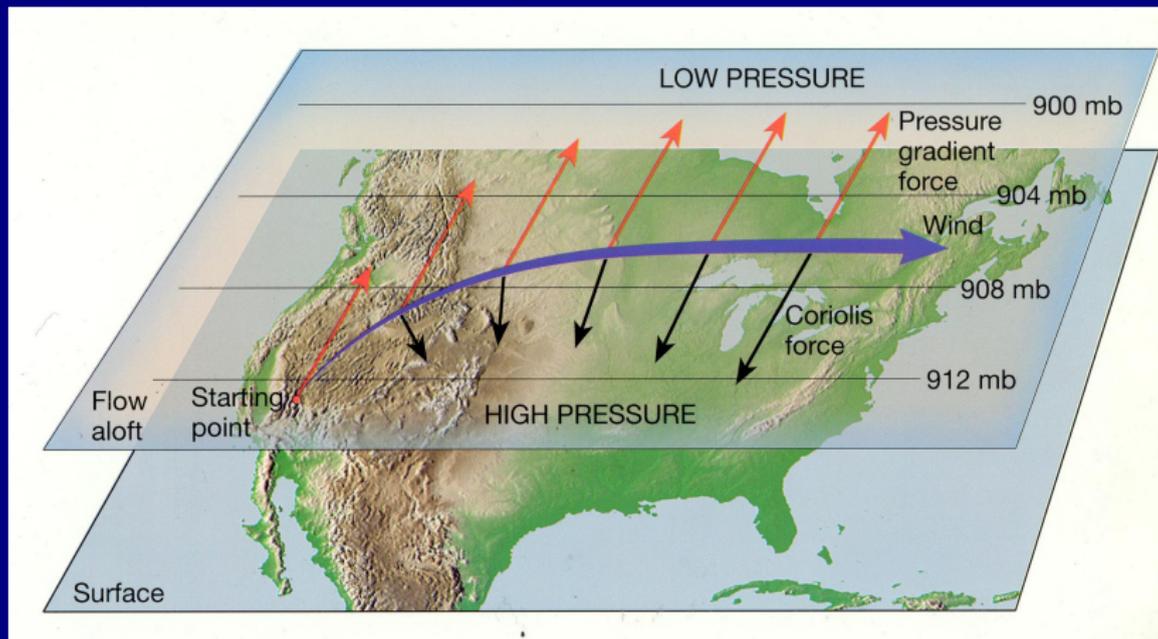


Image from Ackerman & Knox, © Jones & Bartlett Learning.



Examples of Balance in the Atmosphere

- ▶ **Hydrostatic balance**
- ▶ **Geostrophic balance**
- ▶ **Quasi-nondivergence**
- ▶ **Quasi-incompressibility**
- ▶ **Ocean atmosphere balance**
- ▶ **Energy balance**
- ▶ **Ice sheet balance**
- ▶ **Etc., etc., etc.**



Geostrophic Flow is Quasi-nondivergent

$$\mathbf{V}_{\text{GEO}} = \frac{1}{f\rho} \mathbf{k} \times \nabla \rho$$

Ignore variations in f and ρ :

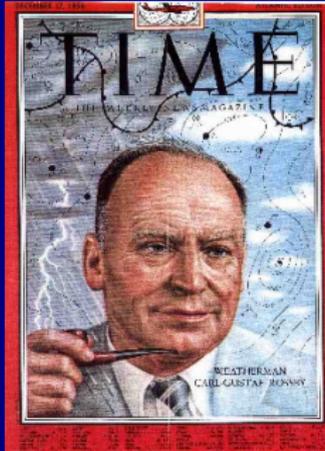
$$\mathbf{V}_{\text{GEO}} = \mathbf{k} \times \nabla \left(\frac{\rho}{f\rho} \right) = \nabla \times \left(-\frac{\rho}{f\rho} \right) \mathbf{k}$$

Divergence of a curl is zero:

$$\nabla \cdot \mathbf{V}_{\text{GEO}} = 0$$



The Rossby Number



C. G. Rossby in *Time*

$$Ro = \frac{\text{Centrifugal Force}}{\text{Coriolis Force}} = \frac{V}{fL}$$

$$Ro = \frac{\text{Spin of the Flow}}{\text{Spin of the Earth}} = \frac{\zeta}{f}$$



500 mb geopotential and wind field

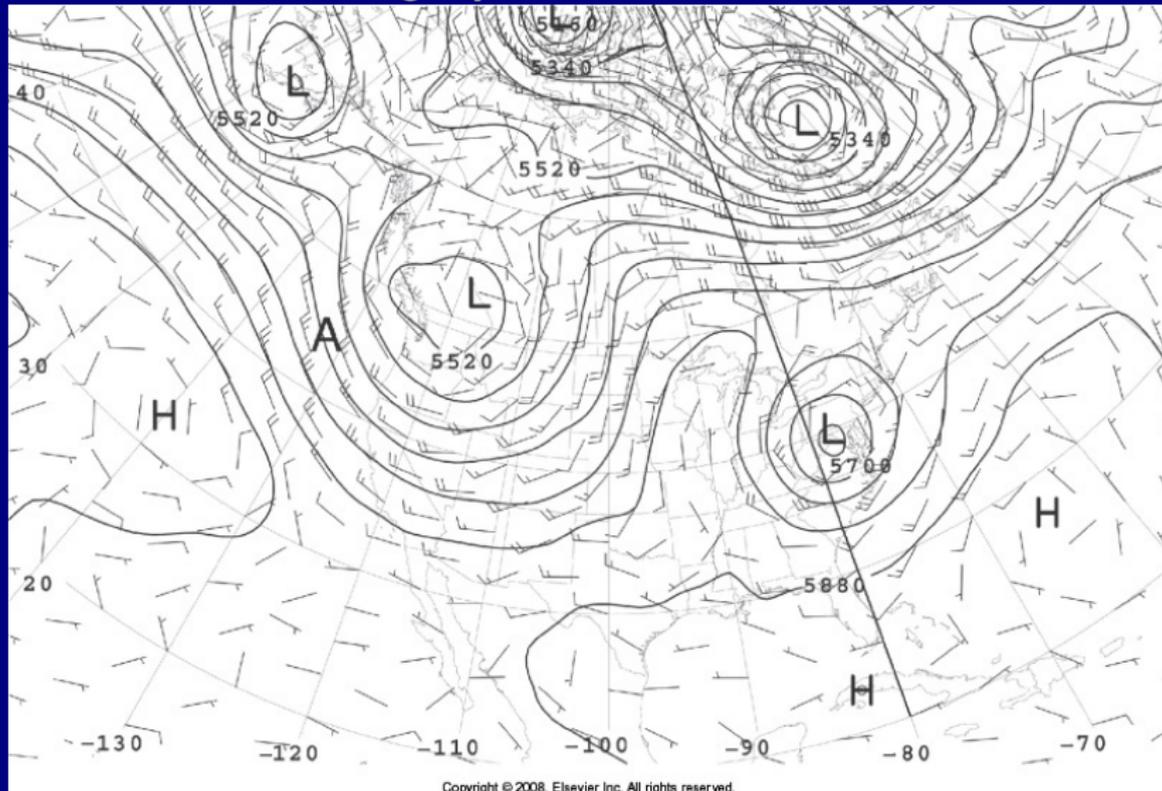


Image from Marshall & Plumb, ©Elsevier.



500 mb Rossby Number $|V \cdot \nabla V| / fV$

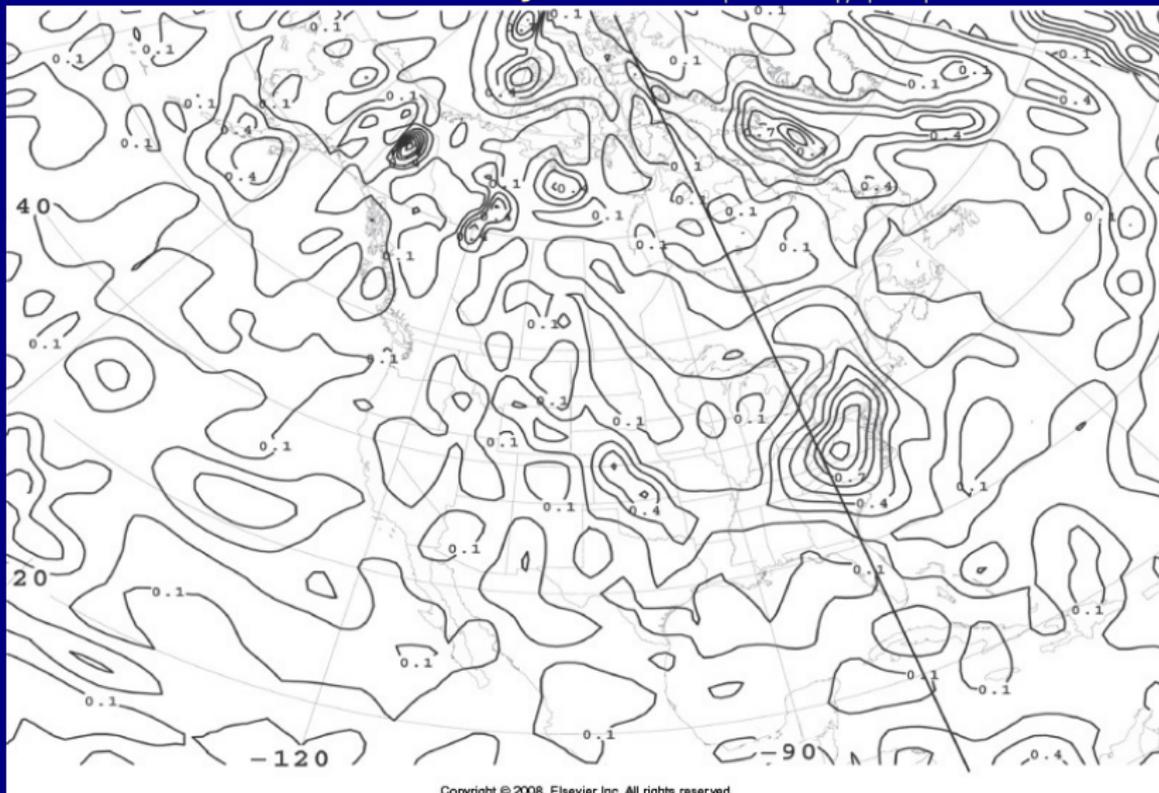


Image from Marshall & Plumb, ©Elsevier.



Balance at Different Scales

- ▶ **Extra-tropical Depressions**
- ▶ **Tropical Cyclones**
- ▶ **Tornadoes**
- ▶ **Domestic.**



Balance at Different Scales: Depressions



Extra-tropical Depression

$$Ro \approx \frac{1}{10}$$

Geostrophic Balance Good

Gradient Balance Better.

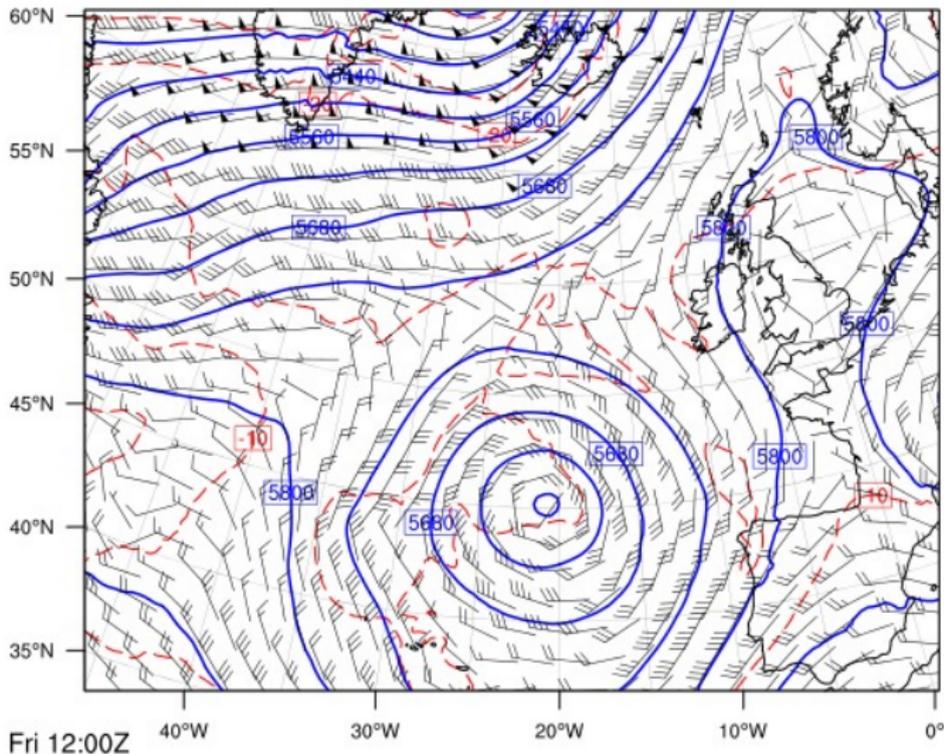


Temperature (C) at 500 hPa
Height (m) at 500 hPa
Wind (kts) at 500 hPa

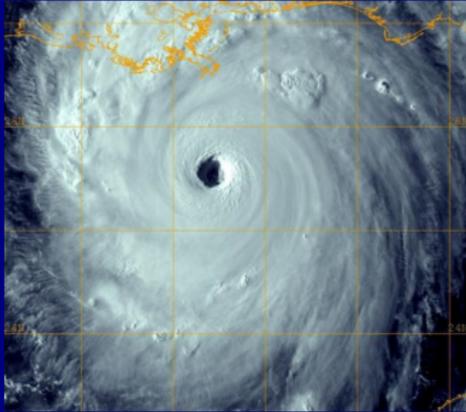
UCD MCC FORECAST

IT: 2014-09-12_00:00

VT: 2014-09-12_12:00



Balance at Different Scales: Tropical Cyclones



Hurricane

$Ro \approx 10 - 100$

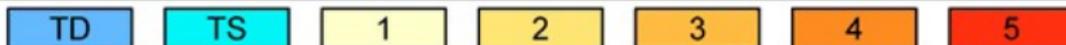
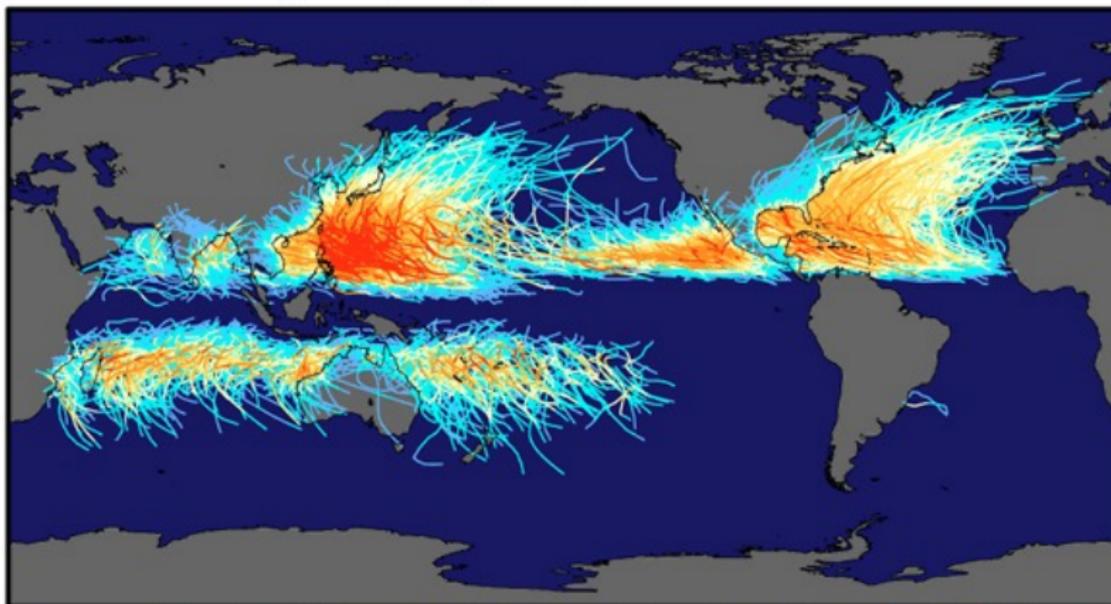
Geostrophic Balance Bad

Gradient Balance Better.



Tropical Cyclone Tracks

Tracks and Intensity of Tropical Cyclones, 1851-2006

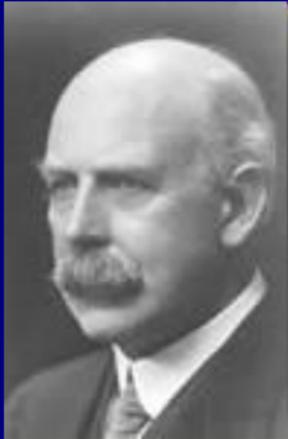


Saffir-Simpson Hurricane Intensity Scale

NASA



Distribution of Tropical Cyclones



“I always find my pen sticks to the paper and refuses to move when I try to draw an isobar across the equator.”

Napier Shaw (1923): *The air and its ways*. CUP, pg. 51.



Balance at Different Scales: Tornadoes



Tornado

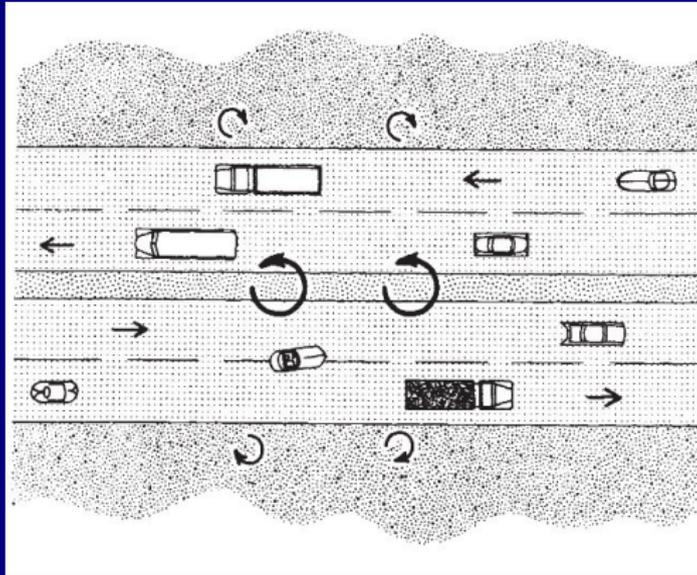
$Ro \approx 10,000$

Close to Cyclostrophic Balance

Coriolis effect influences background flow



Traffic Flow and Vorticity



Effect of vorticity pollution by motor vehicles on tornadoes.
Isaacs, J. D., J. W. Stork, D. B. Goldstein & G. L. Wick
Nature, 253, 254–255 (1975).



> 98% of tornadoes are cyclonic, **but ...**



> 98% of tornadoes are cyclonic, **but ...**



WIKIPEDIA
The Free Encyclopedia

Article **Talk** Read

Anticyclonic tornado

From Wikipedia, the free encyclopedia

▶ **Notable Anticyclonic Tornadoes:**

- ▶ West Bend tornado
- ▶ Grand Island tornado
- ▶ Woodward, Oklahoma April 10th 2012
- ▶ Aurora Nebraska, 2009
- ▶ Freedom, Oklahoma, June 6, 1975
- ▶ Sunnyvale, California, May 4, 1998
- ▶ El Reno, Oklahoma, May 31, 2013

▶ **Anticyclonic tornadoes rotate clockwise (in NHS)**



Balance at Different Scales: Domestic



Down the Plughole

$$Ro \approx 100,000$$

Cyclostrophic Balance



Balance at Different Scales: Domestic



Down the Plughole

$$Ro \approx 100,000$$

Cyclostrophic Balance

Coriolis Effect Completely Irrelevant

... unless you believe Homer Simpson



Review of Dynamical Balance

When the forces acting on a parcel sum to zero, a balance is achieved.

With balance, there is **steady flow**.

- ▶ **Hydrostatic Balance**
- ▶ **Geostrophic Balance**
- ▶ **Gradient Balance**
- ▶ **Cyclostrophic Balance**



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



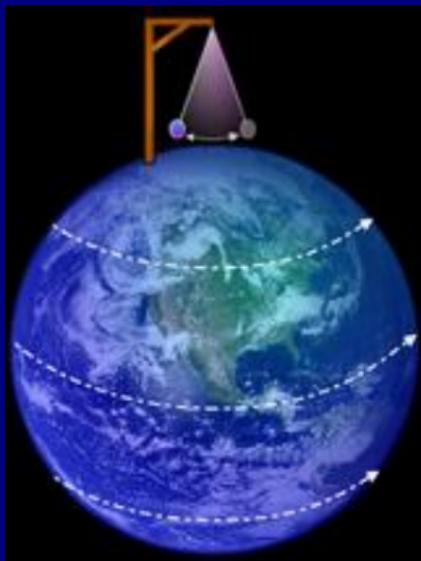
Foucault's Pendulum

Foucault's pendulum experiment in 1851 was the first simple terrestrial demonstration of the rotation of the Earth.

Not entirely true: Laplace, Gauss and falling objects.



Foucault's Pendulum

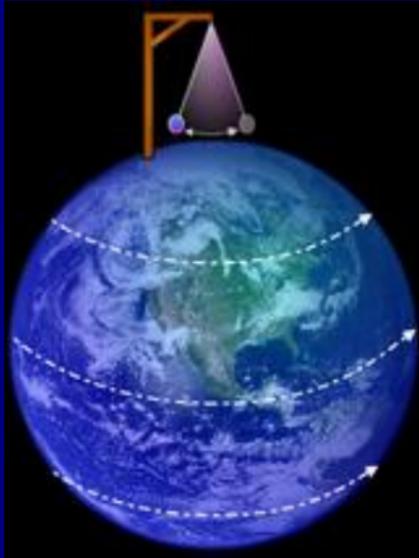


Basic Idea:

The pendulum swings in a fixed plane while the Earth spins beneath it.



Foucault's Pendulum



Basic Idea:

The pendulum swings in a fixed plane while the Earth spins beneath it.

Things are not so simple!



Foucault's Pendulum at the Panthéon



Intro

Balance

Foucault

Coriolis

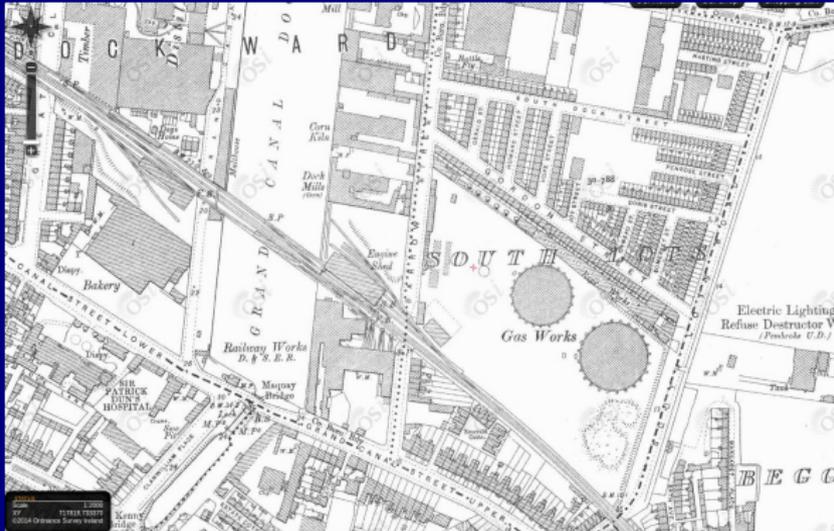
LFR

ENIAC

Foucault's Pendulum in Dublin

Foucault's experiment triggered **pendulum mania**.

The experiment was repeated in Dublin in 1851 by Joseph Galbraith and Samuel Haughton of TCD.



Foucault's Pendulum in Dublin

Galbraith and Haughton published a quite complete mathematical analysis in the *Philosophical Magazine*.

Another analysis by Matthew O'Brien, was one of the first applications of vector analysis.



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



Newtonian Mechanics

Newtonian mechanics assumes the existence of an absolute, unaccelerated frame of reference.

Newton's laws are **covariant in all inertial frames.**

They keep the same mathematical form under Galilean transformations.

They are **not covariant in accelerating frames:** there are additional terms.



Newtonian Mechanics

The second law of motion in vector form is

$$\frac{d\mathbf{p}}{dt} = \mathbf{F}$$

This equation is valid in all inertial frames.



Newtonian Mechanics

The second law of motion in vector form is

$$\frac{d\mathbf{p}}{dt} = \mathbf{F}$$

This equation is valid in all inertial frames.

However, the **component form** of the equation,

$$\frac{dp_i}{dt} = F_i$$

is true only for cartesian coordinates.



Newtonian Mechanics

In cartesian coordinates in two dimensions:

$$\frac{d^2x}{dt^2} = \frac{F_x}{m}$$

$$\frac{d^2y}{dt^2} = \frac{F_y}{m}$$



Newtonian Mechanics

In cartesian coordinates in two dimensions:

$$\frac{d^2x}{dt^2} = \frac{F_x}{m} \qquad \frac{d^2y}{dt^2} = \frac{F_y}{m}$$

In polar coordinates (r, ϕ) additional terms appear:

$$\frac{d^2r}{dt^2} - r \left(\frac{d\phi}{dt} \right)^2 = \frac{F_r}{m},$$
$$r \frac{d^2\phi}{dt^2} + 2 \frac{dr}{dt} \frac{d\phi}{dt} = \frac{F_\phi}{m}$$

The equations are **not covariant**.



Velocity in Rotating Frame

For a point \mathbf{x}' fixed in a rotating frame:

$$\mathbf{v} = \boldsymbol{\Omega} \times \mathbf{x}'$$



Velocity in Rotating Frame

For a point \mathbf{x}' fixed in a rotating frame:

$$\mathbf{v} = \boldsymbol{\Omega} \times \mathbf{x}'$$

The vector product is the root of the difficulty in understanding the Coriolis effect.



Velocity in Rotating Frame

For a point \mathbf{x}' fixed in a rotating frame:

$$\mathbf{v} = \boldsymbol{\Omega} \times \mathbf{x}'$$

The vector product is the root of the difficulty in understanding the Coriolis effect.

For a particle with velocity \mathbf{v}' in the rotating frame,

$$\underbrace{\mathbf{v}}_{\text{ABS}} = \underbrace{\mathbf{v}'}_{\text{REL}} + \underbrace{\boldsymbol{\Omega} \times \mathbf{x}'}_{\text{FRAME}}$$

We just add the two contributions to velocity.



O'Brien's Equation

- ▶ Let \mathbf{A} be a vector in an inertial frame
- ▶ \mathbf{A}' the same vector in a frame with rotation Ω .

The rates of change are related:

$$\frac{d\mathbf{A}}{dt} = \frac{d\mathbf{A}'}{dt} + \Omega \times \mathbf{A}'$$



Matthew O'Brien (1814-1855)



O'Brien's Equation

- ▶ Let \mathbf{A} be a vector in an inertial frame
- ▶ \mathbf{A}' the same vector in a frame with rotation Ω .



Matthew O'Brien (1814-1855)

The rates of change are related:

$$\frac{d\mathbf{A}}{dt} = \frac{d\mathbf{A}'}{dt} + \Omega \times \mathbf{A}'$$

This expression is fundamental. It was first expressed in vector form by Matthew O'Brien.

I propose to call it O'Brien's equation.

Paper to appear in *Bulletin of Irish Mathematical Society*.



Applying O'Brien's equation to the position vectors,

$$\frac{d\mathbf{x}}{dt} = \frac{d\mathbf{x}'}{dt} + \boldsymbol{\Omega} \times \mathbf{x}' ,$$

or

$$\mathbf{v} = \mathbf{v}' + \boldsymbol{\Omega} \times \mathbf{x}' .$$



Applying O'Brien's equation to the position vectors,

$$\frac{d\mathbf{x}}{dt} = \frac{d\mathbf{x}'}{dt} + \boldsymbol{\Omega} \times \mathbf{x}' ,$$

or

$$\mathbf{v} = \mathbf{v}' + \boldsymbol{\Omega} \times \mathbf{x}' .$$

Now applying the relationship again

$$\frac{d\mathbf{v}}{dt} = \frac{d\mathbf{v}'}{dt} + \underbrace{2\boldsymbol{\Omega} \times \mathbf{v}'}_{\text{COR}} + \underbrace{\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{x}')}_{\text{CFF}} + \underbrace{\dot{\boldsymbol{\Omega}} \times \mathbf{x}'}_{\text{EUL}}$$



Applying O'Brien's equation to the position vectors,

$$\frac{d\mathbf{x}}{dt} = \frac{d\mathbf{x}'}{dt} + \boldsymbol{\Omega} \times \mathbf{x}' ,$$

or

$$\mathbf{v} = \mathbf{v}' + \boldsymbol{\Omega} \times \mathbf{x}' .$$

Now applying the relationship again

$$\frac{d\mathbf{v}}{dt} = \frac{d\mathbf{v}'}{dt} + \underbrace{2\boldsymbol{\Omega} \times \mathbf{v}'}_{\text{COR}} + \underbrace{\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{x}')}_{\text{CFF}} + \underbrace{\dot{\boldsymbol{\Omega}} \times \mathbf{x}'}_{\text{EUL}}$$

The acceleration has three additional terms:

- ▶ The Coriolis acceleration $2\boldsymbol{\Omega} \times \mathbf{v}'$
- ▶ The centrifugal acceleration $\boldsymbol{\Omega} \times (\boldsymbol{\Omega} \times \mathbf{x}')$
- ▶ The Euler term $\dot{\boldsymbol{\Omega}} \times \mathbf{x}'$.



Transforming the Equations

Assume Ω constant ($\dot{\Omega} = 0$) and drop the Euler term.



Transforming the Equations

Assume Ω constant ($\dot{\Omega} = 0$) and drop the Euler term.

Newton's equation may then be written

$$m \frac{d\mathbf{v}'}{dt} = \mathbf{F}' - \underbrace{2m\Omega \times \mathbf{v}'}_{\text{COR}} - \underbrace{m\Omega \times (\Omega \times \mathbf{x}')}_{\text{CFF}}$$

where \mathbf{F}' is the physical force in the rotating frame.

The two additional terms now appear as forces.



Covariant form of Newton's equations

We can express Newton's equations so that they are **covariant under rotations**.

We define a new time derivative

$$\frac{D\mathbf{A}}{Dt} \equiv \frac{d\mathbf{A}}{dt} + \boldsymbol{\Omega} \times \mathbf{A}$$



Covariant form of Newton's equations

We can express Newton's equations so that they are **covariant under rotations**.

We define a new time derivative

$$\frac{D\mathbf{A}}{Dt} \equiv \frac{d\mathbf{A}}{dt} + \boldsymbol{\Omega} \times \mathbf{A}$$

We then write the equation of motion as

$$\mathbf{p} = m \frac{D\mathbf{x}}{Dt} \quad \frac{D\mathbf{p}}{Dt} = \mathbf{F}.$$

These equations keep the same mathematical form under all rotational transformations.



Lagrange's Equations

We define the **Lagrangian**:

$$L = \left[\begin{array}{c} \text{Kinetic} \\ \text{Energy} \end{array} \right] - \left[\begin{array}{c} \text{Potential} \\ \text{Energy} \end{array} \right]$$



Lagrange's Equations

We define the **Lagrangian**:

$$L = \left[\begin{array}{c} \text{Kinetic} \\ \text{Energy} \end{array} \right] - \left[\begin{array}{c} \text{Potential} \\ \text{Energy} \end{array} \right]$$

Then Lagrange's equation of motion are

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{q}_\rho} = \frac{\partial L}{\partial q_\rho}$$

These equations are in covariant form:
They are valid in all frames of reference.



The Principle of Relativity

The equations expressing the laws of physics have
the same form in all admissible frames of reference.



The Principle of Relativity

The equations expressing the laws of physics have **the same form in all admissible frames of reference.**

The Coriolis effect arises through rotation of the reference frame.

Can we use the *Principle of Relativity* to obtain the Coriolis terms?



The Principle of Relativity

The equations expressing the laws of physics have **the same form in all admissible frames of reference.**

The Coriolis effect arises through rotation of the reference frame.

Can we use the *Principle of Relativity* to obtain the Coriolis terms?

Yes!

Warning: Not quite as easy as A, B, C!



Tensorial Formulation of Equations

Three very recent papers in the *Quarterly Journal of the Royal Meteorological Society*:

Charron, Martin, Ayrton Zadra, and Claude Girard, 2014: Four-dimensional tensor equations for a classical fluid in an external gravitational field.

Quart. J. Roy. Met. Soc. 140 (680), 908–916.

Fundamental equations in tensor form:

$$T^{\mu\nu}{}_{;\nu} = -\rho h^{\mu\nu} \Phi_{,\nu}$$

where the mass-momentum-stress tensor is

$$T^{\mu\nu} = \rho u^\mu u^\nu + h^{\mu\nu} p + \sigma^{\mu\nu}$$



The General Geodesic Equation

In an inertial frame with cartesian coordinates,

$$ds^2 = dx^2 + dy^2 = g_{\mu\nu} dx^\mu dx^\nu$$

The line element is invariant.



The General Geodesic Equation

In an inertial frame with cartesian coordinates,

$$ds^2 = dx^2 + dy^2 = g_{\mu\nu} dx^\mu dx^\nu$$

The line element is invariant.

The rotating coordinates (X, Y) are

$$X = \cos \Omega t x + \sin \Omega t y$$

$$Y = -\sin \Omega t x + \cos \Omega t y$$

In the rotating frame

$$ds^2 = dX^2 + dY^2 - 2\Omega dx dT + 2\Omega dX dT + \Omega^2 (X^2 + Y^2) dT^2$$



We write this as

$$ds^2 = g'_{\mu\nu} dX^\mu dX^\nu$$

where the metric tensor is

$$g'_{\mu\nu} = \begin{bmatrix} 1 & 0 & -\Omega Y \\ 0 & 1 & \Omega X \\ -\Omega Y & \Omega X & \Omega^2(X^2 + Y^2) \end{bmatrix}$$

Note that $g'_{\mu\nu}$ is **singular**: inverse $g'^{\mu\nu}$ does not exist.



The geodesic equation is:

$$\frac{d}{dt} \left(g'_{\sigma\nu} \frac{dX^\nu}{dt} \right) - \frac{1}{2} \frac{\partial g'_{\mu\nu}}{\partial X^\sigma} \frac{dX^\mu}{dt} \frac{dX^\nu}{dt} = 0$$



The geodesic equation is:

$$\frac{d}{dt} \left(g'_{\sigma\nu} \frac{dX^\nu}{dt} \right) - \frac{1}{2} \frac{\partial g'_{\mu\nu}}{\partial X^\sigma} \frac{dX^\mu}{dt} \frac{dX^\nu}{dt} = 0$$

Writing this explicitly, we get

$$\begin{aligned} \ddot{X} - 2\Omega \dot{Y} - \Omega^2 X &= 0 \\ \ddot{Y} + 2\Omega \dot{X} - \Omega^2 Y &= 0 \end{aligned}$$

These are the equations derived already by more conventional means.



Why use the Tensor Formulation?

- ▶ Tensor equations are **covariant**: they preserve their form in all coordinate systems;
- ▶ **Transformations** are handled systematically;
- ▶ **Approximations** are derived rigourously;
- ▶ **Conservation** properties are preserved.



An alternative equation for the geodesics is

$$\frac{d^2 X^\rho}{ds^2} + \Gamma^\rho_{\mu\nu} \frac{dX^\mu}{ds} \frac{dX^\nu}{ds} = 0$$

The Christoffel symbols of the first kind are

$$[\sigma|\mu\nu] = \Gamma_{\sigma|\mu\nu} = \frac{1}{2} \left[\frac{\partial g'_{\sigma\nu}}{\partial X^\mu} + \frac{\partial g'_{\mu\sigma}}{\partial X^\nu} - \frac{\partial g'_{\mu\nu}}{\partial X^\sigma} \right]$$

There are ten non-vanishing symbols:

$$\begin{aligned} [1, 33] &= -\Omega^2 X & [2, 33] &= -\Omega^2 Y \\ [1, 23] &= [1, 32] = -\Omega & [2, 13] &= [2, 31] = +\Omega \\ [3, 13] &= [3, 31] = \Omega^2 X & [3, 23] &= [3, 32] = \Omega^2 Y \end{aligned}$$

where the variables are $(X^1, X^2, X^3) = (X, Y, T)$.



The Christoffel symbols of the second kind are

$$\Gamma^{\rho}{}_{\mu\nu} = g^{\rho\sigma}\Gamma_{\sigma|\mu\nu}$$

To regularise $g_{\mu\nu}$, we write the metric as

$$ds^2 = dx^2 + dy^2 + \epsilon dt^2$$

and consider the limiting case $\epsilon \rightarrow 0$.

The $\Gamma^{\rho}{}_{\mu\nu}$ are independent of ϵ . The non-zero ones are

$$\begin{aligned} \Gamma^1{}_{23} = \Gamma^1{}_{32} = -\Omega & & \Gamma^2{}_{13} = \Gamma^2{}_{31} = +\Omega \\ \Gamma^1{}_{33} = -\Omega^2 X & & \Gamma^2{}_{33} = -\Omega^2 Y \end{aligned}$$

These yield the same equations as obtained above.

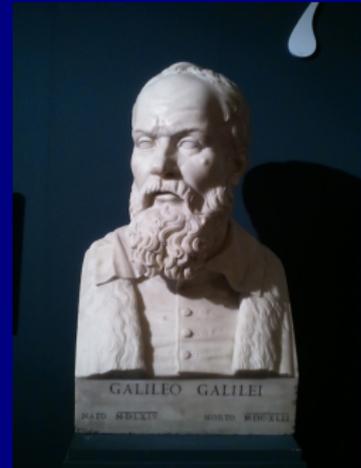
The curvature tensor vanishes: $R^{\rho}{}_{\sigma\mu\nu} \equiv 0$.



Galileo on Mathematics

[The universe] ... is written in the language of mathematics ... without which it is ... impossible to understand a single word of it.

Without this understanding, one is wandering around in a dark labyrinth.



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

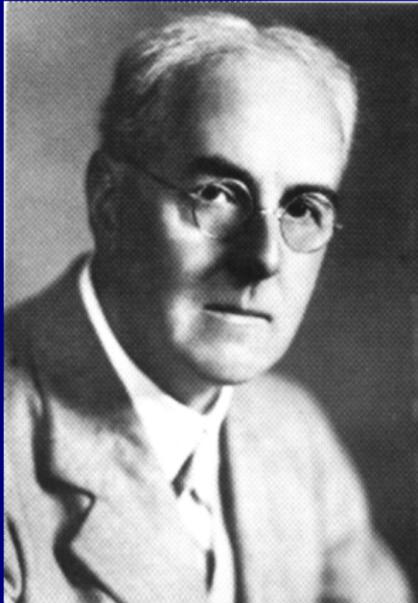
Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



Lewis Fry Richardson, 1881–1953.



During WWI, Richardson computed the pressure change at a single point.

It took him **two years** !

His ‘forecast’ was a catastrophic failure:

$$\Delta p = 145 \text{ hPa in 6 hrs}$$

But Richardson’s **method** was scientifically sound.

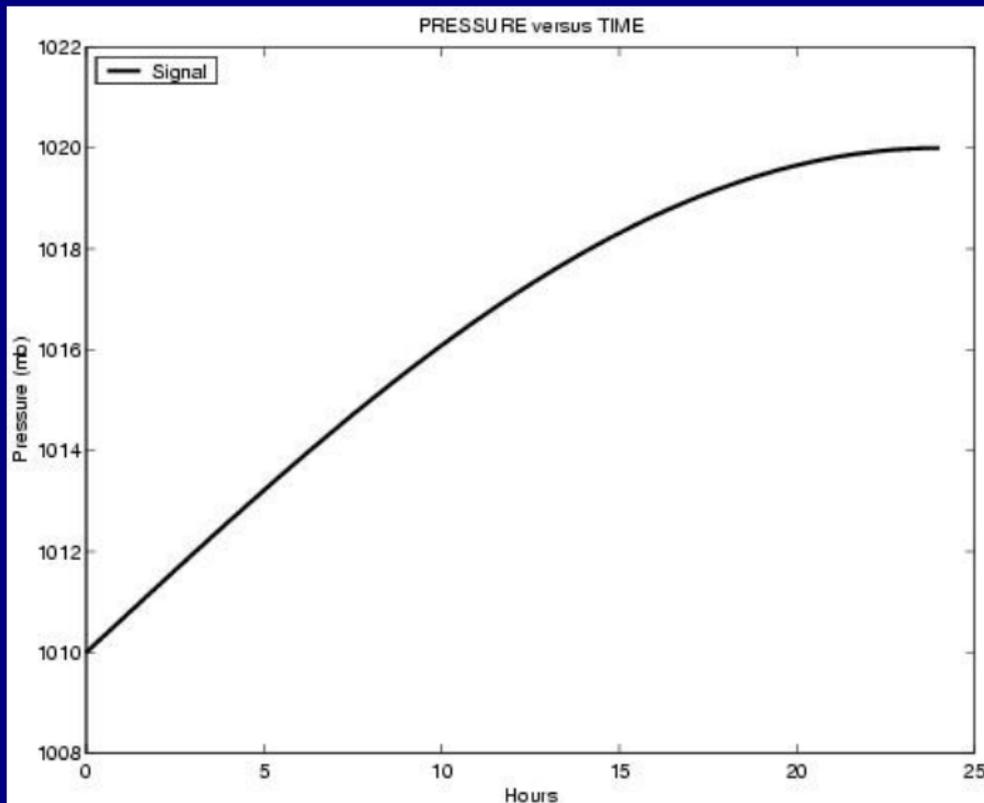


A Richardsonian Limerick

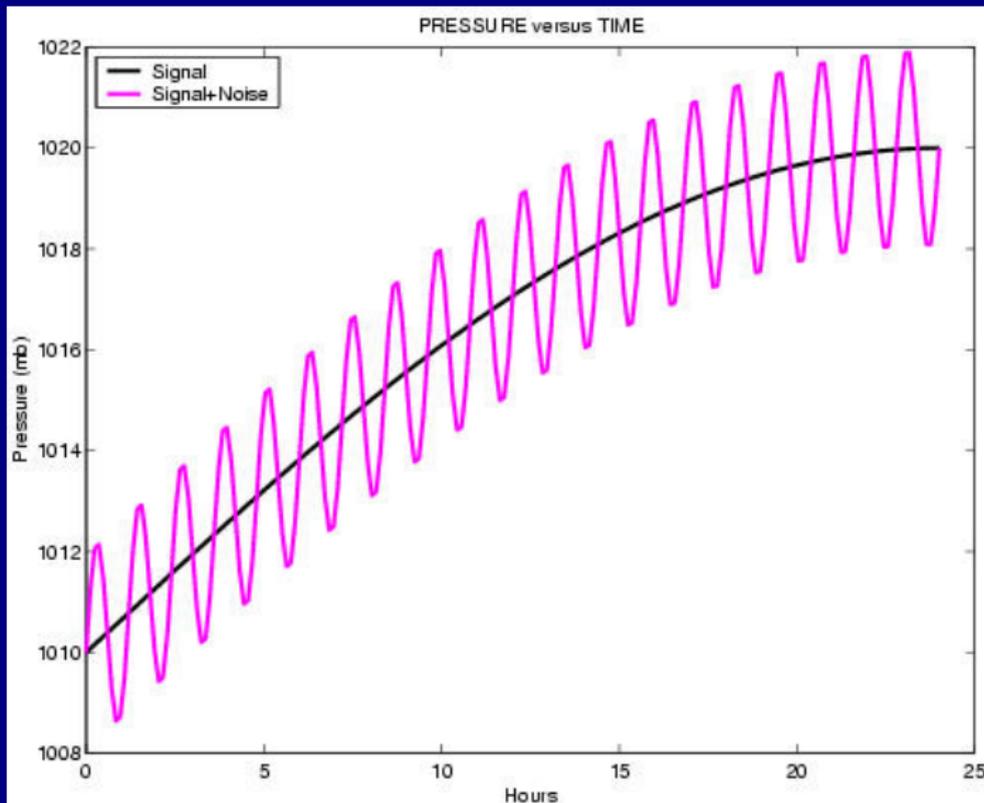
*Young Richardson wanted to know
How quickly the pressure would grow.
But, what a surprise, 'cos
The six-hourly rise was,
In Pascals, One Four Five — Oh Oh!*



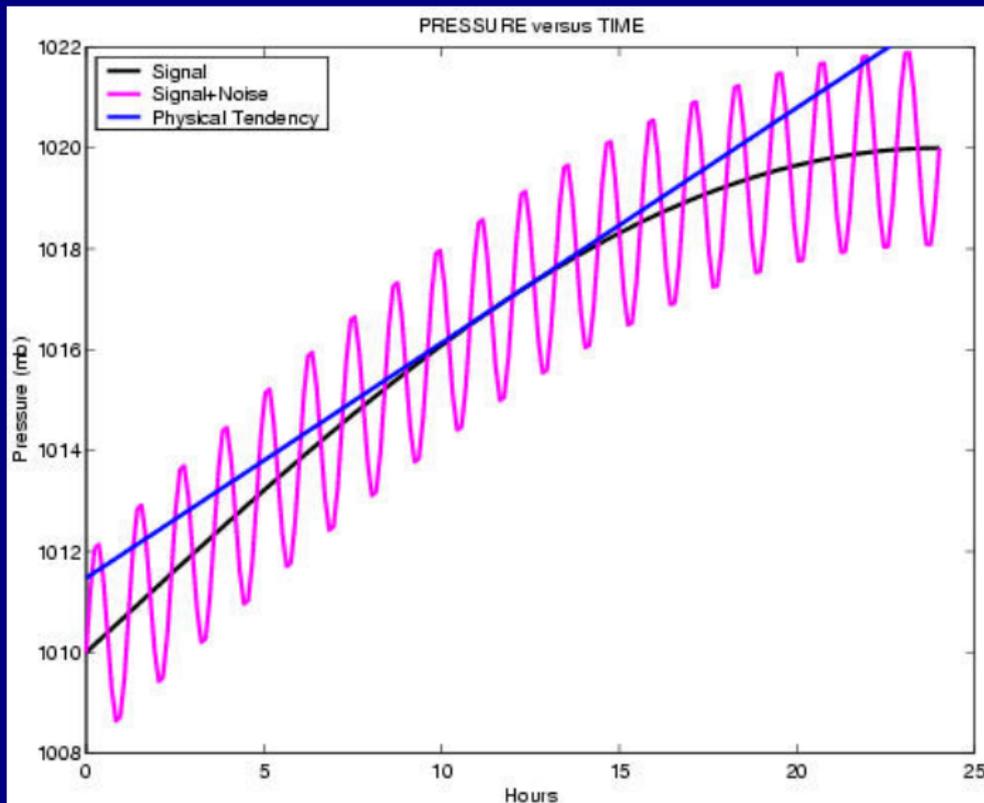
Smooth Evolution of Pressure



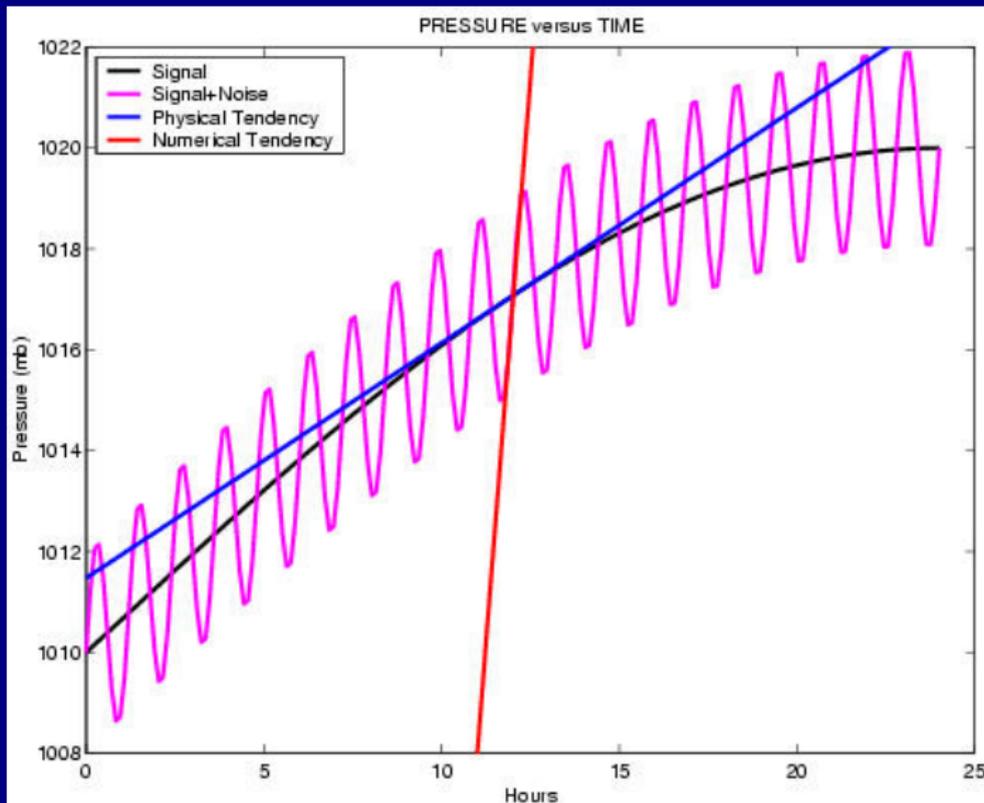
Noisy Evolution of Pressure



Tendency of a Smooth Signal



Tendency of a Noisy Signal



Initialization of Richardson's Forecast

Richardson's Forecast has been repeated.

**The atmospheric observations for 20 May, 1910
*were recovered from original sources.***



Initialization of Richardson's Forecast

Richardson's Forecast has been repeated.

**The atmospheric observations for 20 May, 1910
*were recovered from original sources.***

- ▶ **ORIGINAL:** $\frac{dp_s}{dt} = +145 \text{ hPa}/6 \text{ h}$
- ▶ **INITIALIZED:** $\frac{dp_s}{dt} = -0.9 \text{ hPa}/6 \text{ h}$

Observations: The barometer was steady!



Initialization of Richardson's Forecast

Richardson's Forecast has been repeated.

The atmospheric observations for 20 May, 1910 *were recovered from original sources.*

- ▶ **ORIGINAL:** $\frac{dp_s}{dt} = +145 \text{ hPa}/6 \text{ h}$
- ▶ **INITIALIZED:** $\frac{dp_s}{dt} = -0.9 \text{ hPa}/6 \text{ h}$

Observations: **The barometer was steady!**

BALANCED INITIAL DATA IS ESSENTIAL!



Outline

Introduction

Atmospheric Balance

Foucault Pendulum

Coriolis Effect

Richardson's Forecast

The ENIAC Forecasts



The ENIAC Forecasts



ENIAC: The first multi-purpose programmable electronic digital computer.

- ▶ 18,000 vacuum tubes
- ▶ 70,000 resistors
- ▶ 10,000 capacitors
- ▶ 6,000 switches
- ▶ Power: 140 kWatts



Charney, et al., *Tellus*, 1950.

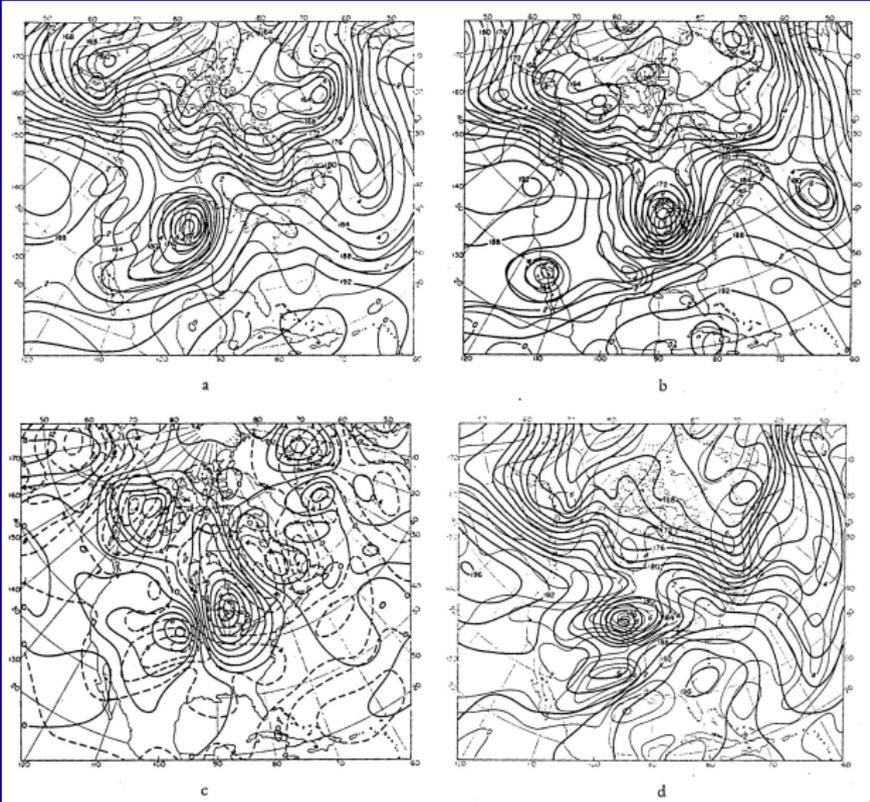
- ▶ The atmosphere is treated as a single layer.
- ▶ The flow is assumed to be nondivergent.
- ▶ Absolute vorticity is conserved.

$$\frac{d(\zeta + f)}{dt} = 0.$$

$$\left[\begin{array}{c} \text{Absolute} \\ \text{Vorticity} \end{array} \right] = \left[\begin{array}{c} \text{Relative} \\ \text{Vorticity} \end{array} \right] + \left[\begin{array}{c} \text{Planetary} \\ \text{Vorticity} \end{array} \right].$$



ENIAC Forecast for Jan 5, 1949



Recreating the ENIAC Forecasts

The ENIAC integrations have been repeated using:

- ▶ A **MATLAB** program to solve the BVE
- ▶ Data from the NCEP/NCAR reanalysis

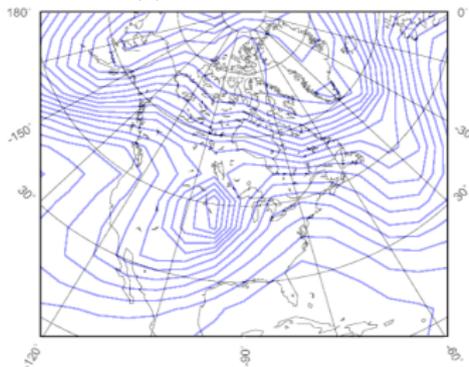
The matlab code is available on the website

<http://maths.ucd.ie/~plynch/eniac>

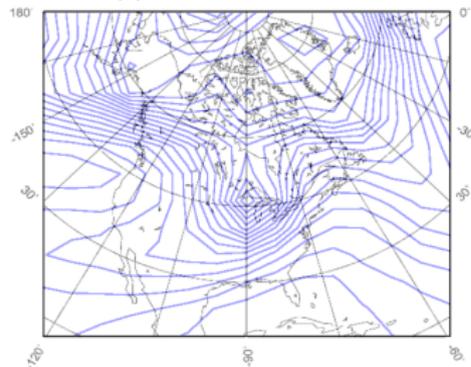


Recreation of the Forecast

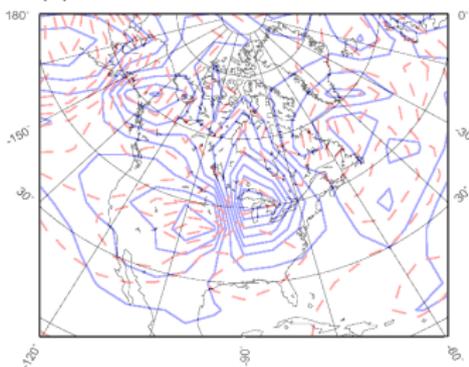
(A) INITIAL ANALYSIS



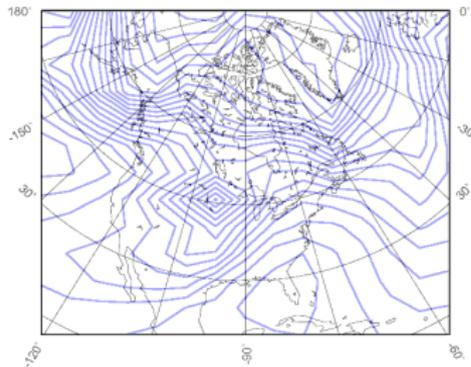
(B) VERIFYING ANALYSIS



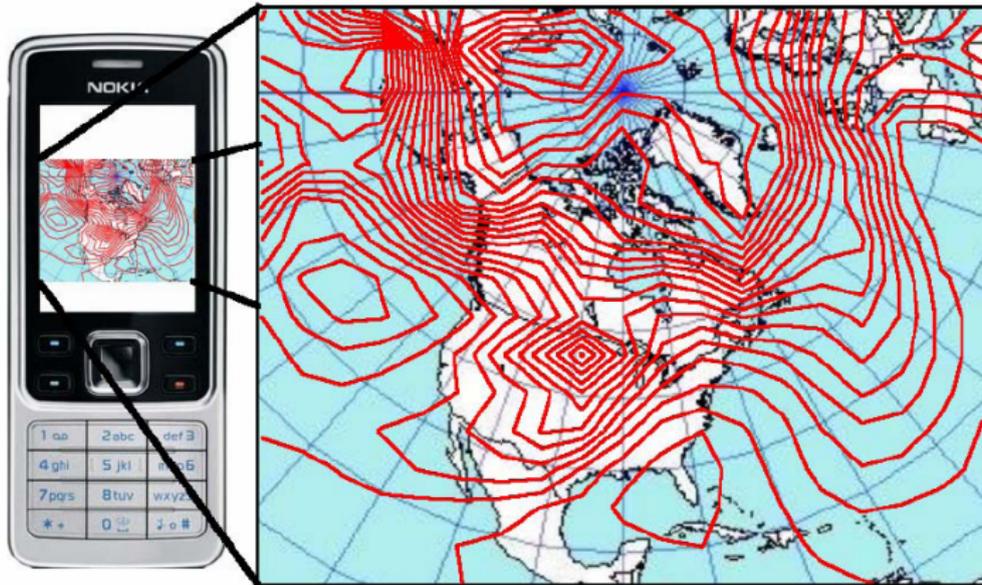
(C) ANALYSED & FORECAST CHANGES



(D) FORECAST HEIGHT



PHONIAC: Portable Hand Operated Numerical Integrator and Computer



Forecasts by PHONIAc

Peter Lynch¹
and Owen Lynch²

¹University College Dublin, Meteorology
and Climate Centre, Dublin

²Dublin Software Laboratory, IBM Ireland

The first computer weather forecasts were made in 1950, using the ENIAC (Electronic Numerical Integrator and Computer). The ENIAC forecasts led to operational numerical weather prediction within five years, and paved the way for the remarkable advances in weather prediction and climate modelling that have been made over the past half century. The basis for the forecasts was the barotropic vorticity equation (BVE). In the present study, we describe the solution of the BVE on a mobile phone (cell-phone), and repeat one of the ENIAC forecasts. We speculate on the possible applications of mobile phones for micro-scale numerical weather prediction.

The ENIAC Integrations

and John von Neumann (1950; cited below as CFvN). The story of this work was recounted by George Platzman in his Victor P. Starr Memorial Lecture (Platzman, 1979). The atmosphere was treated as a single layer, represented by conditions at the 500 hPa level, modelled by the BVE. This equation, expressing the conservation of absolute vorticity following the flow, gives the rate of change of the Laplacian of height in terms of the advection. The tendency of the height field is obtained by solving a Poisson equation with homogeneous boundary conditions. The height field may then be advanced to the next time level. With a one hour time-step, this cycle is repeated 24 times for a one-day forecast.

The initial data for the forecasts were prepared manually from standard operational 500 hPa analysis charts of the U.S. Weather Bureau, discretised to a grid of 19 by 16 points, with grid interval of 736 km. Centred spatial finite differences and a leapfrog time-scheme were used. The boundary conditions for height were held constant throughout each 24-hour integration. The forecast starting at 0300 UTC, January 5, 1949 is shown in

vorticity. The forecast height and vorticity are shown in the right panel. The feature of primary interest was an intense depression over the United States. This deepened, moving NE to the 90°W meridian in 24 hours. A discussion of this forecast, which underestimated the development of the depression, may be found in CFvN and in Lynch (2008).

Dramatic growth in computing power

The oft-cited paper in *Tellus* (CFvN) gives a complete account of the computational algorithm and discusses four forecast cases. The ENIAC, which had been completed in 1945, was the first programmable electronic digital computer ever built. It was a gigantic machine, with 18,000 thermionic valves, filling a large room and consuming 140 kW of power. Input and output was by means of punch-cards. McCartney (1999) provides an absorbing account of the origins, design, development and destiny of ENIAC.

Advances in computer technology over the past half-century have been spectacular. The increase in computing power is encap-

Notices of the (other) AMS



September 2013 issue of
*Notices of the American
Mathematical Society.*



Thank you

