

# Some remarks on spectral analysis of Markov chains: from Protein Dynamics to MCMC algorithms

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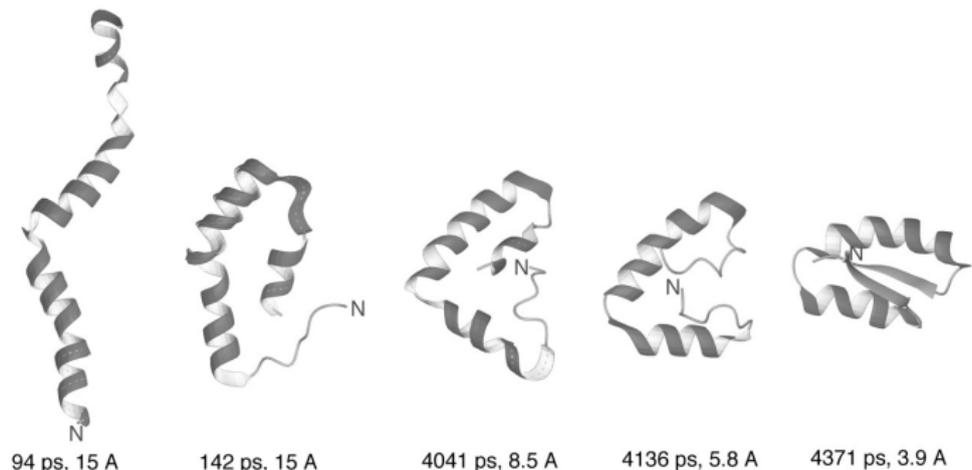
- 1 Introduction
- 2 Protein Dynamics and Approximations
- 3 Connection discrete time – continuous time Markov chain
- 4 Spectral Analysis of Markov operators

# Outline

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## Context

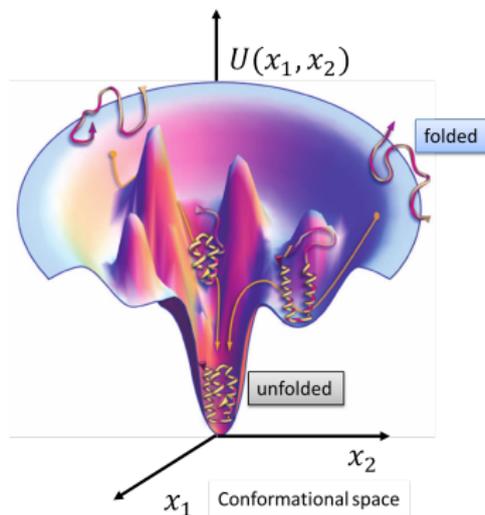
- Protein is a chain of amino acids
- Proteins are assumed to be found in a (possibly large) number of states  $X \in \mathcal{S}$
- A state encompasses information such as the distance/angle between amino acids, energy levels... (so  $|\mathcal{S}| \gg 1$ )



**Figure:** Protein folding pathway of 1E0G obtained in Langevin dynamics simulations (A. Liwo et al, PNAS, 2005)

## Inference of protein dynamics

A model for protein dynamics (typically a Markovian process  $\{X_t\}$  ( $t > 0$ )) can be fitted from experimental measurements



**Figure:** Model of a protein energy landscape (H. Ma et al, PNAS, 2006)

Equilibrium distribution of  $\{X_t\}$  is

$$\Pr(X_t = x) \propto \exp\{-H(x)\}$$

# Inference of protein dynamics

Interest lies in getting summaries from the model

- Is there a simpler model? *i.e.*  $X \in \mathcal{S}'$  and  $|\mathcal{S}'| \ll |\mathcal{S}|$   
 $\Rightarrow$  Is there an equivalent two state model  
 $\mathcal{S} = \{\text{Folded}\}, \{\text{Unfolded}\}$ ?
- If the protein is at an intermediate state what is the chance that it will first Fold before Unfold?

Spectral representation of the Markov operator ruling the dynamic of  $\{X_t\}$  reveals some answers to those questions.

# Markov chain Monte Carlo algorithms

Class of algorithms that simulate discrete time Markov chains  $\{X_k\}$  ( $k \in \mathbb{N}$ ) to perform numerical integration (Bayesian inference in particular)

Quantitatively, those algorithms are usually assessed/ranked according to:

- Speed of convergence of the law of the chain to the target distribution
- Asymptotic variance of Monte Carlo estimators

Here again, those quantities can be revealed by the spectral analysis of the Markov operator that simulates  $\{X_k\}$

# Purpose of the talk

- Gain some insight on how to diagnose when a two state approximation of the protein dynamics model is relevant
- Quantify this approximation
- Compare spectral analysis of continuous time Markov chain vs discrete time Markov chains in general state space
- How the approximation of the chain distribution established in CTMC can be interpreted in the context of MCMC algorithms

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## Specification of the model

- Even though protein dynamics are naturally parameterized by continuous parameters (angles, momenta, energy), it is assumed to be discretized *i.e.*

$$X_t \in \{1, 2, \dots, N\}$$

with  $N \gg 1$

- Each state  $i$  corresponds to a given range of angle, energy...
- The process  $\{X_t\}$  is assumed to be memoryless (Master equation)

$$\begin{aligned} \Delta\{\text{particles in state } i \text{ during } t, t+dt\} = \\ \# \{\text{new particles in state } i \text{ during in } t, t+dt\} \\ - \# \{\text{particles leaving state } i \text{ during in } t, t+dt\} \end{aligned}$$

and

$$\begin{aligned} \# \{\text{new particles in state } i \text{ during in } t, t+dt\} \\ = \sum_{k \neq i} \alpha_{k,i} \# \{\text{particles in state } k \text{ at time } t\} \end{aligned}$$

## Continuous Time Markov chain (CTMC)

The previous assumptions are equivalent to CTMC model.

Let  $\{X_t\}$  ( $t \geq 0$ ) be a stochastic process on the discrete state  $\mathcal{S} = \{1, \dots, N\}$  that evolves as follows:

$$\begin{aligned}\mathbb{P}(X_{t+dt} = j \mid X_{0:t}, X_t = i) = \\ \mathbb{P}(X_{t+dt} = j \mid X_t = i) = Q_{i,j}dt + o(dt), \quad i \neq j \quad (1)\end{aligned}$$

In this assumption  $\{X_t\}$  is a Continuous Time Markov chain (CTMC) and is characterized by:

- 1 an initial distribution  $\pi_0$  on  $(\mathcal{S}, \mathfrak{S})$
- 2 an infinitesimal generator  $Q \in \mathcal{M}_N(\mathbb{R})$  defined as

- $i \neq j$

$$Q_{i,j} = \frac{d}{dt} \lim_{t \downarrow 0} \mathbb{P}(X_t = j \mid X_0 = j)$$

- $i = j$

$$Q_{i,i} = - \sum_{i \neq j} Q_{i,j}$$

## Distribution of the CTMC

We denote by  $P(t) = \{\mathbb{P}(X_t = j | X_0 = i)\}_{i,j}$  the matrix of probability of transition of the CTMC. It is related to  $Q$  by:

$$P(t) = \exp\{Qt\}$$

(from Kolmogorov equations)

Denoting  $\pi(t) = \mathbb{P}(X_t \in \cdot) = \sum_{i=1}^N \mathbb{P}(X_t \in \cdot, X_0 = i)$ , we have:

$$\pi(t) = \pi(0) \exp\{Qt\}$$

### Assumption 1

We will assume that  $\{X_t\}$

- has an unique stationary distribution  $\pi$
- is time reversible:

$$\pi_i P_{i,j}(t) = \pi_j P_{j,i}(t), \quad \forall t > 0.$$

## Spectral decomposition of $Q$

Under Assumption 1, we have:

$$\pi P(t) = \pi, \quad P(t)\mathbf{1} = \mathbf{1}, \quad \text{sp}(P(t)) \subseteq (-1, 1).$$

Spectral properties of  $P(t)$  propagates to that of  $Q$ .

### Proposition 1

*sp(P(t)) and sp(Q) are connected:*

$$\lambda \in \text{sp}(Q) \Leftrightarrow \exp\{\lambda t\} \in \text{sp}(P(t))$$

Indeed, let  $y$  be a right eigenvector of  $Q$  with eigenvalue  $\lambda$ , then

$$P(t)y = \sum_{k=0}^N \frac{t^k}{k!} Q^k y = \sum_{k=0}^N \frac{(\lambda t)^k}{k!} y = \exp\{\lambda t\}y.$$

- $\mathbf{1} \in \text{sp}(P(t))$  implies  $0 \in \text{sp}(Q)$
- $\pi Q = 0$  so  $y_1^{(L)} = \pi$
- $Q\mathbf{1} = 0$  so  $y_1^{(R)} = \mathbf{1}$
- And in fact  $\text{sp}(Q) \subset (-\infty, 0)$

## Spectral decomposition of $Q$

Under Assumption 1,  $Q$  is diagonalizable:

$$Q = UDU^{-1}$$

where:

- $U = [y_1^{(R)} \cdots y_N^{(R)}]$  and  $y_1^{(R)}, \dots, y_N^{(R)}$  are  $Q$  right eigenvectors, with  $y_1^{(R)} = \mathbf{1}_n$
- $U^{-1} = [y_1^{(L)'} \cdots y_N^{(L)'}]'$  and  $y_1^{(L)}, \dots, y_N^{(L)}$  are  $Q$  left eigenvectors, with  $y_1^{(L)} = \pi$
- $D = \text{diag}(\lambda_1, \dots, \lambda_N)$  with  $\lambda_1 = 0$  and  $\lambda_i < 0$  for  $i > 1$

so that:

$$\begin{aligned}\pi(t) &= \sum_{\ell=1}^n \langle \pi_0, y_{\ell}^{(R)} \rangle \exp\{\lambda_{\ell} t\} y_{\ell}^{(L)} \\ &= \pi + \sum_{\ell=2}^n \langle \pi_0, y_{\ell}^{(R)} \rangle \exp\{\lambda_{\ell} t\} y_{\ell}^{(L)}\end{aligned}$$

## Interpretation of the chain distribution

First, note that  $\sum_{i=1}^N y_\ell^{(L)}(i) = 0$

$$\pi(t) = \pi + \sum_{\ell=2}^n \langle \pi_0, y_\ell^{(R)} \rangle \exp\{\lambda_\ell t\} y_\ell^{(L)}$$

- one stationary process ( $\pi$ )
- ( $N - 1$ ) transient processes  $\rho_\ell$

$$\rho_\ell(t) = \langle \pi_0, y_\ell^{(R)} \rangle \exp\{\lambda_\ell t\} y_\ell^{(L)}$$

that act as probability mass transfer.

If  $\pi_0 = \pi + \delta y_\ell^{(L)}$ , it would take  $\tau_\ell := -1/\lambda_\ell$  time to absorb the perturbation (relaxation time).

## Spectral gap and the two state approximation

In protein dynamics, a Physicist is typically interested in knowing if  $\{X_t\}_t$  could be represented by a two state system

$$\tilde{X}_t \in \{\text{Unfolded}, \text{Folded}\}$$

and what the transition rate between those two states look like.

### Rule 1 (Buchete and Hummer, 2008)

*If there is a "large enough" gap in the spectrum of  $Q$ , that is*

$$\gamma = |\lambda_2 - \lambda_3| / |\lambda_1 - \lambda_2| \approx 10$$

*then  $\{\tilde{X}_t\}$  is a "good" approximation of  $\{X_t\}$ .*

Remark:

$$\gamma = \left| \frac{\lambda_2}{\lambda_3} - 1 \right| = \left| \frac{\tau_2}{\tau_3} - 1 \right|.$$

## Interpretation of this rule

Considering the time  $t' = t/\tau_2$

$$\pi(t') = \pi + \beta_2 \exp\{-t'\} y_2^{(L)} + \beta_3 \exp\{-t'(\gamma + 1)\} y_3^{(L)} + \dots$$

with  $\beta_i = \langle \pi_0, y_i^{(R)} \rangle$ .

If  $\gamma \approx 10$  then the  $N - 2$  slowest processes might be neglected without much error

$$\pi(t') \approx \tilde{\pi}(t') = \pi + \beta_2 \exp\{-t'\} y_2^{(L)}$$

### Fact 1

*If  $\gamma \approx 10$ , then the only probability mass transfer is between the two basins of potential, namely {Unfolded} and {Folded}.*

## Simulating random rate matrices

To gain more insight on the 2 state approximations, we resorted to simulating random rate matrix:

- simulate a stationary probability  $\pi$
- simulate the lower triangular rate matrix that has the  $m$ -first diagonals non-zeros filled with

$$Q_{i,j} \sim \text{expo}(\lambda)$$

- fill the upper triangular to ensure reversibility  $Q_{j,i} = Q_{i,j}\pi_i/\pi_j$
- fill the diagonal so that  $Q_{i,i} = -\sum_{j \neq i} Q_{i,j}$

The parameter  $\lambda$  was tuned in order to match the slowest relaxation time usually observed in protein dynamics.

## Two state approximation

We compare on one example where  $\gamma = 7.8$  ( $N = 52$ ), the true distribution  $\pi$  and the two state approximation  $\tilde{\pi}$

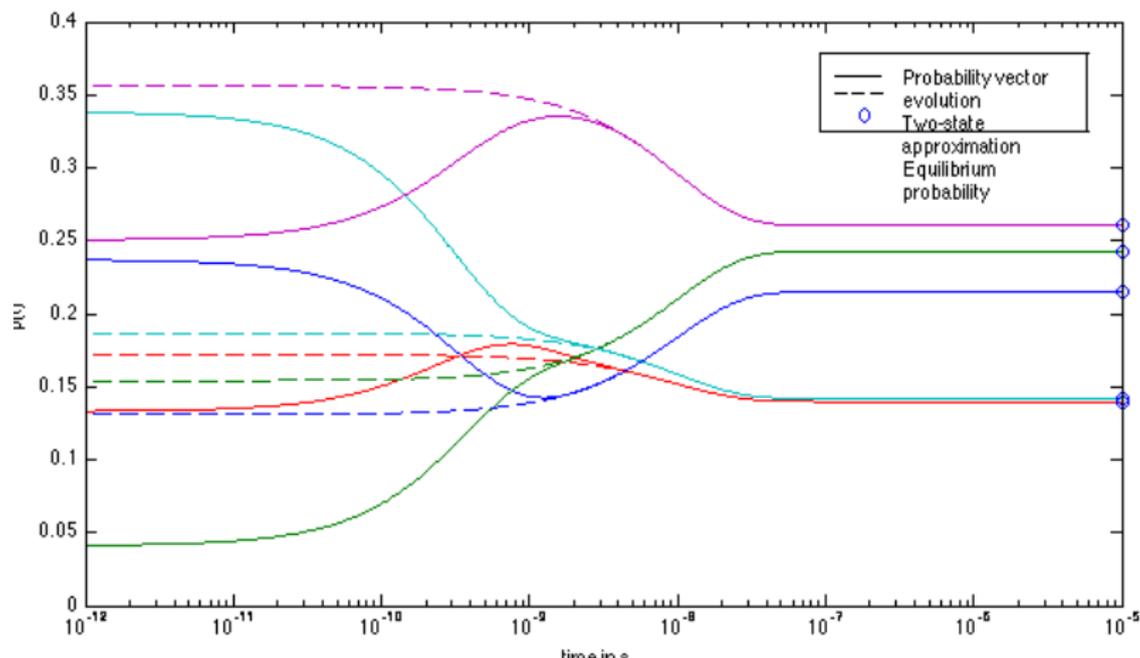


Figure: Representation of the probabilities  $\pi(t)$  and  $\tilde{\pi}(t)$  throughout time (in log scale) for four states.

## Correlation between $\gamma$ and error $L_2$

To quantify the approximation, we define

$$\text{Err} = \int \|\pi(t) - \tilde{\pi}(t)\|_2 dt$$

and try to estimate the correlation with the gap.

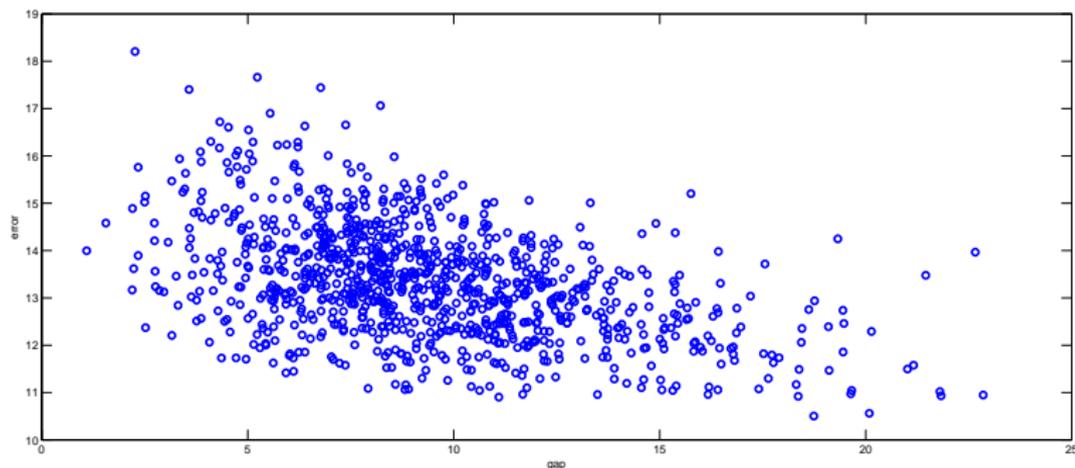
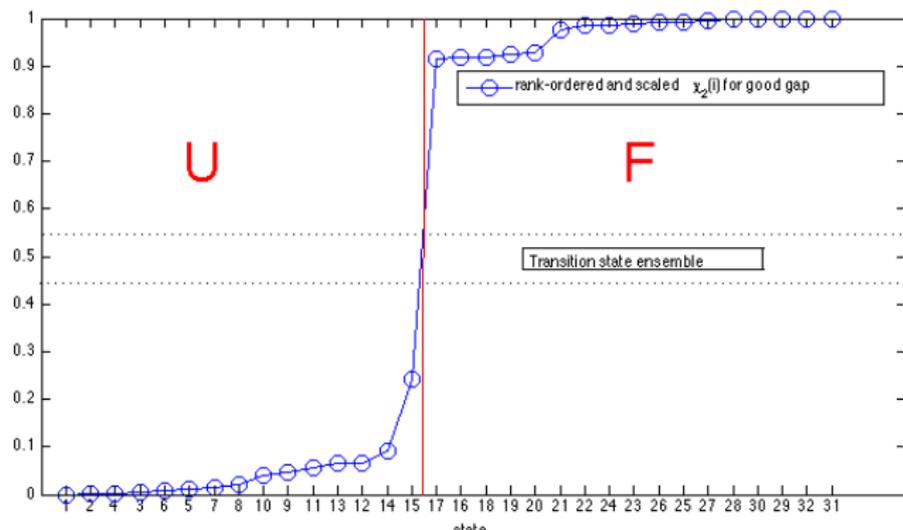


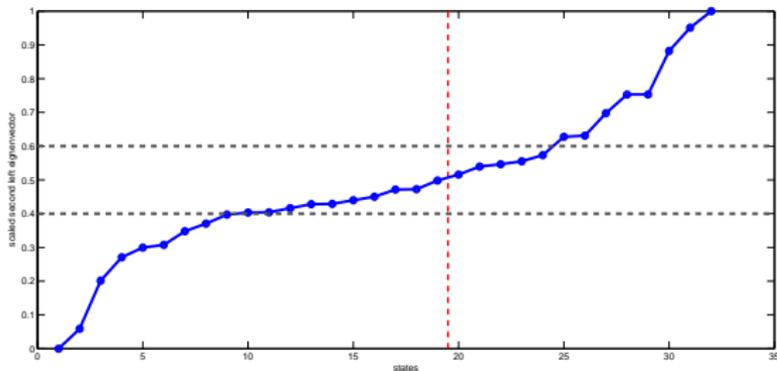
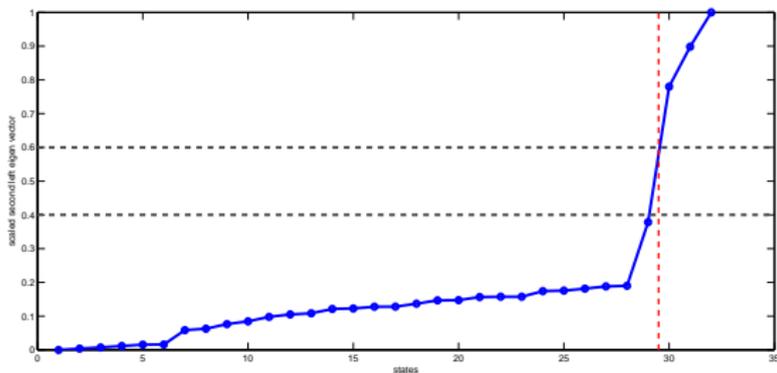
Figure: Realizations of gap plotted against Err.

# Illustration of the two states probability mass transfer channel

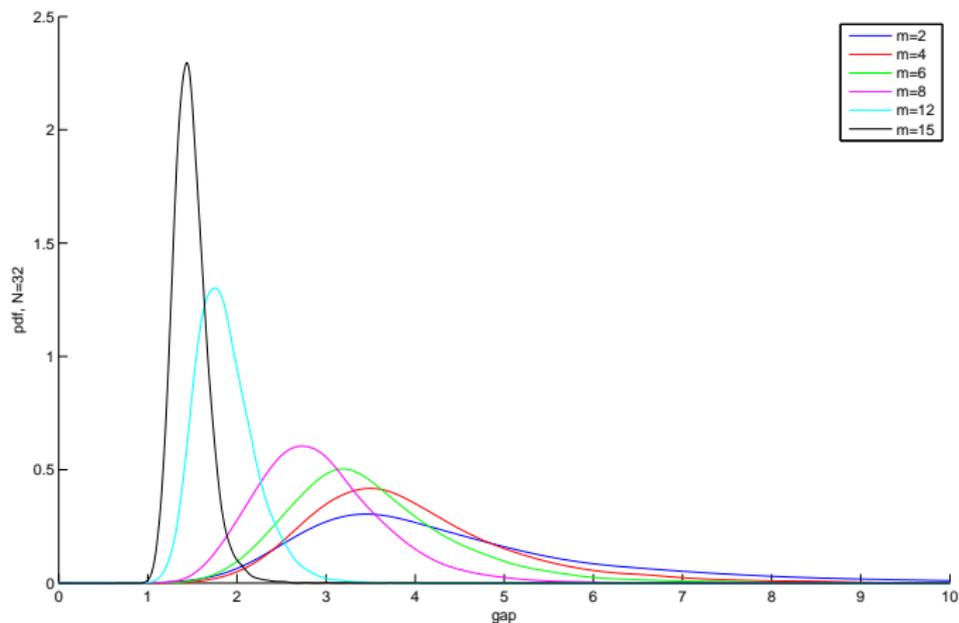
Representation of the scaled second left eigenvector of  $Q$ :  $y_2^{(L)}$  (in this example  $\gamma = 25.2$ ).



Representation of the scaled second left eigenvector of  $Q$ :  
 $y_2^{(L)}$  ( $\gamma = 8.2$  at the top and  $\gamma = 0.65$  at the bottom)

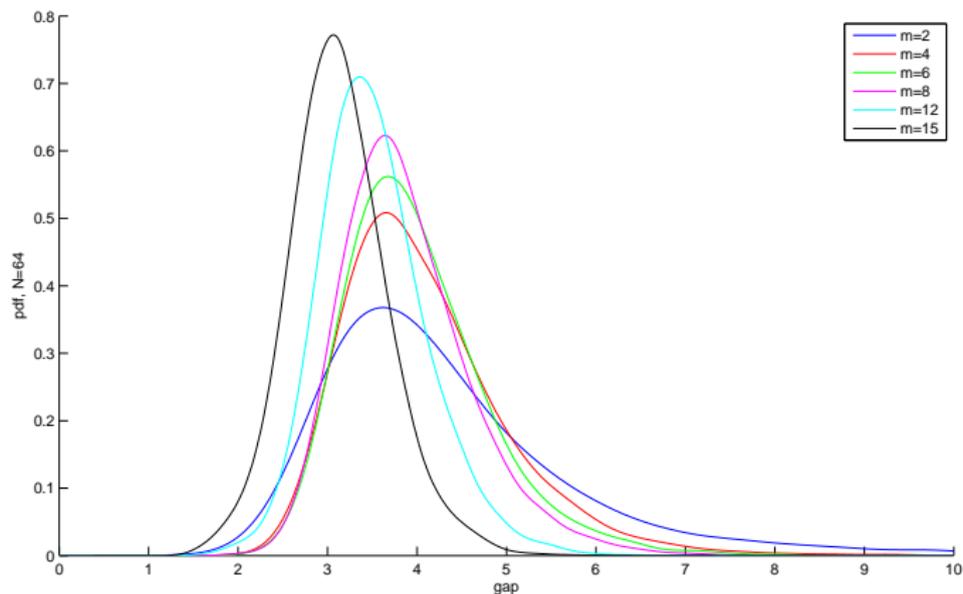


## Gap distribution ( $N = 32$ )



**Figure:** Gap probability density function for band diagonal random rate matrix of different structures.

## Gap distribution ( $N = 64$ )



**Figure:** Gap probability density function for band diagonal random rate matrix of different structures.

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## Discrete Time Markov chain (DTMC)

Let  $\{X_k\}$  ( $k \in \mathbb{N}$ ) be a stochastic process defined on the discrete state space  $\mathcal{S} = \{1, \dots, N\}$  that evolves as follows:

$$\text{for all } (i, j) \in \mathcal{S}^2, \quad \mathbb{P}(X_{k+1} = j \mid X_{0:k-1}, X_k = i) = \\ \mathbb{P}(X_{k+1} = j \mid X_k = i) = P_{i,j},$$

Under this assumption  $\{X_k\}$  is a Discrete Time Markov chain (DTMC) and is characterized by:

- 1 an initial distribution  $\pi_0$  on  $(\mathcal{S}, \mathfrak{G})$
- 2 a probability transition matrix  $P \in \mathcal{M}_N(\mathbb{R})$  satisfying:
  - for all  $(i, j) \in \mathcal{S}^2$

$$P_{i,j} \in (0, 1)$$

- for all  $i \in \mathcal{S}$

$$\sum_{j \in \mathcal{S}} P_{i,j} = 1$$

## A first observation

$$\begin{aligned}\mathbb{P}(X_{t+dt} = j | X_t = i) &= Q_{i,j}dt + o(dt), \quad i \neq j \\ &\updownarrow \\ \mathbb{P}(X_{t+dt} \neq i | X_t = i) &= \underbrace{\sum_{j \neq i} Q_{i,j} dt + o(dt)}_{\lambda_i} \quad (2)\end{aligned}$$

- same type of assumption than in a Poisson Process  
⇒ we know that the time until a change of state (*holding time*) is  $\tau \sim \text{expo}(\lambda_i)$  and

$$X_{t:t+\tau^-} = i$$

- next state satisfies

$$\begin{aligned}\mathbb{P}(X_{t+\tau} = j | X_{t+\tau^-} = i, X_{t+\tau} \neq i) \\ = \begin{cases} 0 & \text{if } i = j \\ \frac{\mathbb{P}(X_{t+\tau} = j | X_{t+\tau^-} = i)}{\mathbb{P}(X_{t+\tau} \neq i | X_{t+\tau^-} = i)} = Q_{i,j}/\lambda_i & \text{otherwise} \end{cases}\end{aligned}$$

# Algorithms to simulate a CTMC

From the previous observation, we deduce that Algorithm 1 simulates the CTMC of interest:

## Algorithm 1

- (1) draw an initial state  $Y_0 \sim \pi_0$  (say  $Y_0 \rightsquigarrow i$ )
- (2) draw a holding time (given  $Y_0 = i$ )  $\tau_0 \sim \text{expo}(\lambda_i)$
- (3) set  $X_{0:\tau_0} = i$  and draw a new state  $Y_1 \sim \bar{P}_{i,\cdot}$ , where

$$\bar{P}_{i,j} = \begin{cases} 0 & \text{if } i = j \\ Q_{i,j}/\lambda_i & \text{otherwise} \end{cases}$$

And then iterate (2)–(3)

# Algorithms to simulate a CTMC

Let  $\lambda^* = \max_i \lambda_i$  and consider the following modification of Algorithm 1:

## Algorithm 2

- (1) draw an initial state  $Y_0 \sim \pi_0$  (say  $Y_0 \rightsquigarrow i$ )
- (2) draw a holding time  $\tau_0 \sim \text{expo}(\lambda^*)$
- (3) set  $X_{0:\tau_0} = i$  and draw a new state  $Y_1 \sim \bar{R}_{i,\cdot}$ , where

$$\bar{R}_{i,j} = \begin{cases} 1 - \lambda_i/\lambda^* & \text{if } i = j \\ Q_{i,j}/\lambda^* & \text{otherwise} \end{cases}$$

And then iterate (2)–(3)

# CTMC-DTMC

Algorithm 2 can be decoupled:

- simulate a Poisson process  $\{N_t\}$  with parameters  $\lambda^*$
- simulate a DTMC  $\{Y_k\}$  with transition matrix  $\bar{R}$

## Proposition 2

Simulate  $\{Y_k\}$  and  $\{N_t\}$  as above. Define  $\{X_t\}$  as

$$\forall t \geq 0, \quad X_t = Y_{N_t}.$$

Then  $\{X_t\}$  is the desired CTMC.

	a counting process: $N_t$	a stochastic matrix: $\bar{R}_{i,j}$ ( $i \neq j$ )
DTMC	$\delta_{\mathbb{N}}$	$P_{i,j}$
CTMC	$PP(\lambda^*)$	$Q_{i,j}/\lambda^*$

## CTMC–DTMC: chain distribution

For the DTMC, it is straightforward to show that

$$\pi_n = \mathbb{P}(X_n \in \cdot) = \sum_{i=1}^N \mathbb{P}(X_n \in \cdot, X_0 = i) = \sum_{i=1}^N \pi_0(i) P^n(i, \cdot) = \pi_0 P^n$$

and for the CTMC, inspired by Algorithm 2, we write:

$$\pi(t) = \sum_{i=1}^N \sum_{n=0}^{\infty} \mathbb{P}(X_n \in \cdot, X_0 = i, N_t = n) = \sum_{n=0}^{\infty} \pi_0 \bar{R}^n \frac{(\lambda^* t)^n}{n!} \exp\{-\lambda^* t\}.$$

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## (Discrete Time) Markov chain on general state space

Let  $\{X_k\}$  ( $k \in \mathbb{N}$ ) be a stochastic process defined on the general state space  $(\mathbb{R}, \mathcal{B})$  that evolves as follows:

$$\text{for all } x_k \in \mathbb{R}, A \in \mathcal{B}, \quad \mathbb{P}(X_{k+1} \in A \mid X_{0:k-1}, X_k = x_k) = \\ \mathbb{P}(X_{k+1} \in A \mid X_k = x_k) = P(x_k, A).$$

Under this assumption  $\{X_k\}$  is a (Discrete Time) Markov chain on general state space and is characterized by:

- 1 an initial distribution  $\pi_0$  on  $(\mathbb{R}, \mathcal{B})$
- 2 a conditional probability distribution  $P(x, \cdot)$  on  $(\mathbb{R}, \mathcal{B})$ , determined by a function  $p : \mathbb{R} \times \mathbb{R} \rightarrow \mathbb{R}^+$ , satisfying:

- for all  $x \in \mathbb{R}$

$$P(x, \mathbb{R}) = \int_{\mathbb{R}} p(x, dy) = 1$$

- for all  $x \in \mathbb{R}, A \in \mathcal{B}$

$$P(x, A) = \int_A p(x, dy) \in (0, 1)$$

## Transition matrix as operator

For DTMC the operator is the transition matrix  $P$  and for CTMC the operator is  $P(t)$

- left operator: let  $\mu$  is a probability measure on  $\mathcal{S}$ ,

$$P : \mu \mapsto \mu P$$

is also a measure on  $\mathcal{S}$ , it is interpreted as

$$\mu \mapsto \mathbb{P}(X_1 \in \cdot \mid X_0 \sim \mu)$$

- right operation: let  $\omega$  be a vector of  $\mathbb{R}^n$ ,

$$P : \omega \mapsto P\omega$$

is interpreted as for all  $i \in \mathcal{S}$ :

$$\omega_i \mapsto \mathbb{E}\{\omega_{X_1} \mid X_0 = i\}$$

## Markov kernel as operator

For Markov chains on general state space, the operator is  $P(x, \cdot)$

- left operation: let  $\mu$  is a probability measure on  $(\mathbb{R}, \mathcal{B})$ ,

$$P : \mu \mapsto \mu P = \int_{\mathbb{R}} \mu(dx) P(x, \cdot)$$

is also a measure on  $(\mathbb{R}, \mathcal{B})$ , it is interpreted as

$$\mu \mapsto \mathbb{P}(X_1 \in \cdot \mid X_0 \sim \mu)$$

- right operation: let  $f : \mathbb{R} \rightarrow \mathbb{R}$  be a function,

$$P : f \mapsto Pf = \int_{\mathbb{R}} P(\cdot, dx) f(x)$$

is interpreted as:

$$f \mapsto \mathbb{E}\{f(X_1) \mid X_0 = \cdot\}$$

## Markov operator on general state space

Consider the (infinite-dimensional) space of functions  $\mathcal{L}^2(\pi)$  defined as:

$$\mathcal{L}^2(\pi) = \left\{ f : \mathbb{R} \rightarrow \mathbb{R}, \int \pi(dx) f(x)^2 < \infty \right\}$$

equipped with the scalar product

$$\langle f, g \rangle = \int \pi(dx) f(x)g(x).$$

### Assumption 2

*We assume that  $\{X_k\}$  is  $\pi$ -reversible:*

$$\int_A \pi(dx) P(x, B) = \int_B \pi(dx) P(x, A)$$

*(ie detailed balance condition holds).*

### Proposition 3

*Under Assumption 1, the operator  $P$  is a self adjoint operator on  $\mathcal{L}^2(\pi)$ .*

## Spectral analysis on $\mathcal{L}^2(\pi)$

### Theorem 1 (spectral theorem)

Let  $P$  be a compact and self-adjoint operator on an Hilbert space  $\mathcal{H}$ , then

- there exists an orthonormal basis of  $\mathcal{H}$  consisting of eigenvectors of  $P$
- the non-zero eigenvalues of  $P$  form a **finite** or **countably infinite** set  $\{\lambda_k\}$  such that

$$P = \sum_{\ell \geq 1} \lambda_\ell \Pi_\ell$$

where  $\Pi_\ell$  is the projection onto the eigenspace with eigenvalue  $\lambda_k$

So for any  $f \in \mathcal{L}^2(\pi)$ ,

$$Pf = \left( \sum_{\ell \geq 1} \lambda_\ell \Pi_\ell \right) f = \sum_{\ell > 1} \lambda_\ell \Pi_\ell f$$

## Case where the Markov kernel is compact

Let  $\{X_k\}$  be generated by the Metropolis-Hastings algorithm:

- 1 propose  $\tilde{X} \sim Q(X_k, \cdot)$
- 2 set  $X_{k+1} = \tilde{X}$  w.p.

$$1 \wedge \frac{\pi(\tilde{X})Q(\tilde{X}, X_k)}{\pi(X_k)Q(X_k, \tilde{X})}$$

and  $X_{k+1} = X_k$  otherwise.

**Proposition 4 (Atchadé and Perron, 2002)**

*If*

$$\iint Q(x, dy)Q(y, dx) < \infty$$

*then  $P$  (the operator induced by M-H) is compact.*

In particular this is always true for Independent M-H  
( $Q(x, dy) = Q(dy)$ )

## Chain distribution

To apply the Spectral Theorem to simplify

$$\mathbb{P}(X_n \in \cdot) = \pi_0 P^n(\cdot)$$

we need to map the left operator  $P$  to its right.

### Proposition 5

Let  $X_0 \sim \pi_0$  be the initial distribution of  $\{X_k\}$ . Assume  $\{X_k\}$  is  $\pi$ -reversible, ( $P$  is self adjoint), then

$$\pi_0 P^n(A) = \int_A \pi(dx) P^n f_0(x),$$

where  $f_0 = d\pi_0/d\pi$ .

### Corollary 1

Let  $\{X_k\}$  be a  $\pi$ -reversible Markov chain and  $P$  compact, then for any  $A \in \mathcal{B}(\mathbb{R})$

$$\pi_n(A) = \int_A \pi(dx) \sum_{\ell \geq 1} \lambda_\ell^n f_0^{(\ell)}(x), \quad f_0^{(\ell)} = \Pi_\ell f_0$$

## Chain distribution

Denote by  $\{\varepsilon_\ell\}_{\ell \geq 1}$  the eigenvectors of  $P$ .

Since  $\lambda_1 = 1 \in \text{sp}(P)$  and its eigenvector is the constant function,  $\varepsilon_1 = 1$ , we have:

$$\pi_n(A) = \int_A \pi(dx) \langle f_0, 1 \rangle 1(x) + \sum_{\ell \geq 2} \lambda_\ell^n \int_A \pi(dx) \langle f_0, \varepsilon_\ell \rangle \varepsilon_\ell(x),$$

but  $\langle f_0, 1 \rangle = \int \pi(dx) f_0(x) 1 = \int \pi_0(dx) = 1$  so that the first term is simply  $\pi_n(A)$ .

Comparing discrete and continuous contexts we have for  $i \in \{1, \dots, n\}$  and  $A \in \mathcal{B}(\mathbb{R})$ :

$$\pi_t(i) = \pi(i) + \sum_{\ell=2}^n \exp\{\lambda_\ell t\} \langle \pi_0, y_\ell^{(R)} \rangle y_\ell^{(L)}(i) \quad (3)$$

$$\pi_n(A) = \pi(A) + \sum_{\ell \geq 2} \lambda_\ell^n \langle f_0, \varepsilon_\ell \rangle \int_A \pi(dx) \varepsilon_\ell(x) \quad (4)$$

$\Rightarrow$  Under the self adjoint and compact assumption, a same interpretation of the probability mass transfer can be given to Markov chain on general state space.

## Case where the Markov kernel is not compact

The spectral analysis is much more complicated.

### Theorem 2 (Von Neumann's Spectral Theorem)

If  $P$  is self-adjoint and  $(f, g) \in \mathcal{L}^2(\pi)$ , then

$$\langle \phi(P)f, g \rangle = \int_{sp(P)} \phi(\lambda) d\mu_{f,g,P}(\lambda)$$

where  $d\mu_{f,g,P}$  is a measure defined as  $d\mu_{f,g,P}(\lambda) = \langle E_P(\lambda)f, g \rangle$  where  $E_P$  is the spectral measure (a projection valued measure) of  $P$ .

### Corrolary 2

If  $P$  is self-adjoint, then:

$$\pi_n(A) = \pi_0 P^n(A) = \int_{sp(P)} \lambda^n d \langle E(\lambda)f_0, \mathbb{1}_A \rangle$$

$\Rightarrow$  since the spectrum is continuous it is not clear how to exhibit a similar decomposition than in the compact case.

## Discussion

- spectral analysis of CTMC allows to derive reliable statistics giving insight on the process of interest:
  - a spectral gap between the second and third eigenvalue (a ratio of 10) seems to support a two state approximation (simulation)
  - but it is not true that if two models have two spectral gaps  $\gamma_1, \gamma_2$  such that  $\gamma_1 < \gamma_2$  then model 2 will yield a better two state approximation than model 1
- Interestingly, this analysis is very similar for Markov operator used to simulate Markov chains in MCMC algorithms
  - looks promising if the operator is compact
  - more work if we remove this assumption
- the decomposition of the MCMC distribution allows an interpretation of the chain dynamics in terms of flux of probability mass between high density regions