

§3.2. Initial Value Problems

Hyperbolic and parabolic PDEs are **initial value problems** or **marching problems** (a term introduced by Richardson). The solution is obtained by using the known initial values and marching or advancing in time.

If boundary values are necessary, they are called **mixed initial-boundary value problems**.

The simplest prototypes of these initial value problems are:

- The **advection equation** (with solution $u(x, t) = u(x - ct, 0)$)

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0,$$

which is a hyperbolic equation.

- The **diffusion equation**,

$$\frac{\partial u}{\partial t} = \sigma \frac{\partial^2 u}{\partial x^2}$$

which is a parabolic equation.

Space-Time Grid:

	Space axis horizontal Time axis vertical						
$n = 5$	+	+	+	+	+	+	+
	I	I	I	I	I	I	I
$n = 4$	+	+	+	+	+	+	+
	I	I	I	I	I	I	I
$n = 3$	+	+	+	+	+	+	+
	I	I	I	I	I	I	I
$n = 2$	+	+	+	+	+	+	+
	I	I	I	I	I	I	I
$n = 1$	+	+	+	+	+	+	+
	I	I	I	I	I	I	I
$n = 0$	+	+	+	+	+	+	+
	j=-3	j=-2	j=-1	j=0	j=1	j=2	j=3

The Finite Difference Method

We take *discrete values* for x and t : $x_j = j\Delta x$ and $t_n = n\Delta t$.

The solution of the finite difference equation is also defined at the discrete points $(x_j, t_n) = (j\Delta x, n\Delta t)$:

$$U_j^n = U(j\Delta x, n\Delta t) = U(x_j, t_n).$$

That is, we use a small u to denote the solution of the PDE (*continuous*) and a capital U to denote the solution of the finite difference equation (FDE, a *discrete solution*).

Consider again the advection equation

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0,$$

Suppose we choose to approximate this **PDE** with the **FDE**

$$\frac{U_j^{n+1} - U_j^n}{\Delta t} + c \frac{U_j^n - U_{j-1}^n}{\Delta x} = 0$$

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To repeat:

$$\frac{U_j^{n+1} - U_j^n}{\Delta t} + c \frac{U_j^n - U_{j-1}^n}{\Delta x} = 0$$

This is called an **upstream scheme** (we are assuming $c > 0$).

Note that both differences are **non-centered** with respect to the point $(x_j, t_n) = (j\Delta x, n\Delta t)$.

We can ask two fundamental questions:

- [1] Is the FDE **consistent** with the PDE?
- [2] For a given time $t > 0$, will the solution of the FDE **converge** to that of the PDE as $\Delta x \rightarrow 0$ and $\Delta t \rightarrow 0$?

We will clarify these questions below.

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Warning: Sometimes superscript n denotes a *power*; sometimes it is just an index. Be careful!

Truncation Errors and Consistency

The FDE is **consistent** with the PDE if, in the limit $\Delta x \rightarrow 0$ and $\Delta t \rightarrow 0$, the FDE coincides with the PDE.

Obviously, this is a basic requirement that the FDE should fulfill if its solutions are going to be good approximations of the solutions of the PDE.

Definition: The difference between the PDE and the FDE is called the **discretization error** or **local truncation error**.

Consistency is rather simple to verify:

- Substitute u instead of U in the FDE.
- Evaluate all terms using a Taylor series expansion centered on the point (x_j, t_n) .
- Subtract the PDE from the FDE.

If the difference (local truncation error) goes to zero as $\Delta x \rightarrow 0, \Delta t \rightarrow 0$, then the FDE is consistent with the PDE.

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Again, the local truncation error is:

$$\tau = \left(\frac{u_{tt}}{2}\right) \Delta t - \left(\frac{cu_{xx}}{2}\right) \Delta x + \text{H.O.T.} = O(\Delta t) + O(\Delta x)$$

Clearly, as $\Delta t \rightarrow 0$ and $\Delta x \rightarrow 0$, we have $\tau \rightarrow 0$.

Therefore, the FDE is consistent.

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Note that both the time and the space truncation errors are of first order, because the finite differences are uncentered in both space and time.

Truncation errors for centered differences are second order.

Therefore, in general, centered differences are more accurate than uncentered differences.

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Truncation errors are a crucial factor in determining forecast accuracy in NWP.

Example

Let us verify the consistency of the upstream scheme for the advection equation.

First, consider the Taylor series expansion:

$$\left. \begin{aligned} u_j^{n+1} &= \left(u + u_t \Delta t + \frac{1}{2} u_{tt} \Delta t^2 + \dots \right)_j^n \\ u_{j-1}^n &= \left(u - u_x \Delta x + \frac{1}{2} u_{xx} \Delta x^2 - \dots \right)_j^n \end{aligned} \right\}$$

We substitute this in the FDE and obtain

$$\left(u_t + \frac{1}{2} u_{tt} \Delta t + \dots \right)_j^n + c \left(u_x - \frac{1}{2} u_{xx} \Delta x + \dots \right)_j^n \simeq 0$$

Subtracting the PDE gives the **local truncation error**:

$$\tau = \left(\frac{u_{tt}}{2}\right) \Delta t - \left(\frac{cu_{xx}}{2}\right) \Delta x + \text{H.O.T.} = O(\Delta t) + O(\Delta x)$$

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Convergence and Stability

The second question posed above was whether the solution of the FDE **converges** to the PDE solution.

That is, if we let $\Delta x \rightarrow 0$ and $\Delta t \rightarrow 0$, so that $j\Delta x \rightarrow x$ and $n\Delta t \rightarrow t$, does $U(j\Delta x, n\Delta t) \rightarrow u(x, t)$?

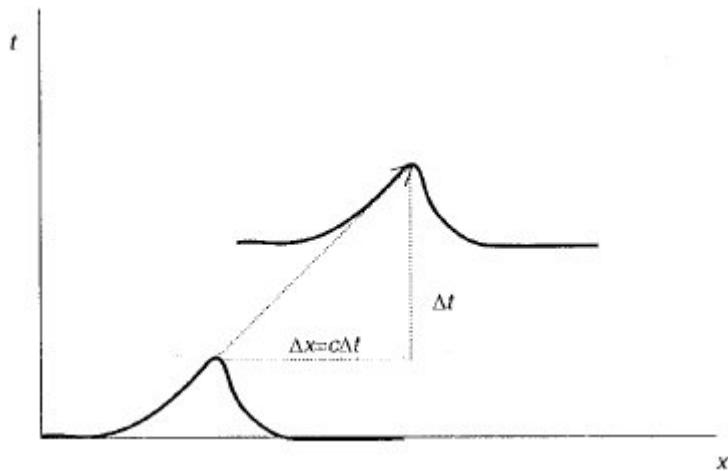
This is clearly important, and can be answered by considering another problem, that of **computational stability**.

Consider again the advection equation

$$\frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0,$$

which has the solution $u(x, t) = u(x - ct, 0)$.

The shape of the solution $u(x, 0)$ translates along the x -axis with velocity c (see Figure below).



Schematic of the solution of the advection equation (for $c > 0$).

The FDE for the upstream scheme can be written as

$$U_j^{n+1} = (1 - \mu)U_j^n + \mu U_{j-1}^n$$

where

$$\mu \equiv \frac{c\Delta t}{\Delta x}$$

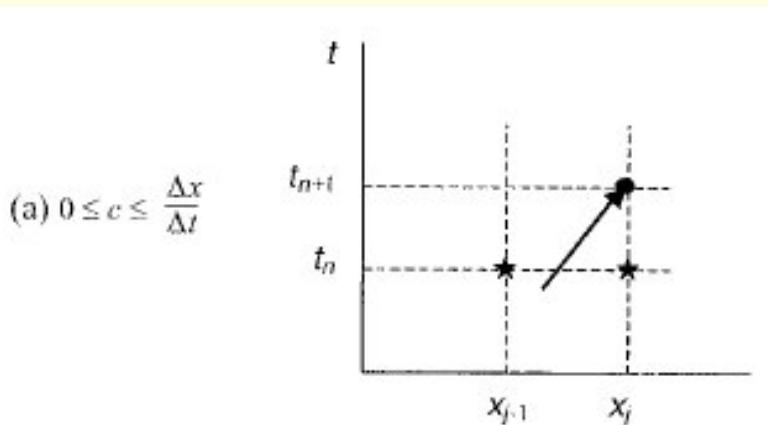
is the *Courant number* (or *Lewy Number*).

Let us suppose that $0 \leq \mu \leq 1$.

Then the FDE solution at the new time level U_j^{n+1} is **interpolated** between the values U_j^n and U_{j-1}^n .

In this case the advection scheme works the way it should, because the true solution lies in between those values.

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Schematic of the relationship between Δx , Δt and c leading to **interpolation** of the solution at time-level $n + 1$.

$$0 < \mu \equiv \frac{c\Delta t}{\Delta x} < 1$$

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However, suppose this condition is not satisfied, so that

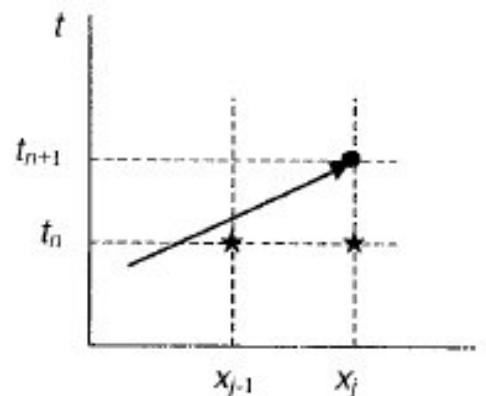
$$\mu = \frac{c\Delta t}{\Delta x} > 1 \quad \text{or else} \quad \mu = \frac{c\Delta t}{\Delta x} < 0.$$

Then the parcel arriving at point x_j at time t_{n+1} comes from somewhere outside the interval (x_{j-1}, x_j) at time t_n .

[Recall that $\partial u / \partial t + c \partial u / \partial x = 0$ is a linear approximation to $du / dt = 0$.]

Thus, the value of U_j^{n+1} is **extrapolated** from the values U_j^n and U_{j-1}^n .

$$(b) 0 \leq \frac{\Delta x}{\Delta t} \leq c$$



Schematic of the relationship between Δx , Δt and c leading to **extrapolation** of the solution at time-level $n + 1$.

$$\mu \equiv \frac{c\Delta t}{\Delta x} > 1$$

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The problem with extrapolation is that the maximum absolute value of the solution U_j^n increases with each time step.

Taking absolute values of the FDE we get

$$|U_j^{n+1}| \leq |U_j^n| |1 - \mu| + |U_{j-1}^n| |\mu|$$

Now defining $\Upsilon^n = \max_j |U_j^n|$, we have

$$\Upsilon^{n+1} \leq \{|1 - \mu| + |\mu|\} \Upsilon^n$$

and $\Upsilon^{n+1} \leq \Upsilon^n$ if and only if $0 \leq \mu \leq 1$.

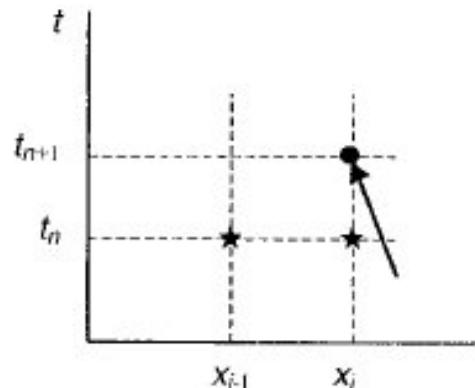
If the condition $0 \leq \mu \leq 1$ is not satisfied, then the solution is not bounded and it grows with n .

If we let $\Delta t \rightarrow 0$ and $\Delta x \rightarrow 0$ with $\mu = \text{const.}$, it only makes things worse, because then $n \rightarrow \infty$.

In practice, if the condition $0 \leq \mu \leq 1$ is not satisfied, the FDE **blows up** in a few time steps.

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$$(c) c \leq 0 \leq \frac{\Delta x}{\Delta t}$$



Schematic of the relationship between Δx , Δt and c leading to **extrapolation** of the solution at time-level $n + 1$.

$$\mu \equiv \frac{c\Delta t}{\Delta x} < 0$$

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Practical Exercise:

Use the simple model SLAM to explore the phenomenon of computational instability.

- Choose Δt small and carry out a complete integration.
- Increase Δt and watch the model solution “blowing up”.
- By numerical experiment, determine approximately the maximum value of Δt which yields stable integrations.
- Relate this maximum to the Courant Number.
What does it imply about the maximum phase-speed of the system?

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Computational Stability

We now define **computational stability**.

Definition: An FDE is **computationally stable** if the solution of the FDE at a fixed time $t = n\Delta t$ is bounded as $\Delta t \rightarrow 0$.

Note that with $n\Delta t$ fixed, $\Delta t \rightarrow 0$ implies $n \rightarrow \infty$.

We will derive a condition for stability which involves the **Courant Number**.

The condition on the Courant number is usually known as the **Courant-Friedrichs-Lowy criterion** or simply the **CFL condition**.

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Recall the story of Courant, Friedrichs and Lewy in Göttingen.

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Example: We use the **criterion of the maximum** method to study the stability condition of the diffusion equation

$$\frac{\partial u}{\partial t} = \sigma \frac{\partial^2 u}{\partial x^2}$$

A FDE approximation (FTCS scheme) is given by

$$\frac{U_j^{n+1} - U_j^n}{\Delta t} = \sigma \frac{U_{j+1}^n - 2U_j^n + U_{j-1}^n}{\Delta x^2}$$

(verification of **consistency** of this FDE is immediate).

Note: Since the differences are **centered in space** but **forward in time**, the truncation error is first order in time and second order in space

$$\tau = O(\Delta t) + O(\Delta x)^2.$$

We can write the FDE in the form

$$U_j^{n+1} = \mu U_{j+1}^n + (1 - 2\mu)U_j^n + \mu U_{j-1}^n$$

where $\mu = \sigma \Delta t / \Delta x^2$.

The Lax-Richtmyer Theorem

We can now state the fundamental Lax–Richtmyer theorem:

*Given a properly posed linear initial value problem, and a finite difference scheme that satisfies the consistency condition, then the **stability** of the FDE is the necessary and sufficient condition for **convergence**.*

$$\frac{\text{Stability}}{\text{For consistent systems}} \iff \text{Convergence}$$

This theorem allows us to establish **convergence** by examining the easier questions of **consistency** and **stability**.

We are interested in convergence not because we want to let $\Delta t, \Delta x \rightarrow 0$, but because we want to make sure that the errors $[u(j\Delta x, n\Delta t) - U_j^n]$ are acceptably small.

Definition: $[u(j\Delta x, n\Delta t) - U_j^n]$ is the **global truncation error**.

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Again,

$$U_j^{n+1} = \mu U_{j+1}^n + (1 - 2\mu)U_j^n + \mu U_{j-1}^n$$

If we take absolute values, and let $\Upsilon^n = \max_j |U_j^n|$, we get

$$\Upsilon^{n+1} \leq \{|\mu| + |1 - 2\mu| + |\mu|\} \Upsilon^n$$

Thus, we obtain the condition

$$0 \leq \mu \leq 1/2$$

to insure that the solution remains bounded as $n \rightarrow \infty$.

This is the necessary condition for stability of the FDE.

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Unfortunately, the **criterion of the maximum** can only be applied in very few cases.

In most FDEs some coefficients of the equations are negative, and the criterion cannot be applied.

We need a **more powerful method** of establishing stability.

The von Neumann Method

Another stability criterion that has much wider application is **the von Neumann stability criterion**.

We assume that we can expand the solution of the FDE in an appropriate set of eigenfunctions.

For simplicity we assume an expansion into **Fourier series**:

$$U(x, t) = \sum_k Z_k(t) e^{ikx}$$

The space variable x and the wavenumber k can be multi-dimensional, e.g., $\mathbf{x} = (x_1, x_2, x_3)$, $\mathbf{k} = (k_1, k_2, k_3)$ but, for simplicity, we will consider the scalar case.

We have $x_j = j\Delta x$ and $t_n = n\Delta t$.

We define the wavenumber for the Fourier series: $p = k\Delta x$.

Then the Fourier expansion is

$$U_j^n = \sum_p Z_p^n e^{ipj} \quad (\text{Note: } kx = kj\Delta x = pj)$$

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Aside: Spectral Radius

For the multi-dimensional case, the modulus of ρ is replaced by the **norm of the matrix G** and the stability condition becomes $\|G^n\| < M$ for all p , as $n \rightarrow \infty$.

The norm $\|G\|$ is a measure of the size of a matrix G .

Let $\sigma(G)$ be the **spectral radius** of G , i.e., $\sigma(G) = \max_i |\lambda_i|$, where λ_i are the eigenvalues of G ,

Then “it can be shown that”, for any norm,

$$[\sigma(G)]^n \leq \|G^n\| \leq \|G\|^n$$

The equal sign is valid if G is **normal**, i.e., if $G^* = G^*G$, where G^* is the transpose-conjugate of G , but in general the amplification matrices arising from FDEs are not normal.

End of digression

When we substitute this Fourier expansion into a linear FDE, we obtain a system of equations

$$Z_p^{n+1} = \rho_p Z_p^n$$

Here ρ_p is an **amplification factor** that, applied to the p -th Fourier component of the solution at time $n\Delta t$, advances it to the time $(n+1)\Delta t$; ρ_p depends on $p, \Delta t$ and Δx .

If we know the initial conditions

$$U_j^0 = \sum_p Z_p^0 e^{ipj}$$

then the solution of the FDE is (remember warning about superscripts)

$$Z_p^n = \rho_p^n Z_p^0$$

Therefore, stability is guaranteed if the factor ρ^n is bounded for all p when $\Delta t \rightarrow 0$ and $n \rightarrow \infty$.

So, we must have $|\rho_p|^n < M$ for all p as $n \rightarrow \infty$.

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We found that, for stability, we must have $|\rho|^n < M$ for all p as $n \rightarrow \infty$. Clearly, this requires

$$|\rho|^n \leq \exp \alpha \quad \text{for some constant } \alpha$$

Thus, a **necessary condition for stability**, and therefore a **necessary condition for convergence**, is that

$$\lim_{\Delta t \rightarrow 0, n\Delta t \rightarrow t} |\rho|^n = \text{finite} = e^\alpha$$

Then

$$|\rho| \leq [|\rho|^n]^{1/n} \leq e^{\alpha/n} = e^{\alpha\Delta t/t} \approx 1 + \frac{\alpha\Delta t}{t}$$

or simply

$$|\rho| \leq 1 + O(\Delta t)$$

This is the **von Neumann necessary condition** for computational stability.

Comment: The term $O(\Delta t)$ allows for bounded growth which may arise from a physical instability present in the PDE.

If the exact solution grows with time, then the FDE cannot both satisfy $|\rho| \leq 1$ and be consistent with the PDE.

Sufficient conditions are very complicated, and are known only for special cases.

In practice, we usually require $|\rho| \leq 1$ to guarantee computational stability.

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For more complicated equations, the von Neumann criterion involves a matrix G rather than the amplification factor ρ .

The stability criterion then involves the eigenvalues of the amplification matrix, and the von Neumann stability criterion is $\|G\| \leq 1 + O(\Delta t)$.

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We substitute $U_j^n = A\rho^n e^{ipj}$ in the equation and divide by U_j^n to obtain

$$\frac{\rho - 1}{\Delta t} + c \frac{(1 - e^{-ip})}{\Delta x} = 0 \quad \text{for all } p$$

The amplification factor ρ is the same as a 1×1 amplification matrix G , and the stability condition is $|\rho| \leq 1$ for all wavenumbers p .

We need to estimate the **maximum value of ρ** .

$$\rho = 1 - \mu(1 - e^{-ip}) = 1 - \mu(1 - \cos p + i \sin p)$$

Then the modulus squared is just

$$|\rho|^2 = [1 - \mu(1 - \cos p)]^2 + \mu^2 \sin^2 p$$

Application to Advection Equation

$$\text{PDE: } \frac{\partial u}{\partial t} + c \frac{\partial u}{\partial x} = 0$$

$$\text{FDE: } \frac{U_j^{n+1} - U_j^n}{\Delta t} + c \frac{U_j^n - U_{j-1}^n}{\Delta x} = 0 \quad (\text{upstream scheme})$$

We have already studied consistency, and used the criterion of the maximum to get a sufficient condition for stability.

Let us now apply the **von Neumann criterion**. Assume that

$$U_j^n = \sum_p Z_p^n e^{ipj} = \sum_p A_p \rho_p^n e^{ipj}$$

Since the equation is linear we can consider a single term

$$U_j^n = A_p \rho_p^n e^{ipj} = A_p \rho^n e^{ipj}$$

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To repeat,

$$|\rho|^2 = [1 - \mu(1 - \cos p)]^2 + \mu^2 \sin^2 p$$

We make use of the trigonometrical relationships

$$\cos p = \cos^2 \frac{p}{2} - \sin^2 \frac{p}{2} = c^2 - s^2 \quad \sin p = 2 \sin \frac{p}{2} \cos \frac{p}{2} = 2sc$$

Substituting these into $|\rho|^2$ we have

$$\begin{aligned} |\rho|^2 &= [1 - \mu(1 - c^2 + s^2)]^2 + 4\mu^2 s^2 c^2 \\ &= [1 - 2\mu s^2]^2 + 4\mu^2 s^2 (1 - s^2) \\ &= [1 - 4\mu s^2 + 4\mu^2 s^4] + 4\mu^2 s^2 - 4\mu^2 s^4 \\ &= 1 - 4\mu s^2 + 4\mu^2 s^2 \\ &= 1 - 4\mu(1 - \mu)s^2 \end{aligned}$$

Thus we obtain

$$|\rho|^2 = 1 - 4\mu(1 - \mu) \sin^2 \frac{p}{2}$$

To repeat,

$$|\rho|^2 = 1 - 4\mu(1 - \mu) \sin^2 \frac{p}{2}$$

First, consider the $\sin^2 p/2$ term: The shortest wave that can be present in the finite difference solution is $L = 2\Delta x$.

Therefore the maximum value that $p = k\Delta x = 2\pi\Delta x/L$ can take is $p = \pi$, and the maximum value of $\sin^2 p/2$ is 1.

Second, consider the factor $4\mu(1 - \mu)$. This is a parabola, whose maximum value is 1 when $\mu = 0.5$.

So the von Neumann condition is satisfied provided

$$0 \leq \mu \leq 1.$$

This coincides with the criterion of the maximum result.

It is also consistent with the idea that we should not extrapolate but always interpolate to get the new values.

Damping Effects of Scheme

The amplification factor ρ indicates how much the amplitude of each wavenumber will decrease or increase each time step.

The upstream scheme decreases the amplitude of all Fourier wave components of the solution, since $0 < \mu < 1 \Rightarrow |\rho| < 1$.

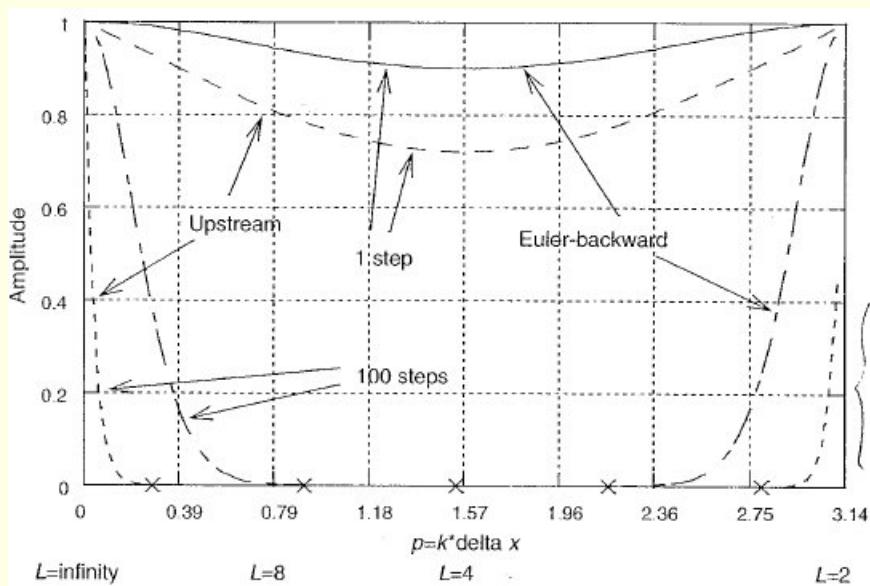
It is therefore a very dissipative FDE: it has **strong numerical diffusion**.

The figure below shows the decrease in amplitude when using the upstream scheme after one time step and after 100 time steps (the Courant number is $\mu = 0.5$).

Since the truncation error is large, the upstream scheme is in general not recommended except for special situations (e.g., for outflow boundary conditions).

An alternative, less damping scheme known as the **Matsuno** or **Euler-backward scheme** is also shown.

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