

# Extraction of very high precision post-Newtonian parameters from self-force computations

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# Introduction to Radiation gauge

- We start with the Teukolsky equation which has the form,

$$\mathcal{OT}(h) = \mathcal{SE}(h).$$

- From this we want to extract the perturbed metric,  $h$ .

# Teukolsky equation

Newman-Penrose equations (Bianchi identities):

Derivative operators acting on Weyl scalars

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$$\mathcal{OT}(h) = \mathcal{SE}(h)$$

# Teukolsky equation

$$\mathcal{O}[\mathcal{T}(h)] = \mathcal{S}\mathcal{E}(h)$$

perturbed Weyl tensor dotted with  
the background null tetrad, i.e.,  $\psi_0$

# Teukolsky equation

$$\mathcal{OT}(h) = \mathcal{S}\boxed{\mathcal{E}(h)}$$

Einstein operator acting on  $h = 8\pi T_{\mu\nu}$

# Teukolsky equation

$$\mathcal{OT}(h) = \boxed{\mathcal{S}}\mathcal{E}(h)$$

2<sup>nd</sup> order derivative operator acting on  $T_{\mu\nu}$

# Teukolsky equation

$$\mathcal{O}\mathcal{T}(h) = \mathcal{S}\mathcal{E}(h)$$

$\mathcal{O}$

2<sup>nd</sup> order derivative operator acting on  $\psi_0$

# Finding $h^{\text{ren}}$ from $\psi_0^{\text{ren}}$

**Chrzanowski-Cohen-Kegeles-Wald**

**Theorem:** Suppose  $\mathcal{SE} = \mathcal{OT}$  holds,  
and suppose  $\Psi$  satisfies  $\mathcal{O}^\dagger \Psi = 0$ .  
If  $\mathcal{E}$  is self-adjoint, then  $\mathcal{S}^\dagger \Psi$  satisfies  $\mathcal{E}(f) = 0$ .

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**Proof:** Taking adjoint of  $\mathcal{SE} = \mathcal{OT}$ , gives us

$$\mathcal{E}^\dagger \mathcal{S}^\dagger = \mathcal{T}^\dagger \mathcal{O}^\dagger$$

$$\mathcal{E} \mathcal{S}^\dagger = \mathcal{T}^\dagger \mathcal{O}^\dagger$$

If  $\mathcal{O}^\dagger \Psi = 0$ , then  $\mathcal{E}(\mathcal{S}^\dagger \Psi) = 0$ , i.e.,  $h = \mathcal{S}^\dagger \Psi$

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How do we connect the solution to the Teukolsky equation,  $\psi_0$  or  $\mathcal{T}(h)$ , to this  $\Psi$ ?

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$$\begin{aligned}\mathcal{SE}(\mathcal{S}^\dagger \Psi) &= \mathcal{OT}(\mathcal{S}^\dagger \Psi) \\ 0 &= \mathcal{O}[\mathcal{T}\mathcal{S}^\dagger \Psi]\end{aligned}$$

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$\mathcal{T} \mathcal{S}^\dagger$  maps solutions of  $\mathcal{O}^\dagger \Psi = 0$  to  $\mathcal{O} \psi = 0$ .

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$$\psi_0 = \mathcal{T}\mathcal{S}^\dagger \Psi$$

$$h = \mathcal{S}^\dagger \Psi$$

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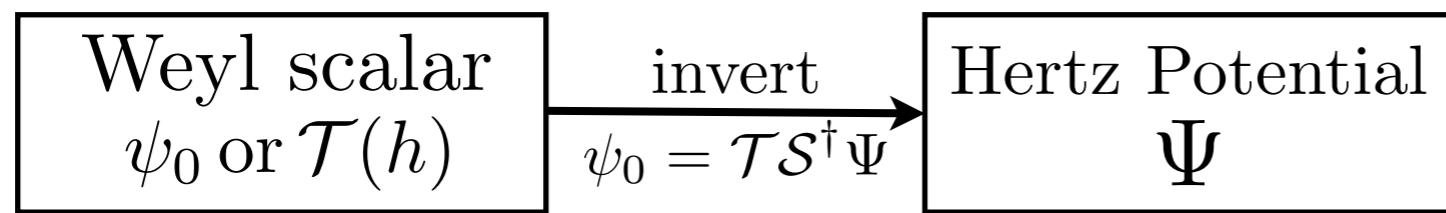
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# Summary

Weyl scalar  
 $\psi_0$  or  $\mathcal{T}(h)$

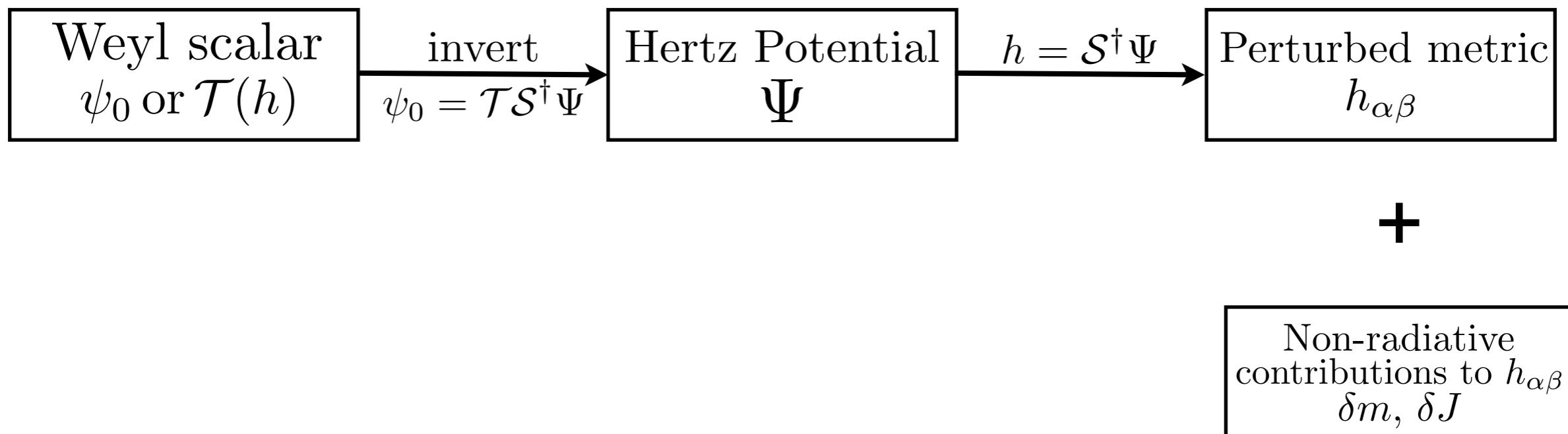
# Summary



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# Summary



+

$$u^\alpha u^\beta (g_{\alpha\beta} + h_{\alpha\beta}) = 1$$

$$u^\alpha = [u_0^t + u_1^t + O(\mu^2)] k^\alpha$$

Non-radiative  
contributions to  $h_{\alpha\beta}$   
 $\delta m, \delta J$

$$\Delta U = u_1^t = u_0^t H$$

$$H := \frac{1}{2} h_{\alpha\beta}^R u_0^\alpha u_0^\beta$$

||

$u_1^t$

# Blanchet et al

coeff.	value
$\alpha_4$	-114.34747(5)
$\alpha_5$	-245.53(1)
$\alpha_6$	-695(2)
$\beta_6$	+339.3(5)
$\alpha_7$	-5837(16)

# How its done?

MST (Mano-Suzuki-Takasugi) algorithm for the radial harmonics as a sum over known analytic functions with good convergence

Analytical form of the angular harmonics,  ${}_s Y_{\ell m}(\frac{\pi}{2}, 0)$

We use *Mathematica* which can handle very high precision computations

The accuracy is

about 1 part in  $10^{227}$  for  $r = 10^{20} M$   
about 1 part in  $10^{242}$  for  $r = 10^{25} M$   
about 1 part in  $10^{252}$  for  $r = 10^{30} M$

## (Expected) pN expansion of $u_1^t$

$$\begin{aligned} u_1^t = & \frac{\alpha_0}{r} + \frac{\alpha_1}{r^2} + \frac{\alpha_2}{r^3} + \frac{\alpha_3}{r^4} + \frac{\alpha_4}{r^5} + \frac{\beta_4 \log(r)}{r^5} \\ & + \frac{\alpha_5}{r^6} + \frac{\beta_5 \log(r)}{r^6} + \frac{\alpha_6}{r^7} + \frac{\beta_6 \log(r)}{r^7} + \dots \end{aligned}$$

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Notice the absence of 5.5-pN term

no  $\frac{\alpha_{5.5}}{r^{6.5}}$

# Analytically known pN coefficients (from literature)

$$u_1^t = \frac{-1}{r} + \frac{-2}{r^2} + \frac{-5}{r^3} + \frac{\left(\frac{-121}{3} + \frac{41}{32}\pi^2\right)}{r^4}$$
$$+ \frac{\frac{-592384 - 196608\gamma + 10155\pi^2 - 393216\log(2)}{7680}}{r^5}$$
$$+ \frac{\frac{64}{5}\log(r)}{r^5} + \frac{\frac{956}{105}\log(r)}{r^6} + \dots$$

Bini & Damour '13\*

\* Thanks to Alexandre Le Tiec

$u_1^t$  at  $r = 10^{30} M$

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Read off the first 4 coefficients to  
30 places of accuracy

$$u_1^t \text{ at } r = 10^{30}M$$

Lets look at

$$u_1^t - \left( \frac{-1}{r} + \frac{-2}{r^2} + \frac{-5}{r^3} + \frac{\left(\frac{-121}{3} + \frac{41}{32}\pi^2\right)}{r^4} + \frac{\frac{64}{5}\log(r)}{r^5} + \frac{\frac{956}{105}\log(r)}{r^6} \right)$$

$$\begin{aligned} &= -114.34895136757260295204000244483653876441286 \\ &\quad 5284407038869234848092925596369282766597634376 \\ &\quad 1937212552305416054218993704569382600204271482 \\ &\quad 5538690979057075189\dots \times 10^{-150} \end{aligned}$$

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$$528440703886923484809285596369282766597634376$$

$$193721255218204271482$$

$$55386908204271482$$

Again read off  $\alpha_4$  to 30 decimal places

$$+ \frac{-59238\gamma}{r^5} + \frac{7680}{r^5} + \frac{\pi^2 - 39321\log(2)}{r^5}$$

$$+ \frac{\frac{64}{5}\log(r)}{r^5} + \frac{\frac{956}{105}\log(r)}{r^6} + \dots$$

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in our numerical matching.

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5.5pN comes from  $\ell = 2$  multipole,

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7.5pN comes from  $\ell = 2, 3, 4$  multipole . . .

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7.5pN comes from  $\ell = 2, 3, 4$  multipole . . .

Specifically,

5.5pN comes from  $\ell = 2, m = \pm 2$  multipole,  
6.5pN comes from  $\ell = 2, m = \pm 1, \pm 2$  and  
 $\ell = 3, m = \pm 1, \pm 3$  multipoles.

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This was very puzzling!!

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A successful comparison with the UC-Dublin group (compared the source,  $l=2$  multipole, of  $5.5\text{pN}$  term):

23-24 digits of agreement at  $r = 10^3\text{M}$

43 digits of agreement at  $r = 10^6\text{M}$

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To further confirm its presence, we computed analytical value of the  $5.5\text{pN}$  coefficient using the self-force recipe in a radiation gauge, and found that the numerically extracted value agrees with it to **113** significant digits!!



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pN theory didn't expect n.5-pN terms.  
And we got those terms starting at 5.5 pN.

About a month ago, it was understood that the origin of these n.5pN terms come from the “tails-of-tails” terms in pN theory (Luc Blanchet).

# Revised pN-series

$$\begin{aligned} u_1^t = & \frac{\alpha_0}{r} + \frac{\alpha_1}{r^2} + \frac{\alpha_2}{r^3} + \frac{\alpha_3}{r^4} + \frac{\alpha_4}{r^5} + \frac{\beta_4 \log(r)}{r^5} \\ & + \frac{\alpha_5}{r^6} + \frac{\beta_5 \log(r)}{r^6} + \frac{\alpha_{5.5}}{r^{6.5}} + \frac{\alpha_6}{r^7} + \frac{\beta_6 \log(r)}{r^7} \\ & + \frac{\alpha_{6.5}}{r^{7.5}} + \frac{\alpha_7}{r^8} + \frac{\beta_7 \log(r)}{r^8} + \frac{\gamma_7 \log^2(r)}{r^8} + \dots \end{aligned}$$

# Analytical value of other terms possible from high precision number?

$$u_1^t = \frac{\alpha_0}{r} + \frac{\alpha_1}{r^2} + \frac{\alpha_2}{r^3} + \frac{\alpha_3}{r^4} + \frac{\alpha_4}{r^5} + \frac{\beta_4 \log(r)}{r^5} \\ + \frac{\alpha_5}{r^6} + \frac{\beta_5 \log(r)}{r^6} + \frac{\alpha_{5.5}}{r^{6.5}} + \frac{\alpha_6}{r^7} + \frac{\beta_6 \log(r)}{r^7} \\ + \frac{\alpha_{6.5}}{r^{7.5}} + \frac{\alpha_7}{r^8} + \frac{\beta_7 \log(r)}{r^8} + \frac{\gamma_7 \log^2(r)}{r^8} + \dots$$

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What about other terms?

If possible, it will save us a number of long, tedious, analytical calculations, and help extract further terms with higher accuracy.

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What about other terms?

**YES!! INDEED!!**

Lets look at the value of the 6-pN log term

$$\begin{aligned} u_1^t = & \frac{\alpha_0}{r} + \frac{\alpha_1}{r^2} + \frac{\alpha_2}{r^3} + \frac{\alpha_3}{r^4} + \frac{\alpha_4}{r^5} + \frac{\beta_4 \log(r)}{r^5} \\ & + \frac{\alpha_5}{r^6} + \frac{\beta_5 \log(r)}{r^6} + \frac{\alpha_{5.5}}{r^{6.5}} + \frac{\alpha_6}{r^7} + \boxed{\frac{\beta_6 \log(r)}{r^7}} \\ & + \frac{\alpha_{6.5}}{r^{7.5}} + \frac{\alpha_7}{r^8} + \frac{\beta_7 \log(r)}{r^8} + \frac{\gamma_7 \log^2(r)}{r^8} + \dots \end{aligned}$$

## Numerically extracted value

$$\beta_6 = -90.398589065255731922398589065255731922\\398589065255731922398589065255731922\\3985890652557319223985890485251879955\dots$$

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$$\beta_6 = -90.398589065255731922398589065255731922$$
$$398589065255731922398589065255731922$$
$$3985890652557319223985890485251879955\dots$$

More than 5 repetition cycles

$$\beta_6 = \frac{-51256}{567}$$

# Another example

$$\begin{aligned} u_1^t = & \frac{\alpha_0}{r} + \frac{\alpha_1}{r^2} + \frac{\alpha_2}{r^3} + \frac{\alpha_3}{r^4} + \frac{\alpha_4}{r^5} + \frac{\beta_4 \log(r)}{r^5} \\ & + \frac{\alpha_5}{r^6} + \frac{\beta_5 \log(r)}{r^6} + \frac{\alpha_{5.5}}{r^{6.5}} + \frac{\alpha_6}{r^7} + \frac{\beta_6 \log(r)}{r^7} \\ & + \frac{\alpha_{6.5}}{r^{7.5}} + \frac{\alpha_7}{r^8} + \frac{\beta_7 \log(r)}{r^8} + \boxed{\frac{\gamma_7 \log^2(r)}{r^8}} + \dots \end{aligned}$$

## Numerically extracted value

$$\gamma_7 = 52.17523809523809523809523809523809  
523809523809523809523809523809523809  
5238095237538043489164331 \dots$$

## Numerically extracted value

$$\gamma_7 = 52.17\boxed{5}23809523809523809523809523809  
523809523809523809523809523809523809  
5238095237538043489164331 \dots$$

## Numerically extracted value

$$\gamma_7 = 52.17523809523809523809523809523809$$
$$523809523809523809523809523809523809$$
$$5238095237538043489164331 \dots$$

More than 11 repetition cycles

$$\gamma_7 = \frac{27392}{525}$$

# 6.5-pN term

$$\alpha_{6.5} = 69.30909049662575956322060886698020553525276282  
80640692917511789688546478292427121038828291  
6578039346534682043068130080764 \dots$$

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$$\frac{\alpha_{6.5}}{\pi} = 22.06176870748299319727891156462585034013605442  
1768707482993197278911564625850340136054787080 \dots$$

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$$1768707482993197278911564625850340136054787080 \dots$$

Almost 2 repetition cycles here

$$\alpha_{6.5} = \frac{81077\pi}{3675}$$

One is not always lucky to have such repetition cycles...

$$\frac{\alpha_{7.5}}{\pi} = 176.4975875153652931430709208486986264764042541820319598 \\ 0973759881572859646496135482085019945242682178 \dots$$

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Multiplying it with appropriate powers of the first few prime numbers clear things up

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Multiplying it with appropriate powers of the first few prime numbers clears things up

Consistent with the accuracy with which we extracted this number

# Final list of analytically known pN terms

$$\begin{aligned} & \frac{-1}{r} + \frac{-2}{r^2} + \frac{-5}{r^3} + \frac{\frac{-121}{3} + \frac{41\pi^2}{32}}{r^4} \\ & + \frac{-592384 - 196608\gamma + 10155\pi^2 - 393216\log(2)}{7680 r^5} \\ & + \frac{64 \log(r)}{5 r^5} + \frac{-956 \log(r)}{105 r^6} + \frac{-13696\pi}{525 r^{6.5}} + \frac{-51256 \log(r)}{567 r^7} \\ & + \frac{81077\pi}{3675 r^{7.5}} + \frac{27392 \log^2(r)}{525 r^8} + \frac{82561159\pi}{467775 r^{8.5}} + \frac{-27016 \log^2(r)}{2205 r^9} \\ & + \frac{-11723776\pi \log(r)}{55125 r^{9.5}} + \frac{-4027582708 \log^2(r)}{9823275 r^{10}} + \frac{99186502\pi \log(r)}{1157625 r^{10.5}} \\ & + \frac{23447552 \log^3(r)}{165375 r^{11}} \end{aligned}$$

# Final list of analytically known pN terms

$$\frac{-1}{r} + \frac{-2}{r^2} + \frac{-5}{r^3} + \frac{\frac{-121}{3} + \frac{41\pi^2}{32}}{r^4}$$

Bini & Damour '13\*

$$\begin{aligned}
 &+ \frac{-592384 - 196608\gamma + 10155\pi^2 - 393216\log(2)}{7680 r^5} \\
 &+ \frac{64 \log(r)}{5 r^5} + \frac{-956 \log(r)}{105 r^6} + \frac{-13696\pi}{525 r^{6.5}} + \frac{-51256 \log(r)}{567 r^7} \\
 &+ \frac{81077\pi}{3675 r^{7.5}} + \frac{27392 \log^2(r)}{525 r^8} + \frac{82561159\pi}{467775 r^{8.5}} + \frac{-27016 \log^2(r)}{2205 r^9} \\
 &+ \frac{-11723776\pi \log(r)}{55125 r^{9.5}} + \frac{-4027582708 \log^2(r)}{9823275 r^{10}} + \frac{99186502\pi \log(r)}{1157625 r^{10.5}} \\
 &+ \frac{23447552 \log^3(r)}{165375 r^{11}}
 \end{aligned}$$

\* Thanks to Alexandre Le Tiec

# Final list of analytically known pN terms

$$\begin{aligned}
& \frac{-1}{r} + \frac{-2}{r^2} + \frac{-5}{r^3} + \frac{\frac{-121}{3} + \frac{41\pi^2}{32}}{r^4} \\
& + \frac{-592384 - 196608\gamma + 10155\pi^2 - 393216\log(2)}{7680 r^5} \\
& + \frac{64 \log(r)}{5 r^5} + \frac{-956 \log(r)}{105 r^6} + \boxed{\frac{-13696\pi}{525 r^{6.5}} + \frac{-51256 \log(r)}{567 r^7}} \\
& + \frac{81077\pi}{3675 r^{7.5}} + \frac{27392 \log^2(r)}{525 r^8} + \frac{82561159\pi}{467775 r^{8.5}} + \frac{-27016 \log^2(r)}{2205 r^9} \\
& + \frac{-11723776\pi \log(r)}{55125 r^{9.5}} + \frac{-4027582708 \log^2(r)}{9823275 r^{10}} + \frac{99186502\pi \log(r)}{1157625 r^{10.5}} \\
& + \frac{23447552 \log^3(r)}{165375 r^{11}}
\end{aligned}$$

# Flux at future horizon

Recent work by Fujita calculated the pN series of flux at future null infinity to 22pN order!

Tagoshi et al calculated the pN series of flux at future horizon to 6pN order.

Using

$$\text{Flux lost} = \text{Flux at infinity} + \text{Flux at horizon}$$

we extract the 7,8,9,10-pN orders of flux at horizon

# Flux lost = Flux at infinity + Flux at horizon

pN order	Lost by particle	$\infty$	$r_+$
1	✓	✓	0
2	✓	✓	0
3	✓	✓	0
4	✓	✓	✓
5	✓	✓	✓
6	✓	✓	✓
7	✓	✓	?
8	✓	✓	?
9	✓	✓	?
10	✓	✓	?

Green tick-marks show what we know from literature and that our result agrees with previous work  
 Red question marks show what is unknown and how we extract the unknown coefficients.

# Flux lost = Flux at infinity + Flux at horizon

pN order	Lost by particle	$\infty$	$r_+$
1	✓	✓	0
2	✓	✓	0
3	✓	✓	0
4	✓	✓	✓
5	✓	✓	✓
6	✓	✓	✓
7	✓	✓	✓
8	✓	✓	✓
9	✓	✓	✓
10	✓	✓	✓

Green tick-marks show what we know from literature and that our result agrees with previous work  
 Red question marks show what is unknown and how we extract the unknown coefficients.

# Matching in Kerr

