Applied Analysis (ACM30020)

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Exercises #4

1. Consider the ODE

$$y'' + p(x)y' + q(x)y = 0.$$

If $y_1(x)$ is a solution, show that a second solution can be written as:

$$y_2(x) = y_1(x) \int_a^x \frac{\mathrm{e}^{-\int_b^{x''} p(x')\mathrm{d}x'}}{[y_1(x'')]^2} \mathrm{d}x''.$$
(1)

Here, a and b are arbitrary.

Sub $y_2(x) = y_1(x)u(x)$ into the ODE to obtain:

$$u'' + \left[p(x) + 2\frac{y'_2}{y_2} \right] u' = 0.$$
⁽²⁾

Let v = u' and reduce Equation (2) to a first-order ODE:

$$v' = -\left[p(x) + 2\frac{y_2'}{y_2}\right]v.$$

The solution is:

$$v(x) = v(b) \frac{\mathrm{e}^{\int_{b}^{x} p(x')\mathrm{d}x]}}{[y_{2}(x)]^{2}},$$

where b is arbitrary. But v(x) = du/dx, hence

$$u(x) = C + v(b) \int_{a}^{x} \frac{e^{-\int_{b}^{x''} p(x')dx'}}{[y_{1}(x'')]^{2}} dx''$$

where C is a constant of integration.

Choose C = 0 and v(b) = 1 to get a second linearly-independent solution:

$$u(x) = \int_{a}^{x} \frac{\mathrm{e}^{-\int_{b}^{x''} p(x')\mathrm{d}x'}}{[y_{1}(x'')]^{2}} \mathrm{d}x''.$$

Hence,

$$y_2(x) = y_1(x) \int_a^x \frac{e^{-\int_b^{x''} p(x') dx'}}{[y_1(x'')]^2} dx''$$

as required.

2. Given that one solution of

$$R'' + \frac{1}{r}R' - \frac{m^2}{r^2}R = 0$$

is $R = r^m$, show that Equation (1) provides a second solution, $R = r^{-m}$.

We read off from Question 1, with p(x) = 1/x. Hence, $\int p(x)dx = \ln x$, and $e^{-\int p(x)dx} = -1/x$. Thus, the second solution (in an obvious notation, and letting $x \to r$) is:

$$R_2(r) = r^m \int \frac{1}{r} \frac{1}{r^{2m}} dr,$$

$$= r^m \left(-\frac{1}{2m} r^{-2m} \right),$$

$$\propto r^{-m}.$$

As the solutions are only defined up to a constant, we choose the second linearly independent solution to be:

$$R_2(r) = r^{-m},$$

as required.

3. Consider Legendre's differential equation:

$$(1 - x2)y'' - 2xy' + n(n+1)y = 0.$$
 (3)

- (a) Solve the equation by by direct series substitution.
- (b) Verify that the indicial equation is:

$$\alpha(\alpha - 1) = 0.$$

(c) Using $\alpha = 0$, obtain the following series of even powers of x $(a_1 = 0)$:

$$y_{even} = a_0 \left[1 - \frac{n(n+1)}{2!} x^2 + \frac{n(n-2)(n+1)(n+3)}{4!} x^4 + \cdots \right],$$

where

$$a_{j+2} = \frac{j(j+1) - n(n+1)}{(j+1)(j+2)}a_j.$$

(d) Using $\alpha = 1$, develop a series of odd powers of x ($a_1 = 0$).

$$y_{odd} = a_0 \left\{ x - \frac{(n-1)(n+2)}{3!} x^3 + \frac{(n-1)(n-3)(n+2)(n+4)}{5!} x^5 + \cdots \right\},$$

where

$$a_{j+2} = \frac{(j+1)(j+2) - n(n+1)}{(j+2)(j+3)}a_j.$$

- (e) Show that both solutions, y_{even} and y_{odd} , diverge for $x = \pm 1$ if the series continue to infinity.
- (f) Finally, show that by an appropriate choice of n, one series at a time may be converted into a polynomial, thereby avoiding the divergence catastrophe.

Because the coefficient functions $p(x) = -2x/(1-x^2)$ and $q(x) = n(n+1)/(1-x^2)$ are regular at x = 0, this is a regular point and hence, the ODE has a simple power-series solution:

$$y(x) = \sum_{p=0}^{\infty} a_p x^p$$

(we can't use n for the index because it's used already as the parameter in the ODE). We substitute this into Equation (3) to get:

$$\sum_{p=0}^{\infty} a_p p(p-1) x^{p-2} = \sum_{p=0}^{\infty} a_p p(p-1) x^p + 2 \sum_{p=0}^{\infty} a_p p x^p - \sum_{p=0}^{\infty} n(n+1) a_p x^p.$$

Re-index. In the term on the LHS we use q = p - 2. In the terms on the RHS we use q = p. This gives:

$$\sum_{q=-2}^{\infty} a_{q+2}(q+2)(q+1)x^q = \sum_{q=0}^{\infty} a_q q(q-1)x^q + 2\sum_{q=0}^{\infty} a_q qx^q - \sum_{q=0}^{\infty} n(n+1)a_q x^q.$$

We equate powers of x^q . At q = -2 we have $a_0 \times 0 \times 1 = 0$, which implies that a_0 is arbitrary. At q = -1 we have $a_1 \times (-1) \times 0 = 0$, meaning that a_1 is arbitrary also. We also have the recurrence relation:

$$a_{q+2} = a_q \frac{q(q-1) + 2q - n(n+1)}{(q+2)(q+1)},$$

= $a_q \frac{q(q+1) - n(n+1)}{(q+2)(q+1)}$

These results mean that we can look at the possibilities $\{a_0 \neq 0, a_1 = 0\}$ and $\{a_0 = 0, a_1 \neq 0\}$ separately. Furthermore, because the recurrence relation steps up in steps of two, this means that we are looking at odd and even series solutions.

Consequently, the series solutions are $y(x) = a_0 + a_2x^2 + a_4x^4 + \cdots$ and $y(x) = a_1x + a_3x^3 + \cdots$, which can be encapsulated as $y(x) = x^{\alpha}(a_0 + a_1x + a_2x^2 + \cdots)$, with $\alpha = 0, 1$. This implies an indicial equation $\alpha(\alpha - 1) = 0$, which answers **Part** (b).

Parts (a) and (c): We look at the possibilities $\{a_0 \neq 0, a_1 = 0\}$ and $\{a_0 = 0, a_1 \neq 0\}$ separately, and generate odd and even series solutions. We use the recurrence relation to generate the first few terms of the even solution:

$$a_{2} = a_{0} \frac{-n(n+1)}{2!},$$

$$a_{4} = a_{2} \left[\frac{6 - n(n+1)}{4 \times 3} \right],$$

$$= a_{0} \frac{1}{4!} (-1) \left[6 - n(n+1) \right] n(n+1),$$

$$= a_{0} \frac{1}{4!} n(n+1)(n-2)(n+3).$$

Hence,

$$y_{even} = a_0 \left[1 - \frac{n(n+1)}{2!} x^2 + \frac{n(n+1)(n-2)(n+3)}{4!} x^4 + \cdots \right],$$

where

$$a_{q+2} = a_q \frac{q(q+1) - n(n+1)}{(q+2)(q+1)}.$$

Thus, the answer to Part (c) is complete.

Part (d) concerns the odd solution. The first few terms are:

$$a_{3} = a_{1} \frac{2 - n(n+1)}{3 \times 2},$$

$$= a_{1} \left[-\frac{(n-1)(n+2)}{3!} \right],$$

$$a_{5} = a_{3} \left[\frac{3 \times 4 - n(n+1)}{5 \times 4} \right],$$

$$= a_{1} \left[-\frac{(n-1)(n+2)}{3!} \right] \left[-\frac{(n-3)(n+4)}{5 \times 4} \right]$$

$$= a_{1} \left[\frac{(n-1)(n-3)(n+2)(n+4)}{5!} \right].$$

Hence, we obtain the odd solution:

$$y_{odd} = a_1 \left[x - \frac{(n-1)(n+2)}{3!} x^3 + \frac{(n-1)(n-3)(n+2)(n+4)}{5!} x^5 + \cdots \right].$$

Again, the ratio between successive terms is given by the recurrence relation:

$$a_{q+2} = a_q \frac{q(q+1) - n(n+1)}{(q+2)(q+1)}$$

For **Part** (e), we look again at the recurrence relation in the limit of large q. This gives

$$\frac{a_{q+2}}{a_q} \sim \frac{q}{q+2} \sim 1.$$

Thus, the tail of the series looks like a geometric progression with alternating sign. The geometric progression with alternating sign is:

$$G(x) = \sum_{q=0}^{\infty} (-1)^q x^q,$$

which diverges as $|x| \to 1$. Thus, and in general, the series y_{even} and y_{odd} will diverge as $|x| \to 1$. This answers Part (e).

For **Part (f)**, we notice that the series y_{even} and y_{odd} will terminate and reduce to polynomial expressions if n is zero or a positive integer, since then the expression

$$j(j+1) = n(n+1)$$

will be satisfied for n = j, and hence $a_{j+2} = 0$. These are the Legendre Polynomials (Figure 1).



Figure 1: The first 6 Legendre Polynomials. From Wikipedia.

4. Obtain two series solutions of the confluent hypergeometric equation

$$xy'' + (c - x)y' - ay = 0.$$

Test your solutions for convergence.

We first look at x[(c-x)/x] = c - x and $x^2[-a/x] = -ax$. The RHS of both these expressions have Taylor expansions around zero, so the singular point x = 0 is regular. Thus, a series solution

$$y(x) = x^{\alpha} \sum_{n=0}^{\infty} a_n x^n,$$

is possible. We substitute this trial solution into the ODE and evaluate:

$$\sum_{n=0}^{\infty} a_n (n+\alpha)(n+\alpha-1)x^{n+\alpha-1} + c \sum_{n=0}^{\infty} a_n (n+\alpha)x^{n+\alpha-1} - \sum_{n=0}^{\infty} a_n (n+\alpha)x^{n+\alpha} - a \sum_{n=0}^{\infty} a_n x^{n+\alpha} = 0.$$

We cancel out a power of x^{α} on both sides. Hence, we have:

$$\sum_{n=0}^{\infty} a_n (n+\alpha)(n+\alpha-1)x^{n-1} + c \sum_{n=0}^{\infty} a_n (n+\alpha)x^{n-1} - \sum_{n=0}^{\infty} a_n (n+\alpha)x^n - a \sum_{n=0}^{\infty} a_n x^n = 0.$$

We then re-index with p = n - 1, hence n = p + 1, and $p_{start} = -1$. Hence, we have:

$$\sum_{p=-1}^{\infty} a_{p+1}(p+1+\alpha)(p+\alpha)x^p + c\sum_{p=-1}^{\infty} a_{p+1}(p+1+\alpha)x^p - \sum_{p=0}^{\infty} a_p(p+\alpha)x^p - a\sum_{p=0}^{\infty} a_px^p = 0.$$

For the p = -1 term we have:

$$a_0 \left[\alpha(\alpha - 1) + c\alpha \right] = 0.$$

Hence, the indicial equation is:

$$\alpha(\alpha - 1) + c\alpha = 0,$$

with solutions

$$\alpha = 0, \qquad \alpha = 1 - c.$$

We look at the analytic solution with $\alpha=0.\,$ We look in particular at the recurrence relation:

$$a_{p+1}[p(p+1) + c(p+1)] = a_p(p+a).$$

or

$$a_n n(n-1+c) = a_{n-1} (n-1+a), \qquad n \ge 1.$$

Hence,

$$a_n = \frac{(n-1+a)}{n(n-1+c)} a_{n-1},$$
(4)

Furthermore,

$$a_n = \frac{(n-1+a)(n-2+a)}{n(n-1)(n-1+c)(n-2+c)}a_{n-2},$$

= $\frac{(n-1+a)(n-2+a)\cdots(1+a)a}{n!(n-1+c)(n-2+c)\cdots(1+c)c}a_0.$

Notice that the recurrence relation will terminate if a is a negative integer or zero (giving a polynomial solution), and also, the recurrence relation will fail if c is a negative integer or zero.

We identify the Pochammer symbol (rising factorial):

$$(a)^n = a(a+1)\cdots(a+n-1) = \prod_{k=0}^{n-1}(a+k),$$

and similarly for c. Hence, the first series solution can be written as:

$$y_1(x) = a_0 \sum_{n=0}^{\infty} \frac{(a)^n}{n!(c)^n} x^n, \qquad a, c \notin -\mathbb{N} \cup \{0\}.$$

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or

$$y_1(x) = a_0 \left[{}_1F_1(a;c;x) \right].$$
(5)

where ${}_{1}F_{1}(a;c;x)$ is the confluent hypergeometric function.

Furthermore, from Equation (4), we identify the ratio

$$\rho_n = \left| \frac{a_n}{a_{n-1}} \right| = \left| \frac{n-1+a}{n(n-1+c)} \right|.$$

Hence,

$$\rho = \lim_{n \to \infty} \rho_n = 0.$$

The radius of convergence of the series is therefore $R = 1/\rho = \infty$. Thus, except for $c = 0, -1, -2, \cdots$, the series in Equation (5) is an entire function, that is, an analytic on the entire complex plane.

For the second solution, we look at $\alpha = 1 - c$. The recurrence relation is:

$$a_n = \frac{n + (a - c)}{n(n + 1 - c)} a_{n-1},$$

$$= \frac{n - 1 + (a - c + 1)}{n[n - 1 + (2 - c)]} a_{n-1},$$

$$= \frac{[n - 1 + (a - c + 1)][n - 2 + (a - c + 1)]}{n(n - 1)[n - 1 + (2 - c)][n - 2 + (2 - c)]} a_{n-2},$$

$$= \cdots$$

$$= a_0 \frac{(a - c + 1)^n}{n!(2 - c)^n} a_0.$$

Hence, the second solution is:

$$y_2(x) = x^{1-c} \left[{}_1F_1(1+a-c;2-c;x) \right], \qquad 1+a-c, 2-c \notin -\mathbb{N} \cup \{0\}.$$

5. Bessel's equation can be written as

$$x^{2}y'' + xy' + (x^{2} - \nu^{2})y = 0.$$

Using power series, find the two linearly independent solutions of Bessel's equation with $\nu = 1/2$.

We use class notes to find the recurrence relations:

$$\alpha^2 - \nu^2 = 0,$$

 $a_1 (\pm 2\nu + 1) = 0.$

Also,

$$a_n n (n \pm 2\nu) = -a_{n-2}, \qquad n \ge 2.$$
 (6)

The indicial equation is $\alpha^2 = \nu^2$, hence $\alpha = \pm 1/2$. There are two distinct roots of the inidical equation.

Next, we have $a_1(2\alpha + 1) = 0$. For $\alpha = 1/2$ this gives $2a_1 = 0$, so $a_1 = 0$ in this case. For $\alpha = -1/2$ this gives $a_10 = 0$, so a_1 is undetermined in this second case. **First Case:** We have $\alpha = 1/2$, and the recurrence relation is well defined for all $n \ge 2$:

$$a_n = -\frac{a_{n-2}}{n(n+1)}, \qquad n \ge 2$$

This gives:

$$a_{2n} = -\frac{1}{2n(2n+1)}a_{2(n-1)}$$

= $\frac{1}{2n(2n+1)}\frac{1}{(2n-2)(2n-1)}a_{2(n-2)} = \cdots (-1)^n \frac{1}{(2n+1)!}a_0,$

so we have first solution

$$J_{\frac{1}{2}}(z) = \left(\frac{z}{2}\right)^{1/2} \left(1 - \frac{1}{3!}z^2 + \frac{1}{5!}z^4 - \cdots\right)$$

where, following convention, we have taken $a_0 = 2^{-1/2}$. Furthermore, we take:

$$J_{\frac{1}{2}}(z) = \left(\frac{z}{2}\right)^{1/2} \left(1 - \frac{1}{3!}z^2 + \frac{1}{5!}z^4 - \cdots\right),$$

$$= \left(\frac{z}{2}\right)^{1/2} \frac{1}{z} \left(z - \frac{1}{3!}x^2 + \frac{1}{5!}z^5 - \cdots\right),$$

$$= \frac{1}{\sqrt{2}} \frac{\sin z}{z^{1/2}}.$$

Second Case: For $\alpha = -\frac{1}{2}$, Equation (6) gives:

$$a_n n (n-1) = -a_{n-2}, \qquad n \ge 2.$$
 (7)

Here, the recurrence relation holds up for $n \ge 2$, so we can proceed to develop a regular series expansion for the second solution.

For the even powers of n, set $n \rightarrow 2n$ in Equation (7):

$$a_{2n} = -\frac{1}{2n(2n-1)}a_{2n-2}$$

Hence,

$$a_{2n} = (-1)^n \frac{1}{(2n)!} a_0.$$
(8)

We know for the second case that a_1 is arbitrary. Furthermore, for the odd powers of n we get the following recurrence relation:

$$a_{2n+1} \stackrel{\text{Eq. (7)}}{=} -\frac{1}{(2n+1)(2n)}a_{(2n+1)-2},$$
$$= -\frac{1}{(2n+1)(2n)}a_{2n-1},$$
$$= \cdots,$$
$$= \frac{(-1)^n}{(2n+1)!}a_1.$$

Thus, the odd coefficients just reproduce our previous solution, so we can set $a_1 = 0$ to produce a linearly independent second solution.

Thus, the linearly independent second solution is given entirely by the even powers (*cf.* Equation (8)):

$$J_{-\frac{1}{2}}(z) = \left(\frac{z}{2}\right)^{-1/2} \left(1 - \frac{1}{2!}z^2 + \frac{1}{4!}z^4 - \cdots\right)$$

where, following convention, we have taken $a_0 = 2^{1/2}$. Hence,

$$J_{-\frac{1}{2}}(z) = \sqrt{2} \frac{\cos z}{z^{1/2}}.$$

Note that our two solutions are just multiples of $z^{-1/2} \sin z$ and $z^{-1/2} \cos z$ (behaving as $z^{1/2}$ and $z^{-1/2}$ respectively as $z \to 0$).