

Applied Analysis (ACM30020)

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Exercises #5

- Given the following second order linear homogeneous ODE:

$$y'' + \omega^2 y = 0,$$

where ω is a real number. The initial conditions are: $y(0) = y_0$ and $y'(0) = 0$.

Transform the problem into a Volterra integral equation. Solve the integral equation using an iterative scheme.

From class notes, the general solution of the IVP is:

$$y(x) = y_0 + p_2(x_0)y'_0 P(x) + \int_{x_0}^x [r(t) - p_0(t)y(t)][P(x) - P(t)]dt,$$

where

$$P(x) = \int_{x_0}^x \frac{ds}{p_2(s)}.$$

Reading off, we have $x_0 = 0$, $r(x) = 0$, $p_2(x) = 1$, and $p_0(x) = \omega^2$. Also, $y'_0 = 0$. Hence, $P(x) = x$, and

$$y(x) = y_0 - \omega^2 \int_0^x y(t)(x-t)dt. \quad (1)$$

This equation is in the form of a Volterra integral equation. We solve the equation iteratively using $y^{(0)}(x) = y_0$, and

$$\begin{aligned} y^{(1)}(x) &= y_0 - \omega^2 \int_0^x y^{(0)}(t)(x-t)dt, \\ &= y_0 - \omega^2 y_0 \int_0^x (x-t)dt, \\ &= y_0 \left(1 - \frac{1}{2}\omega^2 x^2\right). \end{aligned}$$

Next, we have:

$$\begin{aligned} y^{(2)}(x) &= y_0 - \omega^2 \int_0^x y^{(1)}(t)(x-t)dt, \\ &= y_0 - \omega^2 y_0 \int_0^x \left(1 - \frac{1}{2}\omega^2 t^2\right) (x-t)dt, \\ &= y_0 \left(1 - \frac{1}{2}\omega^2 x^2 + \frac{1}{4!}\omega^4 x^4\right). \end{aligned}$$

Guessing the pattern, and letting $n \rightarrow \infty$, we obtain:

$$y(x) = y_0 \cos(\omega x).$$

We can double-check this by considering the integral

$$I = y_0 - \omega^2 y_0 \int_0^x \cos(\omega t)(x-t)dt,$$

which by direct evaluation gives:

$$\begin{aligned} \frac{I}{y_0} &= 1 - \omega^2 t \frac{\sin \omega t}{\omega} \Big|_0^x + \omega^2 \left[\frac{1}{\omega^2} \cos(\omega t) \right]_0^x + \frac{t \sin \omega t}{\omega} \Big|_0^x, \\ &= \cos(\omega x). \end{aligned}$$

Hence, $y(x) = y_0 \cos(\omega x)$ does indeed solve the integral equation (1).

2. Consider the inhomogeneous problem

$$y'' + \omega^2 y = f(x), \quad x \in [0, \pi].$$

Here, the problem is a Boundary Value Problem (BVP), with boundary conditions $y(0) = y(\pi) = 0$.

Use the Green's Function to solve the BVP in the case when:

- (a) $f(x) = 1$;
- (b) $f(x) = \sin(\omega x)$.

Linearly independent solutions of the homogeneous problem are:

$$u(x) = \sin(\omega x),$$

and

$$v(x) = \sin(\omega x) - \tan(\omega \pi) \cos(\omega x).$$

Hence, $u(x)$ satisfies the left-hand Boundary Condition (LHBC) at $x = 0$ and $v(x)$ satisfies the right-hand Boundary Condition (RHBC) at $x = \pi$. The Wronskian is:

$$W = \begin{vmatrix} \sin \omega x & \sin \omega x - \tan \omega \pi \cos \omega x \\ \omega \cos \omega x & \omega \cos x + \omega \tan \omega \pi \sin \omega x \end{vmatrix}.$$

By linear dependence, this simplifies:

$$W = \begin{vmatrix} \sin \omega x & -\tan \omega \pi \cos \omega x \\ \omega \cos \omega x & \omega \tan \omega \pi \sin \omega x \end{vmatrix} = \omega \tan(\omega \pi).$$

Thus, $W = \text{Const.}$, consistent with Abel's Theorem with $p_1/p_2 = 0/1$.

Remark: This approach only works when ω is not an integer! If ω is an integer and equal to n , say, then we have $W \propto \tan(n\pi) = 0$, and we fail to generate two linearly independent solutions to the BVP. In this case, we would have to look at the **Boundary-Value Problem Alternative**, whereby the BVP has either no solutions or infinitely many solutions.

Motivated by the foregoing remark, for the remainder of this exercise we work with $\omega \notin \mathbb{Z}$. The Green's function for the BVP is now:

$$G(x, s) = \frac{[\sin(\omega x) - \tan(\omega\pi) \cos(\omega x)] \sin(\omega s)}{\omega \tan \omega\pi}, \quad 0 \leq s \leq x \leq \pi,$$

and

$$G(x, s) = \frac{\sin(\omega x) [\sin(\omega s) - \tan(\omega\pi) \cos(\omega s)]}{\omega \tan \omega\pi}, \quad 0 \leq x < s \leq \pi.$$

By convolution, the solution to the BVP is:

$$y(x) = \int_0^\pi G(x, s) f(s) ds.$$

Hence,

$$\begin{aligned} y(x) &= \int_0^x \frac{[\sin(\omega x) - \tan(\omega\pi) \cos(\omega x)] \sin(\omega s)}{\omega \tan \omega\pi} f(s) ds \\ &\quad + \int_x^\pi \frac{\sin(\omega x) [\sin(\omega s) - \tan(\omega\pi) \cos(\omega s)]}{\omega \tan \omega\pi} f(s) ds, \end{aligned}$$

or

$$\begin{aligned} y(x) &= \int_0^x \frac{[\sin(\omega x) \cos(\omega\pi) - \sin(\omega\pi) \cos(\omega x)] \sin(\omega s)}{\omega \sin \omega\pi} f(s) ds \\ &\quad + \int_x^\pi \frac{\sin(\omega x) [\sin(\omega s) \cos(\omega\pi) - \sin(\omega\pi) \cos(\omega s)]}{\omega \sin \omega\pi} f(s) ds. \end{aligned}$$

Or again,

$$y(x) = \int_0^x \frac{\sin(\omega(x - \pi)) \sin(\omega s)}{\omega \sin \omega\pi} f(s) ds + \int_x^\pi \frac{\sin(\omega x) \sin(\omega(s - \pi))}{\omega \sin \omega\pi} f(s) ds.$$

Finally,

$$y(x) = \frac{\sin(\omega(x - \pi))}{\omega \sin \omega\pi} \int_0^x \sin(\omega s) f(s) ds + \frac{\sin(\omega x)}{\omega \sin \omega\pi} \int_x^\pi \sin(\omega(s - \pi)) f(s) ds.$$

We fill in for $f(x) = 1$:

$$\begin{aligned} y(x) &= \frac{\sin(\omega(x - \pi))}{\omega \sin \omega\pi} \int_0^x \sin(\omega s) ds + \frac{\sin(\omega x)}{\omega \sin \omega\pi} \int_x^\pi \sin(\omega(s - \pi)) ds, \\ &= \frac{\sin(\omega(x - \pi))}{\omega^2 \sin \omega\pi} [1 - \cos(\omega x)] + \frac{\sin(\omega x)}{\omega^2 \sin \omega\pi} [\cos(\omega(x - \pi)) - 1] \\ &= \frac{1}{\omega^2 \sin(\omega\pi)} [\sin(\omega(x - \pi)) - \sin(\omega x)] \\ &\quad + \frac{1}{\omega^2 \sin \omega\pi} [\sin(\omega x) \cos(\omega(x - \pi)) - \cos(\omega x) \sin(\omega(x - \pi))]. \end{aligned}$$

Hence,

$$y(x) = \frac{1}{\omega^2 \sin(\omega\pi)} [\sin(\omega(x - \pi)) - \sin(\omega x)] + \frac{1}{\omega^2}. \quad (2)$$

Clearly, $y''(x) + \omega^2 y = 1$, and $y(0) = 0$ and $y(\pi) = 0$, so our calculations are correct and Equation (2) does indeed solve the BVP.

Next, we fill in for $f(x) = \sin(\omega x)$:

$$y(x) = \frac{\sin(\omega(x - \pi))}{\omega \sin \omega\pi} \int_0^x \sin(\omega s) \sin(\omega s) ds + \frac{\omega \sin(\omega x)}{\sin \omega\pi} \int_x^\pi \sin(\omega(s - \pi)) \sin(\omega s) ds.$$

This 'simplifies' to:

$$\begin{aligned} y(x) &= \frac{\sin(\omega(x - \pi))}{\omega \sin \omega\pi} \left[\frac{1}{2}x - \frac{\sin(2\omega x)}{4\omega} \right] \\ &\quad + \frac{\sin(\omega x)}{\omega \sin \omega\pi} \left[\frac{1}{2}(\pi - x) \cos(\omega\pi) - \frac{\sin(\omega(\pi - 2x))}{4\omega} - \frac{\sin(\omega\pi)}{4\omega} \right] \end{aligned} \quad (3)$$

We have checked this result using a numerical 'shooting' method (Figure 1). Matlab listings are provided in the Appendix.

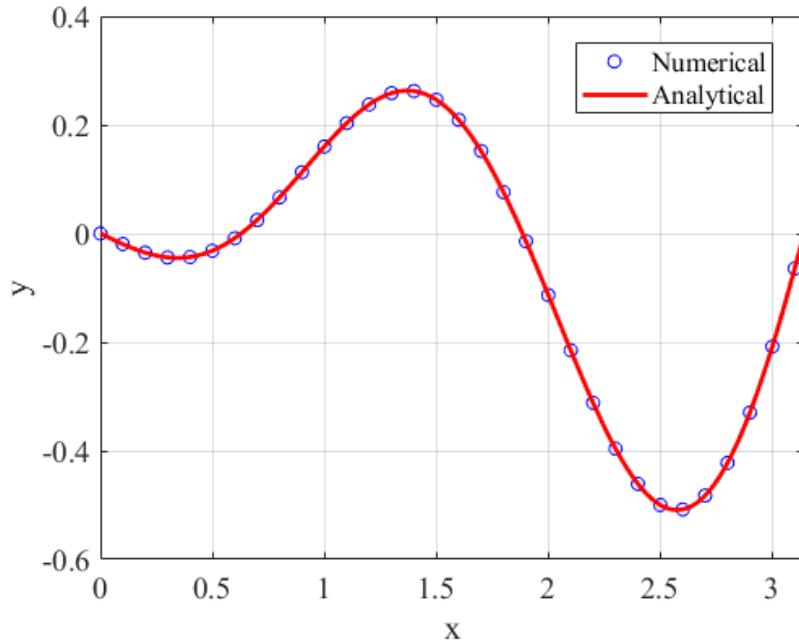


Figure 1: Comparison between numerical shooting method and Equation (3), for $\omega = 2.5$.

3. The stationary temperature distribution in a rod of unit length that has both ends kept at a constant zero temperature, with heat loss through its surface proportional to u , and that is subject to a given non-uniform heat source per unit length $f(x)$, is the solution of

$$-u'' + u = f, \quad u(0) = u(1) = 0.$$

Show that the Green's function of this Boundary Value Problem is given by:

$$G(x, \xi) = \begin{cases} \frac{\sinh(x) \sinh(1-\xi)}{\sinh(1)}, & 0 \leq x \leq \xi \leq 1, \\ \frac{\sinh(\xi) \sinh(1-x)}{\sinh(1)}, & 0 \leq \xi \leq x \leq 1. \end{cases}$$

Linearly independent solutions of the homogeneous problem are:

$$u(x) = \sinh(x)$$

and

$$v(x) = \sinh(1) \cosh(x) - \cosh(1) \sinh(x).$$

Hence, $u(x)$ satisfies the left-hand Boundary Condition (LHBC) at $x = 0$ and $v(x)$ satisfies the right-hand Boundary Condition (RHBC) at $x = 1$. The Wronskian is:

$$\begin{aligned} W &= \begin{vmatrix} u & v \\ u' & v' \end{vmatrix}, \\ &= \begin{vmatrix} \sinh(x) & \sinh(1) \cosh(x) \\ \cosh(x) & \sinh(1) \sinh(x) \end{vmatrix}, \\ &= -\sinh(1). \end{aligned}$$

Also, $p_2 = -1$, hence $p_2 W = \sinh(1)$.

Reading off from class notes, the Green's function is now:

$$G(x, \xi) = \frac{[\sinh(1) \cosh(x) - \cosh(1) \sinh(x)] \sinh(\xi)}{\sinh(1)}, \quad 0 \leq \xi \leq x \leq 1.$$

The term in the square bracket can be written as $\sinh(1-x)$, using a trig identity. Also,

$$G(x, \xi) = \frac{\sinh(x) [\sinh(1) \cosh(\xi) - \cosh(1) \sinh(\xi)]}{\sinh(1)}, \quad 0 \leq x \leq \xi \leq 1.$$

Again, a trig identity can be used on the term in the square bracket, this is $\sinh(1-\xi)$. Hence,

$$G(x, \xi) = \begin{cases} \frac{\sinh(\xi) \sinh(1-x)}{\sinh(1)}, & 0 \leq \xi \leq x \leq 1, \\ \frac{\sinh(x) \sinh(1-\xi)}{\sinh(1)}, & 0 \leq x \leq \xi \leq 1, \end{cases}$$

as required.

4. Consider the Fredholm integral equation

$$y(x) = f(x) + \lambda \int_{-1}^1 (x+s)y(s)ds.$$

- (a) For which values of λ does the equation have a unique solution? Find the solution in this case.
- (b) For each of those values of λ for which the equation does not have a unique solution, state a condition which $f(x)$ must satisfy in order for a solution to exist, and find the general solution when this is satisfied.

We have:

$$K(x, s) = u_1(x)v_1(s) + u_2(x)v_2(s),$$

where $u_1(x) = x$, $v_1(s) = 1$, $u_2(x) = 1$, and $v_2(s) = 1$. Thus, the FIE can be written as:

$$y(x) = f(x) + \lambda x c_1 + \lambda c_2 \quad (4)$$

where

$$c_1 = \int_{-1}^1 y(s)ds, \quad c_2 = \int_{-1}^1 s y(s)ds.$$

Let

$$f_1 = \int_{-1}^1 f(s)ds, \quad f_2 = \int_{-1}^1 s f(s)ds.$$

We multiply the FIE (4) by v_1 and integrate to obtain:

$$c_1 = f_1 + 2\lambda c_2. \quad (5a)$$

Next, we multiply the FIE (4) by v_2 and integrate to obtain:

$$c_2 = f_2 + 2\lambda c_1/3. \quad (5b)$$

In doing these integrals, it is helpful to remember the properties of integrals of odd/even functions over symmetric intervals, e.g.

$$\int_{-1}^1 x dx = 0,$$

etc. Equations (5) can be simplified and written in matrix form:

$$\underbrace{\begin{pmatrix} 1 & -2\lambda \\ -2\lambda/3 & 1 \end{pmatrix}}_{=M(\lambda)} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$$

The characteristic polynomial is

$$\det [M(\lambda)] = 0,$$

hence

$$1 - \frac{4}{3}\lambda^2 = 0.$$

Thus, the roots of the characteristic polynomial are given by:

$$\lambda \in \mathcal{S} = \{\lambda_-, \lambda_+\} = \{-\frac{\sqrt{3}}{2}, \frac{\sqrt{3}}{2}\}.$$

Thus, the answer to **part (a)** is that there is a unique solution if $\lambda \neq \mathcal{S}$.

For **part (b)**, we look at the two separate cases.

- Case 1, $\lambda = -\sqrt{3}/2$. Here, we require $f_1 = \sqrt{3}f_2$, that is

$$\int_{-1}^1 f(s)(1 - \sqrt{3}x)dx = 0,$$

in which case the general solution is given for arbitrary f_2 by

$$\begin{aligned} y(x) &= f(x) + \lambda_- \left[\left(-\sqrt{3}c_2 + b_1 \right) x + c_2 \right], \\ &= f(x) + \lambda_- b_1 x + \lambda_- c_2 \underbrace{\left(1 - \sqrt{3}x \right)}_{=y_-}. \end{aligned}$$

Thus, the general solution is:

$$y(x) = f(x) + \lambda_- b_1 x + \alpha y_-(x),$$

where $y_-(x)$ satisfies the homogeneous FIE:

$$y_-(x) = \lambda_- \int_{-1}^1 K(x, s)y_-(s)ds.$$

The compatibility condition is:

$$\int_{-1}^1 y_-(x)f(x)dx = 0.$$

- Case 2, $\lambda = \sqrt{3}/2$. Here, we require $f_1 = -\sqrt{3}f_2$, that is

$$\int_{-1}^1 f(s)(1 + \sqrt{3}x)dx = 0,$$

in which case the general solution is given for arbitrary f_2 by

$$\begin{aligned} y(x) &= f(x) + \lambda_+ \left[\left(\sqrt{3}c_2 + b_1 \right) x + c_2 \right], \\ &= f(x) + \lambda_+ b_1 x + \lambda_+ c_2 \underbrace{\left(1 + \sqrt{3}x \right)}_{=y_-}. \end{aligned}$$

Thus, the general solution is:

$$y(x) = f(x) + \lambda_+ b_1 x + \alpha y_+(x),$$

where $y_+(x)$ satisfies the homogeneous FIE:

$$y_+(x) = \lambda_- \int_{-1}^1 K(x, s) y_+(s) ds.$$

The compatibility condition is:

$$\int_{-1}^1 y_+(x) f(x) dx = 0.$$

Remark: When we look at Hilbert–Schmidt theory, we will find, for $\lambda \notin \mathcal{S}$, that the solution of the FIE can be written as:

$$y(x) = f(x) + \frac{\lambda}{\lambda_+ - \lambda} f_+ y_+(x) + \frac{\lambda}{\lambda_- - \lambda} f_- y_-(x),$$

where

$$\lambda_{\pm} = \pm \sqrt{32}.$$

Also,

$$y_{\pm}(x) = \frac{1 \pm \sqrt{3}x}{2}$$

are the *normalized* eigenfunctions of the FIE, with

$$\int_{-1}^1 y_i(x) y_j(x) dx = \delta_{ij}, \quad i, j \in \{-, +\},$$

Furthermore,

$$f_{\pm} = \int_{-1}^1 f(x) y_{\pm}(x) dx.$$

Thus, we see transparently that in part (b), the condition which f must satisfy in order for a solution to exist when $\lambda = \lambda_+$ (respectively $\lambda = \lambda_-$) is $f_+ = 0$ (respectively $f_- = 0$).

5. Solve for $\phi(x)$ in the integral equation

$$\phi(x) = f(x) + \lambda \int_0^1 \left[\left(\frac{x}{y} \right)^n + \left(\frac{y}{x} \right)^n \right] \phi(y) dy,$$

where $f(x)$ is bounded for $0 < x < 1$ and $-1/2 < n < 1/2$, expressing your answer in terms of the quantities $F_m = \int_0^1 f(y) y^m dy$.

- (a) Give the explicit solution when $\lambda = 1$.
- (b) For what values of λ are there no solutions unless $F_{\pm n}$ are in a particular ratio? What is this ratio?

Solution: We are looking at:

$$\phi(x) = f(x) + \lambda \int_0^1 K(x, y) \phi(y) dy,$$

hence a Fredholm integral equation with kernel $K(x, y) = (x/y)^n + (y/x)^n$. This is a separable kernel,

$$K(x, s) = u_1(x)v_1(s) + u_2(x)v_2(s),$$

with $u_1(x) = x^n$, $v_1(s) = s^{-n}$, $u_2(x) = x^{-n}$, and $v_2(s) = s^n$. We have:

$$\begin{aligned} c_1 &= \int_0^1 v_1(s) \phi(s) ds = \int_0^1 s^{-n} \phi(s) ds, \\ c_2 &= \int_0^1 v_2(s) \phi(s) ds = \int_0^1 s^n \phi(s) ds. \end{aligned}$$

Also (using notation from class notes):

$$\begin{aligned} b_1 &= \int_0^1 v_1(s) f(s) ds = \int_0^1 s^{-n} f(s) ds = F_{-n}, \\ b_2 &= \int_0^1 v_2(s) f(s) ds = \int_0^1 s^n f(s) ds = F_n. \end{aligned}$$

We project the integral equation on to v_1 and v_2 to get:

$$c_i = b_i + \lambda \sum_{j=1}^2 A_{ij} c_j,$$

where $A_{ij} = \int_0^1 v_i(s)u_j(s)ds$. Hence, it remains to compute:

$$\begin{aligned} A_{11} &= \int_0^1 1 \, ds = 1, \\ A_{12} &= \int_0^1 s^{-2n} \, ds = \frac{1}{1-2n}, \\ A_{21} &= \int_0^1 s^{2n} \, ds = \frac{1}{2n+1}, \\ A_{22} &= \int_0^1 1 \, ds = 1. \end{aligned}$$

We don't encounter any divisions by zero here because $|n| < 1/2$. Thus, the linear problem to solve is:

$$\underbrace{\begin{pmatrix} 1-\lambda & -\frac{\lambda}{1-2n} \\ -\frac{\lambda}{1+2n} & 1-\lambda \end{pmatrix}}_M \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}. \quad (6)$$

We focus on the characteristic equation:

$$\begin{vmatrix} 1-\lambda & -\frac{\lambda}{1-2n} \\ -\frac{\lambda}{1+2n} & 1-\lambda \end{vmatrix} = 0.$$

The solutions are:

$$\lambda_+ = \frac{\sqrt{1-4n^2}}{1+\sqrt{1-4n^2}}, \quad \lambda_- = \frac{\sqrt{1-4n^2}}{-1+\sqrt{1-4n^2}}.$$

Thus, when $\lambda = 1$, it is the case that $\lambda \neq \lambda_{\pm}$, and hence, Equation (6) has a unique solution, characterized by:

$$M^{-1} = \begin{pmatrix} 0 & -(2n+1) \\ 2n-1 & 0 \end{pmatrix}.$$

Hence,

$$\begin{aligned} (c_1, c_2)^T &= \begin{pmatrix} 0 & -(2n+1) \\ 2n-1 & 0 \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix} \\ &= \begin{pmatrix} 0 & -(2n+1) \\ 2n-1 & 0 \end{pmatrix} \begin{pmatrix} b_1 \\ b_2 \end{pmatrix}. \end{aligned}$$

Hence,

$$c_1 = -(1+2n)b_2, \quad c_2 = -(1-2n)b_1.$$

The solution in the case $\lambda = 1$ is therefore:

$$\begin{aligned} \phi_{\lambda=1}(x) &= f(x) + c_1 u_1(x) + c_2 u_2(x), \\ &= f(x) - (1+2n)b_2 x^n - (1-2n)b_1 x^{-n}. \end{aligned}$$

Or, in terms of the notation in the question,

$$\phi_{\lambda=1}(x) = f(x) - (1 + 2n)F_n x^n - (1 - 2n)F_{-n} x^{-n}.$$

For part (b), we look again at Equation (6) with $\lambda = \lambda_{\pm}$. In an obvious notation, we have:

$$M(\lambda_{\pm})\mathbf{c} = \mathbf{b}, \quad (7)$$

and $\det M(\lambda_{\pm}) = 0$. Hence, for almost all values of n , Equation (6) with $\lambda = \lambda_{\pm}$ will have no solution.

However, there is one special case to look at. Hence, we re-write Equation (7) as:

$$\begin{aligned} M(\lambda_{\pm})\mathbf{c} &= \mathbf{b}, \\ \lambda_{\pm} \begin{pmatrix} \frac{1}{\lambda_{\pm}} - 1 & -\frac{1}{1-2n} \\ -\frac{1}{1+2n} & \frac{1}{\lambda_{\pm}} - 1 \end{pmatrix} \mathbf{c} &= \mathbf{b}. \end{aligned}$$

Hence:

$$\lambda_{\pm} \begin{pmatrix} \pm \frac{1}{\sqrt{1-4n^2}} & -\frac{1}{1-2n} \\ -\frac{1}{1+2n} & \pm \frac{1}{\sqrt{1-4n^2}} \end{pmatrix} \mathbf{c} = \mathbf{b}.$$

In other words,

$$\begin{aligned} \pm \frac{1}{\sqrt{1-4n^2}} c_1 - \frac{1}{1-2n} c_2 &= b_1/\lambda_{\pm}, \\ (\mp) \frac{1+2n}{\sqrt{1-4n^2}} \times \left[-\frac{1}{1+2n} c_1 \pm \frac{1}{\sqrt{1-4n^2}} c_2 \right] &= (\mp) \frac{1+2n}{\sqrt{1-4n^2}} \times b_2/\lambda_{\pm}. \end{aligned}$$

The second row can now be re-written as:

$$\pm \frac{1}{\sqrt{1-4n^2}} c_1 - \frac{1}{1-2n} c_2 = \mp \sqrt{\frac{1+2n}{1-2n}} (b_2/\lambda_{\pm}).$$

Compare the two rows again:

$$\begin{aligned} \pm \frac{1}{\sqrt{1-4n^2}} c_1 - \frac{1}{1-2n} c_2 &= b_1/\lambda_{\pm}, \\ \pm \frac{1}{\sqrt{1-4n^2}} c_1 - \frac{1}{1-2n} c_2 &= \mp \sqrt{\frac{1+2n}{1-2n}} b_2/\lambda_{\pm}. \end{aligned}$$

If the two rows are the same, there are infinitely many solutions. Otherwise, there are no solutions. Hence, to find the case with solutions, we require:

$$\frac{b_1}{b_2} = \mp \sqrt{\frac{1+2n}{1-2n}}.$$

Lastly, go back to the notation in the question ($b_1 \rightarrow F_{-n}$, $b_2 \rightarrow F_n$):

$$\frac{F_n}{F_{-n}} = \mp \sqrt{\frac{1-2n}{1+2n}}.$$

6. Consider the FIE

$$y(x) = f(x) + \lambda \int_0^1 \cosh(x-s)y(s)ds. \quad (8)$$

- (a) Show that the eigenvalues of (8) are given by $2/(1 \pm \sinh 1)$.
- (b) Using the Hilbert–Schmidt eigenfunction expansion, or otherwise, find the solution for $\lambda \notin \{2/(1 \pm \sinh 1)\}$.
- (c) Find a necessary and sufficient condition on f for the equation

$$y(x) = f(x) + \frac{2}{1 + \sinh 1} \int_0^1 \cosh(x-s)y(s)ds$$

to have a solution and find all solutions when this condition is satisfied.

For part (a), note that the kernel is separable, since

$$K(x, s) = \cosh(x-s) = \frac{1}{2}e^x e^{-s} + \frac{1}{2}e^{-x} e^s.$$

To find the eigenvalues, let us write the homogeneous equation as

$$y(x) = \frac{1}{2}e^x \lambda c_1 + \frac{1}{2}e^{-x} \lambda c_2,$$

where

$$c_1 = \int_0^1 e^{-s} y(s) ds, \quad c_2 = \int_0^1 e^s y(s) ds.$$

Then

$$\begin{aligned} c_1 &= \int_0^1 e^{-x} y(x) dx = \frac{1}{2} \lambda c_1 \int_0^1 e^{-x} e^x dx + \frac{1}{2} \lambda c_2 \int_0^1 e^{-x} e^{-x} dx \\ &= \frac{1}{2} \lambda c_1 + \frac{1}{2} \lambda c_2 \frac{1}{2} (1 - e^{-2}), \end{aligned}$$

and

$$\begin{aligned} c_2 &= \int_0^1 e^x y(x) dx = \frac{1}{2} \lambda c_1 \int_0^1 e^x e^x dx + \frac{1}{2} \lambda c_2 \int_0^1 e^x e^{-x} dx \\ &= \frac{1}{2} \lambda c_1 \frac{1}{2} (e^2 - 1) + \frac{1}{2} \lambda c_2. \end{aligned}$$

These equations may be written in matrix form as

$$\underbrace{\begin{pmatrix} \lambda - 2 & \lambda \frac{1}{2} (1 - e^{-2}) \\ \lambda \frac{1}{2} (e^2 - 1) & \lambda - 2 \end{pmatrix}}_{=M(\lambda)} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}. \quad (9)$$

The characteristic polynomial is:

$$\det [M(\lambda)] = 0.$$

This evaluates to:

$$(\lambda - 2)^2 = \frac{1}{4}\lambda^2(e^2 - 1)(1 - e^{-2}) = \frac{1}{4}\lambda^2(e - e^{-1})^2 = \lambda^2 \sinh^2(1).$$

Hence, the eigenvalues of the FIE are:

$$\lambda \in \mathcal{S} = \{\lambda_1, \lambda_2\} = \left\{ \frac{2}{1 + \sinh 1}, \frac{2}{1 - \sinh 1} \right\}.$$

For **part (b)**, the pertinent eigenfunctions are obtained by solving for $(c_1, c_2)^T$ in Equation (9). We look at the two cases:

- Case 1. We look at $\lambda = \lambda_1$. Notice that $\lambda_1 \sinh 1 = -(\lambda_1 - 2)$ so, for $\lambda = \lambda_1$, Equation (9) becomes:

$$\begin{pmatrix} \lambda_1 - 2 & -(\lambda_1 - 2)e^{-1} \\ -(\lambda_1 - 2)e & \lambda_1 - 2 \end{pmatrix} \begin{pmatrix} c_1 \\ c_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}.$$

This implies that $c_2 = ec_1$, so the corresponding eigenfunction is a multiple of

$$e^x + ee^{-x} = 2e^{1/2} \frac{1}{2}(e^{x-1/2} + e^{-(x-1/2)}),$$

or equivalently:

$$y_1(x) \propto \cosh(x - \frac{1}{2}).$$

- For $\lambda = \lambda_2$, we have $c_2 = -ec_1$, and the corresponding eigenfunction is

$$y_2(x) \propto \sinh(x - \frac{1}{2}).$$

We look at normalized eigenfunctions:

$$y_1(x) = \sqrt{\lambda_1} \cosh(x - \frac{1}{2}), \quad y_2(x) = \sqrt{-\lambda_2} \sinh(x - \frac{1}{2}).$$

Notice that:

$$\int_0^1 y_i(x) y_j(x) dx = \delta_{ij}.$$

In this case, Hilbert–Schmidt theory tells us that the solution to the FIE can be expanded in terms of eigenfunctions:

$$y(x) = f(x) + \frac{\lambda}{\lambda_1 - \lambda} f_1 y_1(x) + \frac{\lambda}{\lambda_2 - \lambda} f_2 y_2(x), \quad \lambda \notin \mathcal{S}.$$

Here,

$$f_i = \int_0^1 f(x) y_i(x) dx.$$

For **part (c)**, Hilbert–Schmidt theory tells us that if $\lambda \in \mathcal{S}$ (in this case, $\lambda = \lambda_1$), then a necessary necessary and sufficient condition on f for the FIE to have a solution is:

$$\int_0^1 f(x) y_1(x) dx = 0 \implies f_1 = 0.$$

In this case, the general solution is

$$y(x) = f(x) + A y_1(x) + \frac{\lambda_1}{\lambda_2 - \lambda_1} f_2 y_2(x).$$

A Code Listings

In this code, I solve the ODE $y''(x) = -\omega^2 y(x) + \sin(x)$ for the initial conditions $y(0) = 0$ and $y'(0) = p$. Here, p is a parameter that can be varied.

```
function [x,y]=odesolve(param)
omega=2.5;
[x,Y]=ode45(@myfun,[0 pi], [0 param]);
y=Y(:,1);

function dYdx=myfun(x,Y)
dYdx=0*Y;
dYdx(1)=Y(2);
dYdx(2)=-omega*omega*Y(1)+sin(omega*x);
end

end
```

In this next code, I look at the cost function

$$J(p) = [y(x = \pi; p)]^2,$$

where $y(x = \pi, p)$ is output from `odesolve.m`, at parameter value p . I solve an optimization problem where I make $J(p)$ as small as possible. In this way, I solve the BVP $y''(x) = -\omega^2 y(x) + \sin(x)$ subject to $y(0) = 0$ and $y(\pi) = 0$. This is called the **shooting method**.

```
p_star=fminbnd(@mycost,-10 ,10000);
display(p_star);
J_min=mycost(p_star);
display(J_min);

% Obtain solution of BVP:
[x,y]=odesolve(p_star);

% Interpolate solution on to regular grid:
xi=0:0.1:pi;
yi=interp1(x,y,xi,'spline');

x=xi;
y=yi;
```

```
% Plot solution:  
plot(x,y, 'o', 'color ', 'blue')  
  
% Calculate analytical solution for omega=2.5:  
  
omega=2.5;  
xa=0:0.01:pi;  
ya=sin(omega*(xa-pi)).*(0.5*xa-(sin(2*omega*xa)/(4*omega)))...  
+sin(omega*xa).*(0.5*(pi-xa)*cos(omega*pi))...  
-(sin(omega*(pi-2*xa))/(4*omega))  
-(sin(omega*pi)/(4*omega));  
ya=ya/(omega*sin(omega*pi));  
  
% Plot analytical and numerical solution on same axes  
% and compare:  
  
hold on  
plot(xa, ya, 'color ', 'red', 'linewidth', 2)  
grid on  
xlabel('x')  
ylabel('y')  
set(gca, 'fontsize', 14, 'fontname', 'times new roman')  
hold off  
  
  
  
function J=mycost(p)  
[~,y]=odesolve(p);  
  
J=y(end)*y(end);  
end
```