## Applied Analysis (ACM30020)

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Exercises #5

1. Given the following second order linear homogeneous ODE:

$$y'' + \omega^2 y = 0,$$

where  $\omega$  is a real number. The initial conditions are:  $y(0) = y_0$  and y'(0) = 0. Transform the problem into a Volterra integral equation. Solve the integral equation using an iterative scheme.

From class notes, the general solution of the IVP is:

$$y(x) = y_0 + p_2(x_0)y'_0P(x) + \int_{x_0}^x [r(t) - p_0(t)y(t)][P(x) - P(t)]dt,$$

where

$$P(x) = \int_{x_0}^x \frac{\mathrm{d}s}{p_2(s)}$$

Reading off, we have  $x_0 = 0$ , r(x) = 0,  $p_2(x) = 1$ , and  $p_0(x) = \omega^2$ . Also,  $y'_0 = 0$ . Hence, P(x) = x, and

$$y(x) = y_0 - \omega^2 \int_0^x y(t)(x-t) dt.$$
 (1)

This equation is in the form of a Volterra integral equation. We solve the equation iteratively using  $y^{(0)}(x) = y_0$ , and

$$y^{(1)}(x) = y_0 - \omega^2 \int_0^x y^{(0)}(t)(x-t) dt,$$
  
=  $y_0 - \omega^2 y_0 \int_0^x (x-t) dt,$   
=  $y_0 \left(1 - \frac{1}{2}\omega^2 x^2\right).$ 

Next, we have:

$$y^{(2)}(x) = y_0 - \omega^2 \int_0^x y^{(1)}(t)(x-t) dt,$$
  
=  $y_0 - \omega^2 y_0 \int_0^x \left(1 - \frac{1}{2}\omega^2 t^2\right)(x-t) dt$   
=  $y_0 \left(1 - \frac{1}{2}\omega^2 x^2 + \frac{1}{4!}\omega^4 x^4\right).$ 

Guessing the pattern, and letting  $n \to \infty$ , we obtain:

$$y(x) = y_0 \cos(\omega x).$$

We can double-check this by considering the integral

$$I = y_0 - \omega^2 y_0 \int_0^x \cos(\omega t) (x - t) \mathrm{d}t,$$

which by direct evaluation gives:

$$\frac{I}{y_0} = 1 - \omega^2 t \frac{\sin \omega t}{\omega} \Big|_0^x + \omega^2 \left[ \frac{1}{\omega^2} \cos(\omega t) \Big|_0^x + \frac{t \sin \omega t}{\omega} \Big|_0^x \right],$$
  
=  $\cos(\omega x).$ 

Hence,  $y(x) = y_0 \cos(\omega x)$  does indeed solve the integral equation (1).

2. Consider the inhomogeneous problem

$$y'' + \omega^2 y = f(x), \qquad x \in [0, \pi].$$

Here, the problem is a Boundary Value Problem (BVP), with boundary conditions  $y(0) = y(\pi) = 0$ ,.

Use the Green's Function to solve the BVP in the case when:

Linearly independent solutions of the homogeneous problem are:

$$u(x) = \sin(\omega x),$$

and

$$v(x) = \sin(\omega x) - \tan(\omega \pi) \cos(\omega x).$$

Hence, u(x) satisfies the left-hand Boundary Condition (LHBC) at x = 0 and v(x) satisfies the right-hand Boundary Condition (RHBC) at  $x = \pi$ . The Wronskian is:

$$W = \begin{vmatrix} \sin \omega x & \sin \omega x - \tan \omega \pi \cos \omega x \\ \omega \cos \omega x & \omega \cos x + \omega \tan \omega \pi \sin \omega x \end{vmatrix}.$$

By linear dependence, this simplifies:

$$W = \begin{vmatrix} \sin \omega x & -\tan \omega \pi \cos \omega x \\ \omega \cos \omega x & \omega \tan \omega \pi \sin \omega x \end{vmatrix} = \omega \tan(\omega \pi).$$

Thus, W = Const., consistent with Abel's Theorem with  $p_1/p_2 = 0/1$ .

The Green's function for the BVP is now:

$$G(x,s) = \frac{[\sin(\omega x) - \tan(\omega \pi) \cos(\omega x)] \sin(\omega s)}{\omega \tan \omega \pi}, \qquad 0 \le s \le x \le \pi,$$

 $\quad \text{and} \quad$ 

$$G(x,s) = \frac{\sin(\omega x) \left[\sin(\omega s) - \tan(\omega \pi) \cos(\omega s)\right]}{\omega \tan \omega \pi}, \qquad 0 \le x < s \le \pi.$$

By convolution, the solution to the BVP is:

$$y(x) = \int_0^{\pi} G(x,s)f(s)\mathrm{d}s.$$

Hence,

$$y(x) = \int_0^x \frac{[\sin(\omega x) - \tan(\omega \pi) \cos(\omega x)] \sin(\omega s)}{\omega \tan \omega \pi} f(s) ds + \int_x^\pi \frac{\sin(\omega x) [\sin(\omega s) - \tan(\omega \pi) \cos(\omega s)]}{\omega \tan \omega \pi} f(s) ds,$$

or

$$y(x) = \int_0^x \frac{[\sin(\omega x)\cos(\omega \pi) - \sin(\omega \pi)\cos(\omega x)]\sin(\omega s)}{\omega \sin \omega \pi} f(s) ds + \int_x^\pi \frac{\sin(\omega x)[\sin(\omega s)\cos(\omega \pi) - \sin(\omega \pi)\cos(\omega s)]}{\omega \sin \omega \pi} f(s) ds.$$

Or again,

$$y(x) = \int_0^x \frac{\sin(\omega(x-\pi))\sin(\omega s)}{\omega\sin\omega\pi} f(s) ds + \int_x^\pi \frac{\sin(\omega x)\sin(\omega(s-\pi))}{\omega\sin\omega\pi} f(s) ds.$$

Finally,

$$y(x) = \frac{\sin(\omega(x-\pi))}{\omega\sin\omega\pi} \int_0^x \sin(\omega s) f(s) ds + \frac{\sin(\omega x)}{\omega\sin\omega\pi} \int_x^\pi \sin(\omega(s-\pi)) f(s) ds.$$

We fill in for f(x) = 1:

$$y(x) = \frac{\sin(\omega(x-\pi))}{\omega \sin \omega \pi} \int_0^x \sin(\omega s) ds + \frac{\sin(\omega x)}{\omega \sin \omega \pi} \int_x^\pi \sin(\omega(s-\pi)) ds,$$
  

$$= \frac{\sin(\omega(x-\pi))}{\omega^2 \sin \omega \pi} [1 - \cos(\omega x)] + \frac{\sin(\omega x)}{\omega^2 \sin \omega \pi} [\cos(\omega(x-\pi)) - 1]$$
  

$$= \frac{1}{\omega^2 \sin(\omega \pi)} [\sin(\omega(x-\pi) - \sin(\omega x))]$$
  

$$+ \frac{1}{\omega^2 \sin \omega \pi} [\sin(\omega x) \cos(\omega(x-\pi)) - \cos(\omega x) \sin(\omega(x-\pi))].$$

Hence,

$$y(x) = \frac{1}{\omega^2 \sin(\omega \pi)} \left[ \sin(\omega(x - \pi) - \sin(\omega x)) \right] + \frac{1}{\omega^2}.$$
 (2)

Clearly,  $y''(x) + \omega^2 y = 1$ , and y(0) = 0 and  $y(\pi) = 0$ , so our calculations are correct and Equation (2) does indeed solve the BVP.

Next, we fill in for  $f(x) = \sin(\omega x)$ :

$$y(x) = \frac{\sin(\omega(x-\pi))}{\omega\sin\omega\pi} \int_0^x \sin(\omega s) \sin(\omega s) ds + \frac{\omega\sin(\omega x)}{\sin\omega\pi} \int_x^\pi \sin(\omega(s-\pi)) \sin(\omega s) ds$$

This 'simplifies' to:

$$y(x) = \frac{\sin(\omega(x-\pi))}{\omega\sin\omega\pi} \left[ \frac{1}{2}x - \frac{\sin(2\omega x)}{4\omega} \right] + \frac{\sin(\omega x)}{\omega\sin\omega\pi} \left[ \frac{1}{2}(\pi-x)\cos(\omega\pi) - \frac{\sin(\omega(\pi-2x))}{4\omega} - \frac{\sin(\omega\pi)}{4\omega} \right]$$
(3)

We have checked this result using a numerical 'shooting' method (Figure 1). Matlab listings are provided in the Appendix.



Figure 1: Comparison between numerical shooting method and Equation (3), for  $\omega = 2.5$ .

3. The stationary temperature distribution in a rod of unit length that has both ends kept at a constant zero temperature, with heat loss through its surface proportional to u, and that is subject to a given non-uniform heat source per unit length f(x), is the solution of

$$-u'' + u = f, \qquad u(0) = u(1) = 0.$$

Show that the Green's function of this Boundary Value Problem is given by:

$$G(x,\xi) = \begin{cases} \frac{\sinh(x)\sinh(1-\xi)}{\sinh(1)}, & 0 \le x \le \xi \le 1, \\ \frac{\sinh(\xi)\sinh(1-x)}{\sinh(1)}, & 0 \le \xi \le x \le 1. \end{cases}$$

Linearly independent solutions of the homogeneous problem are:

$$u(x) = \sinh(x)$$

and

$$v(x) = \sinh(1)\cosh(x) - \cosh(1)\sinh(x)$$

Hence, u(x) satisfies the left-hand Boundary Condition (LHBC) at x = 0 and v(x) satisfies the right-hand Boundary Condition (RHBC) at x = 1. The Wronskian is:

$$W = \begin{vmatrix} u & v \\ u' & v' \end{vmatrix},$$
  
= 
$$\begin{vmatrix} \sinh(x) & \sinh(1)\cosh(x) \\ \cosh(x) & \sinh(1)\sinh(x) \end{vmatrix},$$
  
= 
$$-\sinh(1).$$

Also,  $p_2 = -1$ , hence  $p_2W = \sinh(1)$ .

Reading off from class notes, the Green's function is now:

.

$$G(x,\xi) = \frac{[\sinh(1)\cosh(x) - \cosh(1)\sinh(x)]\sinh(\xi)}{\sinh(1)}, \qquad 0 \le \xi \le x \le 1.$$

The term in the square bracket can be written as  $\sinh(1-x)$ , using a trig identity. Also,

$$G(x,\xi) = \frac{\sinh(x)\left[\sinh(1)\cosh(\xi) - \cosh(1)\sinh(\xi)\right]}{\sinh(1)}, \qquad 0 \le x \le \xi \le 1.$$

Again, a trig identity can be used on the term in the square bracket, this is  $\sinh(1-\xi)$ . Hence,

$$G(x,\xi) = \begin{cases} \frac{\sinh(\xi)\sinh(1-x)}{\sinh(1)}, & 0 \le \xi \le x \le 1, \\ \frac{\sinh(x)\sinh(1-\xi)}{\sinh(1)}, & 0 \le x \le \xi \le 1, \end{cases}$$

as required.

4. Consider the Fredholm integral equation

$$y(x) = f(x) + \lambda \int_{-1}^{1} (x+s)y(s) \mathrm{d}s.$$

- (a) For which values of  $\lambda$  does the equation have a unique solution? Find the solution in this case.
- (b) For each of those values of  $\lambda$  for which the equation does not have a unique solution, state a condition which f(x) must satisfy in order for a solution to exist, and find the general solution when this is satisfied.

We have:

$$K(x,s) = u_1(x)v_1(s) + u_2(x)v_2(s),$$

where  $u_1(x) = x$ ,  $v_1(s) = 1$ ,  $u_2(x) = 1$ , and  $v_2(s) = 1$ . Thus, the FIE can be written as:

$$y(x) = f(x) + \lambda x c_1 + \lambda c_2 \tag{4}$$

where

$$c_1 = \int_{-1}^{1} y(s) ds,$$
  $c_2 = \int_{-1}^{1} sy(s) ds.$ 

Let

$$f_1 = \int_{-1}^{1} f(s) ds, \qquad f_2 = \int_{-1}^{1} sf(s) ds$$

We multiply the FIE (4) by  $v_1$  and integrate to obtain:

$$c_1 = f_1 + 2\lambda c_2. \tag{5a}$$

Next, we multiply the FIE (4) by  $v_2$  and integrate to obtain:

$$c_2 = f_2 + 2\lambda c_1/3. \tag{5b}$$

In doing these integrals, it is helpful to remember the properties of integrals of odd/even functions over symmetric intervals, e.g.

$$\int_{-1}^{1} x \mathrm{d}x = 0,$$

etc. Equations (5) can be simplified and written in matrix form:

$$\underbrace{\begin{pmatrix} 1 & -2\lambda \\ -2\lambda/3 & 1 \end{pmatrix}}_{=M(\lambda)} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$$

The characteristic polynomial is

$$\det\left[M(\lambda)\right] = 0,$$

hence

$$1 - \frac{4}{3}\lambda^2 = 0$$

Thus, the roots of the characteristic polynomial are given by:

$$\lambda \in \mathcal{S} = \{\lambda_{-}, \lambda_{+}\} = \{-\frac{\sqrt{3}}{2}, \frac{\sqrt{3}}{2}\}.$$

Thus, the answer to **part (a)** is that there is a unique solution if  $\lambda \neq S$ . For **part (b)**, we look at the two separate cases.

• Case 1,  $\lambda = -\sqrt{3}/2$ . Here, we require  $f_1 = \sqrt{3}f_2$ , that is

$$\int_{-1}^{1} f(s)(1 - \sqrt{3}x) \mathrm{d}x = 0.$$

in which case the general solution is given for arbitrary  $f_2 \ {\rm by}$ 

$$y(x) = f(x) + \lambda_{-} \left[ \left( -\sqrt{3}c_{2} + b_{1} \right) x + c_{2} \right],$$
  
$$= f(x) + \lambda_{-}b_{1}x + \lambda_{-}c_{2} \underbrace{ \left( 1 - \sqrt{3}x \right)}_{=y_{-}}.$$

Thus, the general solution is:

$$y(x) = f(x) + \lambda_{-}b_{1}x + \alpha y_{-}(x),$$

where  $y_{-}(x)$  satisfies the homogeneous FIE:

$$y_{-}(x) = \lambda_{-} \int_{-1}^{1} K(x,s) y_{-}(s) \mathrm{d}s.$$

The compatibility condition is:

$$\int_{-1}^{1} y_{-}(x) f(x) \mathrm{d}x = 0.$$

• Case 2,  $\lambda=\sqrt{3}/2$ . Here, we require  $f_1=-\sqrt{3}f_2$ , that is

$$\int_{-1}^{1} f(s)(1+\sqrt{3}x) \mathrm{d}x = 0,$$

in which case the general solution is given for arbitrary  $f_2$  by

$$y(x) = f(x) + \lambda_{+} \left[ \left( \sqrt{3}c_{2} + b_{1} \right) x + c_{2} \right],$$
  
=  $f(x) + \lambda_{+}b_{1}x + \lambda_{+}c_{2} \underbrace{ \left( 1 + \sqrt{3}x \right)}_{=y_{-}}.$ 

Thus, the general solution is:

$$y(x) = f(x) + \lambda_+ b_1 x + \alpha y_+(x),$$

where  $y_+(x)$  satisfies the homogeneous FIE:

$$y_{+}(x) = \lambda_{-} \int_{-1}^{1} K(x,s) y_{+}(s) \mathrm{d}s.$$

The compatibility condition is:

$$\int_{-1}^{1} y_{+}(x) f(x) \mathrm{d}x = 0.$$

**Remark:** When we look at Hilbert–Schmidt theory, we will find, for  $\lambda \notin S$ , that the solution of the FIE can be written as:

$$y(x) = f(x) + \frac{\lambda}{\lambda_+ - \lambda} f_+ y_+(x) + \frac{\lambda}{\lambda_- - \lambda} f_- y_+(x),$$

where

$$\lambda_{\pm} = \pm \sqrt{3}2.$$

Also,

$$y_{\pm}(x) = \frac{1 \pm \sqrt{3}x}{2}$$

are the normalized eigenfunctions of the FIE, with

$$\int_{-1}^{1} y_i(x) y_j(x) dx = \delta_{ij}, \qquad i, j \in \{-, +\},$$

Furthermore,

$$f_{\pm} = \int_{-1}^{1} f(x) y_{\pm}(x) \mathrm{d}x.$$

Thus, we see transparently that in for part (b), the condition which f must satisfy in order for a solution to exist when  $\lambda = \lambda_+$  (respectively  $\lambda = \lambda_-$ ) is  $f_+ = 0$ (respectively  $f_- = 0$ ). 5. Solve for  $\phi(x)$  in the integral equation

$$\phi(x) = f(x) + \lambda \int_0^1 \left[ \left(\frac{x}{y}\right)^n + \left(\frac{y}{x}\right)^n \right] \phi(y) \mathrm{d}y,$$

where f(x) is bounded for 0 < x < 1 and -1/2 < n < 1/2, expressing your answer in terms of the quantities  $F_m = \int_0^1 f(y) y^m dy$ .

- (a) Give the explicit solution when  $\lambda = 1$ .
- (b) For what values of  $\lambda$  are there no solutions unless  $F_{\pm n}$  are in a particular ratio? What is this ratio?

Solution: We are looking at:

$$\phi(x) = f(x) + \lambda \int_0^1 K(x, y)\phi(y) \,\mathrm{d}y,$$

hence a Fredholm integral equation with kernel  $K(x, y) = (x/y)^n + (y/x)^n$ . This is a separable kernel,

$$K(x,s) = u_1(x)v_1(s) + u_2(x)v_2(s),$$

with  $u_1(x) = x^n$ ,  $v_1(s) = s^{-n}$ ,  $u_2(x) = x^{-n}$ , and  $v_2(s) = s^n$ . We have:

$$c_{1} = \int_{0}^{1} v_{1}(s)\phi(s)ds = \int_{0}^{1} s^{-n}\phi(s)ds$$
  
$$c_{2} = \int_{0}^{1} v_{2}(s)\phi(s)ds = \int_{0}^{1} s^{n}\phi(s)ds.$$

Also (using notation from class notes):

$$b_{1} = \int_{0}^{1} v_{1}(s)f(s)ds = \int_{0}^{1} s^{-n}f(s)ds = F_{-n},$$
  

$$b_{2} = \int_{0}^{1} v_{2}(s)f(s)ds = \int_{0}^{1} s^{n}f(s)ds = F_{n}.$$

We project the integral equation on to  $v_1$  and  $v_2$  to get:

$$c_i = b_i + \lambda \sum_{j=1}^2 A_{ij} c_j,$$

where  $A_{ij} = \int_0^1 v_i(s) u_j(s) ds$ . Hence, it remains to compute:

$$A_{11} = \int_{0}^{1} 1 \, ds = 1,$$
  

$$A_{12} = \int_{0}^{1} s^{-2n} \, ds = \frac{1}{1 - 2n},$$
  

$$A_{21} = \int_{0}^{1} s^{2n} \, ds = \frac{1}{2n + 1},$$
  

$$A_{22} = \int_{0}^{1} 1 \, ds = 1.$$

We don't encounter any divisions by zero here because |n| < 1/2. Thus, the linear problem to solve is:

$$\underbrace{\begin{pmatrix} 1-\lambda & -\frac{\lambda}{1-2n} \\ -\frac{\lambda}{1+2n} & 1-\lambda \end{pmatrix}}_{=M} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}.$$
 (6)

We focus on the characteristic equation:

$$\left|\begin{array}{cc} 1-\lambda & -\frac{\lambda}{1-2n} \\ -\frac{\lambda}{1+2n} & 1-\lambda \end{array}\right| = 0.$$

The solutions are:

$$\lambda_{+} = \frac{\sqrt{1 - 4n^2}}{1 + \sqrt{1 - 4n^2}}, \qquad \lambda_{-} = \frac{\sqrt{1 - 4n^2}}{-1 + \sqrt{1 - 4n^2}}.$$

Thus, when  $\lambda = 1$ , it is the case that  $\lambda \neq \lambda_{\pm}$ , and hence, Equation (6) has a unique solution, characterized by:

$$M^{-1} = \left(\begin{array}{cc} 0 & -(2n+1) \\ 2n-1 & 0 \end{array}\right).$$

Hence,

$$(c_1, c_2)^T = \begin{pmatrix} 0 & -(2n+1) \\ 2n-1 & 0 \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}$$
$$= \begin{pmatrix} 0 & -(2n+1) \\ 2n-1 & 0 \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}.$$

Hence,

$$c_1 = -(1+2n)b_2, \qquad c_2 = -(1-2n)b_1.$$

The solution in the case  $\lambda = 1$  is therefore:

$$\begin{aligned}
\phi_{\lambda=1}(x) &= f(x) + c_1 u_1(x) + c_2 u_2(x), \\
&= f(x) - (1+2n)b_2 x^n - (1-2n)b_1 x^{-n}.
\end{aligned}$$

Or, in terms of the notation in the question,

$$\phi_{\lambda=1}(x) = f(x) - (1+2n)F_n x^n - (1-2n)F_{-n} x^{-n}.$$

For part (b), we look again at Equation (6) with  $\lambda = \lambda_{\pm}$ . In an obvious notation, we have:

$$M(\lambda_{\pm})\boldsymbol{c} = \boldsymbol{b},\tag{7}$$

and  $\det M(\lambda_{\pm}) = 0$ . Hence, for almost all values of n, Equation (6) with  $\lambda = \lambda_{\pm}$  will have no solution.

However, there is one special case to look at. Hence, we re-write Equation (7) as:

$$M(\lambda_{\pm})\boldsymbol{c} = \boldsymbol{b},$$
  
$$\lambda_{\pm} \begin{pmatrix} \frac{1}{\lambda_{\pm}} - 1 & -\frac{1}{1-2n} \\ -\frac{1}{1+2n} & \frac{1}{\lambda_{\pm}} - 1 \end{pmatrix} \boldsymbol{c} = \boldsymbol{b}.$$

Hence:

$$\lambda_{\pm} \left( \begin{array}{cc} \pm \frac{1}{\sqrt{1-4n^2}} & -\frac{1}{1-2n} \\ -\frac{1}{1+2n} & \pm \frac{1}{\sqrt{1-4n^2}} \end{array} \right) \boldsymbol{c} = \boldsymbol{b}.$$

In other words,

$$\pm \frac{1}{\sqrt{1-4n^2}}c_1 - \frac{1}{1-2n}c_2 = b_1/\lambda_{\pm},$$
  
$$(\mp)\frac{1+2n}{\sqrt{1-4n^2}} \times \left[-\frac{1}{1+2n}c_1 \pm \frac{1}{\sqrt{1-4n^2}}c_2\right] = (\mp)\frac{1+2n}{\sqrt{1-4n^2}} \times b_2/\lambda_{\pm}.$$

The second row can now be re-written as:

$$\pm \frac{1}{\sqrt{1-4n^2}}c_1 - \frac{1}{1-2n}c_2 = \mp \sqrt{\frac{1+2n}{1-2n}}(b_2/\lambda_{\pm}).$$

Compare the two rows again:

$$\pm \frac{1}{\sqrt{1-4n^2}}c_1 - \frac{1}{1-2n}c_2 = b_1/\lambda_{\pm},$$
  
$$\pm \frac{1}{\sqrt{1-4n^2}}c_1 - \frac{1}{1-2n}c_2 = \mp \sqrt{\frac{1+2n}{1-2n}}b_2/\lambda_{\pm}.$$

If the two rows are the same, there are infinitely many solutions. Otherwise, there are no solutions. Hence, to find the case with solutions, we require:

$$\frac{b_1}{b_2} = \mp \sqrt{\frac{1+2n}{1-2n}}.$$

Lastly, go back to the notation in the question  $(b_1 \rightarrow F_{-n}, b_2 \rightarrow F_n)$ :

$$\frac{F_n}{F_{-n}} = \mp \sqrt{\frac{1-2n}{1+2n}}.$$

6. Consider the FIE

$$y(x) = f(x) + \lambda \int_0^1 \cosh(x - s) y(s) \mathrm{d}s.$$
 (8)

- (a) Show that the eigenvalues of (8) are given by  $2/(1 \pm \sinh 1)$ .
- (b) Using the Hilbert–Schmidt eigenfunction expansion, or otherwise, find the solution for  $\lambda \notin \{2/(1 \pm \sinh 1)\}$ .
- (c) Find a necessary and sufficient condition on f for the equation

$$y(x) = f(x) + \frac{2}{1 + \sinh 1} \int_0^1 \cosh(x - s) y(s) ds$$

to have a solution and find all solutions when this condition is satisfied.

For part (a), note that the kernel is separable, since

$$K(x,s) = \cosh(x-s) = \frac{1}{2}e^{x}e^{-s} + \frac{1}{2}e^{-x}e^{s}.$$

To find the eigenvalues, let us write the homogeneous equation as

$$y(x) = \frac{1}{2} \mathrm{e}^x \lambda \, c_1 + \frac{1}{2} \mathrm{e}^{-x} \lambda \, c_2,$$

where

$$c_1 = \int_0^1 e^{-s} y(s) ds,$$
  $c_2 = \int_0^1 e^{s} y(s) ds.$ 

Then

$$c_{1} = \int_{0}^{1} e^{-x} y(x) \, \mathrm{d}x = \frac{1}{2} \lambda c_{1} \int_{0}^{1} e^{-x} \mathrm{e}^{x} \, \mathrm{d} + \frac{1}{2} \lambda c_{2} \int_{0}^{1} \mathrm{e}^{-x} \mathrm{e}^{-x} \, \mathrm{d}x$$
$$= \frac{1}{2} \lambda c_{1} + \frac{1}{2} \lambda c_{2} \frac{1}{2} (1 - \mathrm{e}^{-2}),$$

and

$$c_{2} = \int_{0}^{1} e^{x} y(x) dx = \frac{1}{2} \lambda c_{1} \int_{0}^{1} e^{x} e^{x} dx + \frac{1}{2} \lambda c_{2} \int_{0}^{1} e^{x} e^{-x} dx$$
$$= \frac{1}{2} \lambda c_{1} \frac{1}{2} (e^{2} - 1) + \frac{1}{2} \lambda c_{2}.$$

These equations may be written in matrix form as

$$\underbrace{\begin{pmatrix} \lambda - 2 & \lambda \frac{1}{2}(1 - e^{-2}) \\ \lambda \frac{1}{2}(e^{2} - 1) & \lambda - 2 \end{pmatrix}}_{=M(\lambda)} \begin{pmatrix} c_{1} \\ c_{2} \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$
(9)

The charactrestic polynomial is:

$$\det\left[M(\lambda)\right] = 0.$$

This evaluates to:

$$(\lambda - 2)^2 = \frac{1}{4}\lambda^2(e^2 - 1)(1 - e^{-2}) = \frac{1}{4}\lambda^2(e - e^{-1})^2 = \lambda^2\sinh^2(1).$$

Hence, the eigenvalues of the FIE are:

$$\lambda \in \mathcal{S} = \{\lambda_1, \lambda_2\} = \left\{\frac{2}{1+\sinh 1}, \frac{2}{1-\sinh 1}\right\}$$

For **part (b)**, the pertinent eigenfunctions are obtained by solving for  $(c_1, c_2)^T$  in Equation (9). We look at the two cases:

• Case 1. We look at  $\lambda = \lambda_1$ . Notice that  $\lambda_1 \sinh 1 = -(\lambda_1 - 2)$  so, for  $\lambda = \lambda_1$ , Equation (9) becomes:

$$\begin{pmatrix} \lambda_1 - 2 & -(\lambda_1 - 2)e^{-1} \\ -(\lambda_1 - 2)e & \lambda_1 - 2 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

This imples that  $c_2 = ec_1$ , so the corresponding eigenfunction is a multiple of

$$e^{x} + ee^{-x} = 2e^{1/2}\frac{1}{2}(e^{x-1/2} + e^{-(x-1/2)}),$$

or equivalently:

$$y_1(x) \propto \cosh(x - \frac{1}{2}).$$

• For  $\lambda = \lambda_2$ , we have  $c_2 = -ec_1$ , and the corresponding eigenfunction is

$$y_2(x) \propto \sinh(x - \frac{1}{2}).$$

We look at normalized eigenfunctions:

$$y_1(x) = \sqrt{\lambda_1} \cosh\left(x - \frac{1}{2}\right), \qquad y_2(x) = \sqrt{-\lambda_2} \sinh\left(x - \frac{1}{2}\right).$$

Notice that:

$$\int_0^1 y_i(x) y_j(x) \mathrm{d}x = \delta_{ij}.$$

In this case, Hilbert–Schmidt theory tells us that the solution fo the FIE can be expanded in terms of eigenfunctions:

$$y(x) = f(x) + \frac{\lambda}{\lambda_1 - \lambda} f_1 y_1(x) + \frac{\lambda}{\lambda_2 - \lambda} f_2 y_2(x), \qquad \lambda \notin \mathcal{S}$$

Here,

$$f_i = \int_0^1 f(x) y_i(x) \mathrm{d}x.$$

For part (c), Hilbert–Schmidt theory tells us that if  $\lambda \in S$  (in this case,  $\lambda = \lambda_1$ ), then a necessary necessary and sufficient condition on f for the FIE to have a solution is:

$$\int_0^1 f(x)y_1(x)\mathrm{d}x = 0 \implies f_1 = 0.$$

In this case, the general solution is

$$y(x) = f(x) + Ay_1(x) + \frac{\lambda_1}{\lambda_2 - \lambda_1} f_2 y_2(x).$$

## A Code Listings

In this code, I solve the ODE  $y''(x) = -\omega^2 y(x) + \sin(x)$  for the initial conditions y(0) = 0 and y'(0) = p. Here, p is a parameter that can be varied.

In this next code, I look at the cost function

$$J(p) = [y(x = \pi; p)]^2$$
,

where  $y(x = \pi, p)$  is output from odesolve.m, at paramter value p. I solve an optimization problem where I make J(p) as small as possible. In this way, I solve the BVP  $y''(x) = -\omega^2 y(x) + \sin(x)$  subject to y(0) = 0 and  $y(\pi) = 0$ . This is called the **shooting method**.

```
p_star=fminbnd(@mycost,-10 ,10000);
display(p_star);
J_min=mycost(p_star);
display(J_min);
% Obtain solution of BVP:
[x,y]=odesolve(p_star);
% Interpolate solution on to regular grid:
xi=0:0.1:pi;
yi=interp1(x,y,xi,'spline');
x=xi;
y=yi;
```

```
% Plot solution:
plot(x,y,'o','color','blue')
% Calculate analytical solution for mega = 2.5:
omega = 2.5;
xa = 0:0.01:pi;
y_a = sin(omega * (x_a - pi)) . * (0.5 * x_a - (sin(2 * omega * x_a) / (4 * omega))) ...
      +\sin(\operatorname{omega} \times xa) \times (0.5 \times (\operatorname{pi} - xa) \times \cos(\operatorname{omega} \times \operatorname{pi}) \dots
      -(\sin(\operatorname{omega}*(\operatorname{pi}-2*\operatorname{xa}))/(4*\operatorname{omega}))
-(\sin(\operatorname{omega}*\operatorname{pi})/(4*\operatorname{omega})));
ya=ya / (omega * sin (omega * pi));
% Plot analyical and numerical solution on same axes
% and compare:
hold on
plot(xa,ya,'color','red','linewidth',2)
grid on
xlabel('x')
ylabel('y')
set(gca, 'fontsize',14, 'fontname', 'times new roman')
hold off
function J=mycost(p)
      [^{\circ}, y] = odesolve(p);
      J=y(end)*y(end);
end
```