

The Spectral Method (MAPH 40260)

Part 1: Spectral Analysis

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Outline

Introduction

Fourier Analysis

Vibrating String



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The Gridpoint Method

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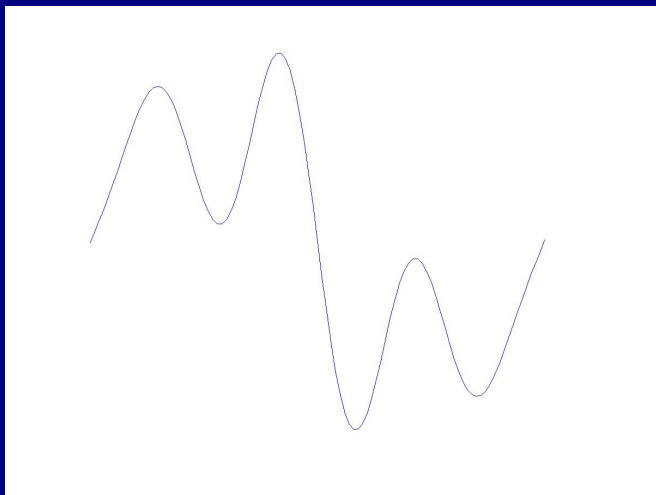
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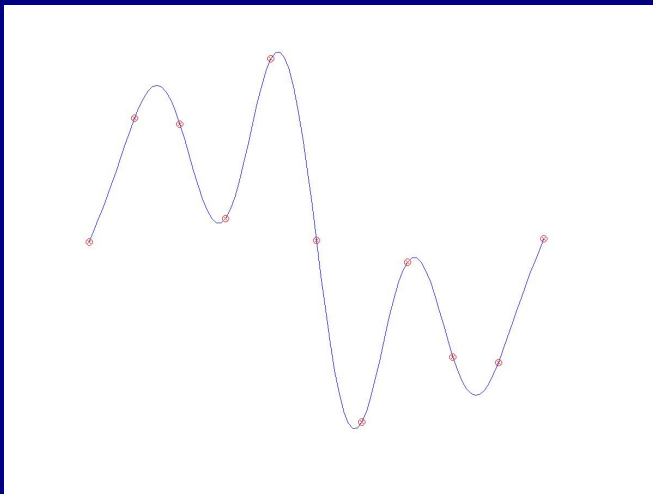
There are several answers to this question.



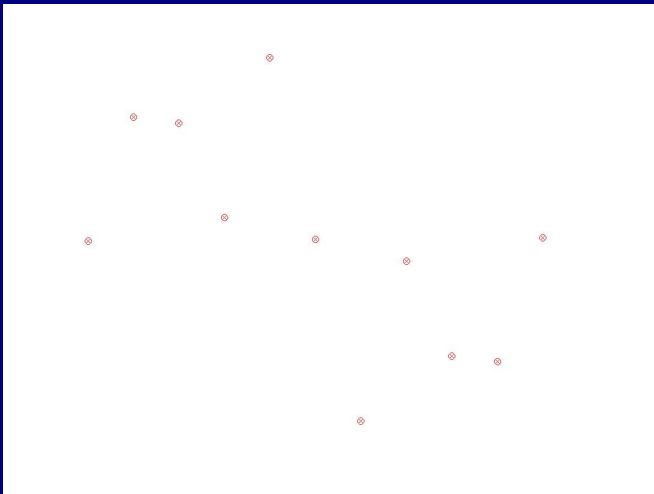


Continuous function of position.





Evaluation on a set of grid points.



Grid point values.



Spectral Analysis

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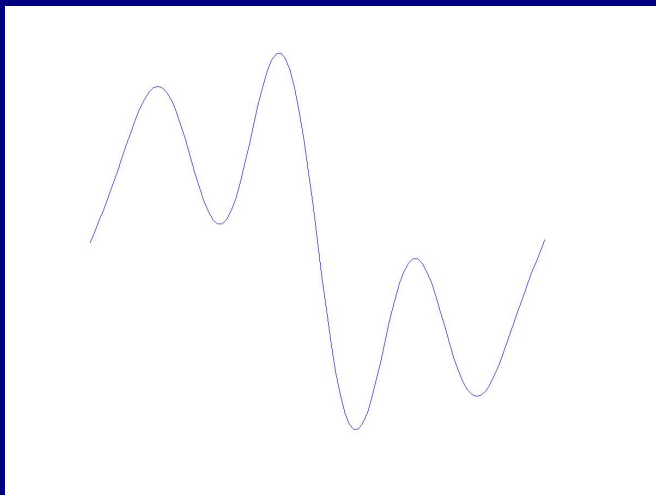
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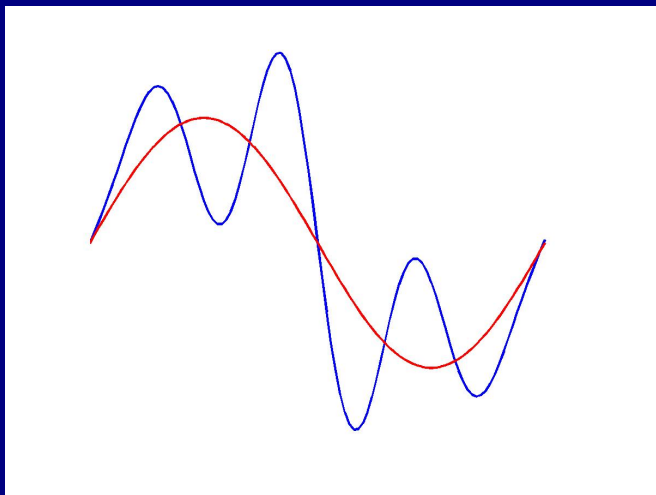
It is somewhat like splitting sunlight into the various colours of the spectrum.





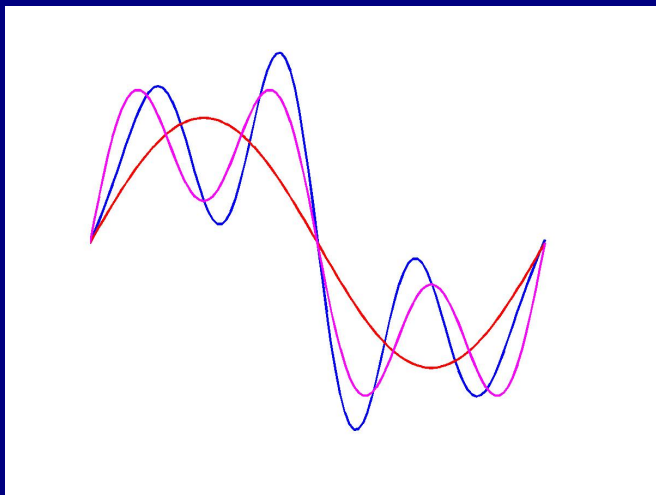
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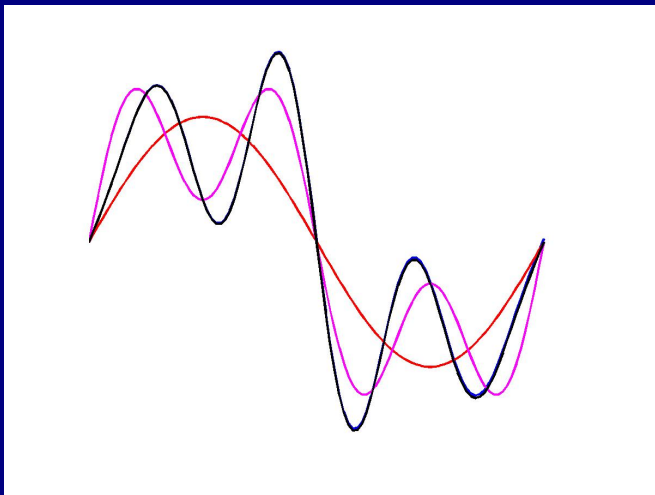
The first spectral component.





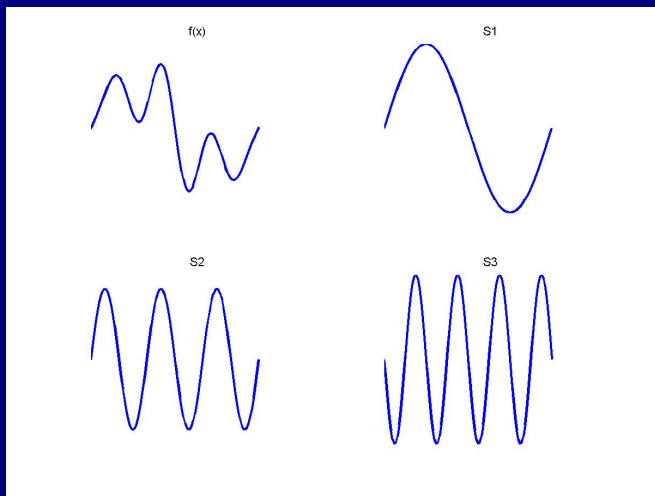
The first and second spectral components.





The first, second and third spectral components.





The original function and its three components.



Gridpoint Method



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- ▶ **Discrete representation**



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- ▶ **Values at geographical locations**



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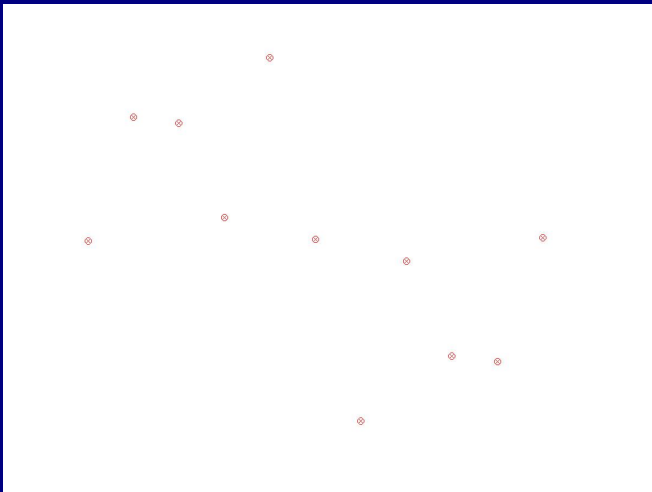


Gridpoint Method

- ▶ Discrete representation
- ▶ Values at geographical locations
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- ▶ Easy to represent graphically.
- ▶ Derivatives evaluated by **finite differences**.

Big drawback: Evaluation of derivatives involves errors.





Grid point values. We have to get derivatives from this set of values.



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- ▶ **Derivatives evaluated **exactly**, by analysis.**

Big advantage: Evaluation of derivatives is exact.



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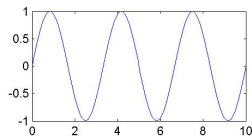
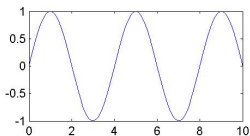
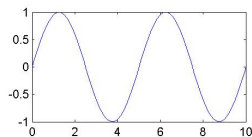
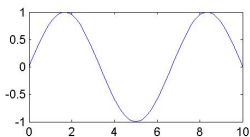
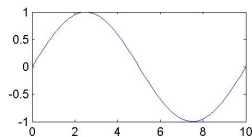
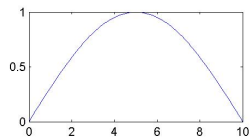
$$f(0) = f(\ell) = 0$$

We note that sinusoidal functions with certain wavelengths also vanish at $x = 0$ and at $x = \ell$:

$$\sin \pi x / \ell, \sin 2\pi x / \ell, \dots, \sin n\pi x / \ell,$$

for all integer values of n .





The first six harmonic components ($l = 10$).



Orthogonality

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We easily show that

$$\begin{aligned}\int_0^\ell [\Psi_n(x)]^2 dx &= \int_0^\ell \sin^2\left(\frac{n\pi}{\ell}x\right) dx \\ &= \frac{1}{2} \int_0^\ell \left[1 - \cos\left(\frac{2n\pi}{\ell}x\right)\right] dx \\ &= \left[\frac{x}{2}\right]_0^\ell - \left[\frac{\ell}{4\pi n} \sin\left(\frac{2n\pi}{\ell}x\right)\right]_0^\ell \\ &= \frac{\ell}{2}.\end{aligned}$$



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$$\begin{aligned} & \int_0^\ell \psi_n(x) \cdot \psi_m(x) dx \\ &= \int_0^\ell \sin\left(\frac{n\pi}{\ell}x\right) \cdot \sin\left(\frac{m\pi}{\ell}x\right) dx \\ &= \frac{1}{2} \int_0^\ell \left[\cos\left(\frac{n-m}{\ell}\pi x\right) - \cos\left(\frac{n+m}{\ell}\pi x\right) \right] dx \\ &= \frac{1}{2\pi} \left[\frac{\ell}{n-m} \sin\left(\frac{n-m}{\ell}\pi x\right) - \frac{\ell}{n+m} \sin\left(\frac{n+m}{\ell}\pi x\right) \right]_0^\ell \\ &= 0. \end{aligned}$$



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We thus have:

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We obtain an **orthonormal** set of functions:

$$\int_0^{\ell} \tilde{\Psi}_n(x) \cdot \tilde{\Psi}_m(x) dx = \delta_{mn}.$$



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We might also choose $\Phi(x, t) = \Psi(x) \cos(\omega t)$ or $\Phi(x, t) = \Psi(x) \sin(\omega t)$.



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k is the inverse of the wavelength, with a 2π factor.



The function

$$\Psi(x) = A \sin kx$$

is a solution of the o.d.e., and satisfies the boundary conditions if

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$\Psi_n(x)$ is an **eigenfunction** of the o.d.e. with **eigenvalue** $k_n = n\pi/\ell$.



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$$\Psi = \sum_{n=1}^{\infty} A_n \Psi_n(x).$$



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Thus,

$$A_m = \frac{2}{\ell} \int_0^{\ell} \Psi_m(x) \Psi(x) dx$$



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The function ψ can be obtained from the expansion

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if the coefficients A_n are known.



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There is a **duality** between $\psi(x)$, a function in **physical space** and $\{A_n\}$, the coefficients in **wavenumber space**. Given either representation, we can obtain the other.



Example: Analysis of a Square Wave

Let

$$\Psi(x) = +1, \text{ for } x \in [0, \frac{\ell}{2}] \quad \Psi(x) = -1, \text{ for } x \in [\frac{\ell}{2}, \ell].$$



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The Fourier coefficients are easily calculated.

$$A_n = \frac{2}{\ell} \int_0^{\ell} \Psi \cdot \Psi_n dx = \frac{2}{\ell} \left(\int_0^{\ell/2} \sin \frac{n\pi}{\ell} x dx - \int_{\ell/2}^{\ell} \sin \frac{n\pi}{\ell} x dx \right).$$



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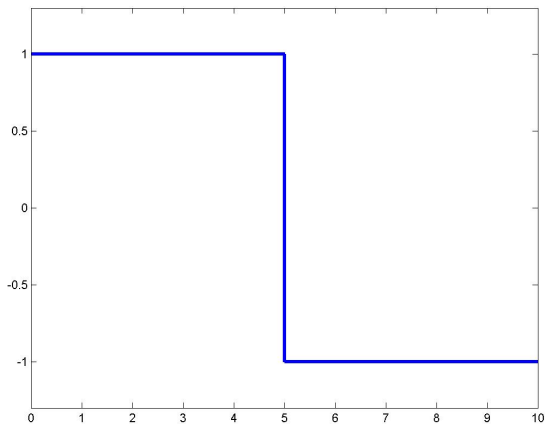
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When we work these out, we find that only every fourth coefficient has a nonzero value:

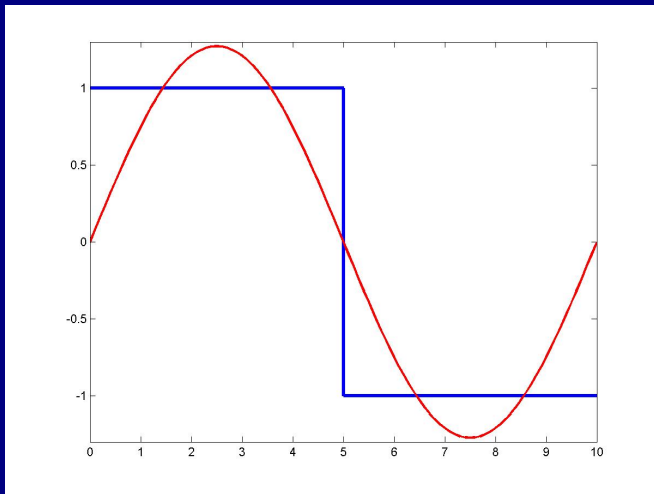
$$A_2 = \frac{1}{2} \left(\frac{8}{\pi} \right), \quad A_6 = \frac{1}{6} \left(\frac{8}{\pi} \right), \quad A_{10} = \frac{1}{10} \left(\frac{8}{\pi} \right), \quad \text{etc.}$$





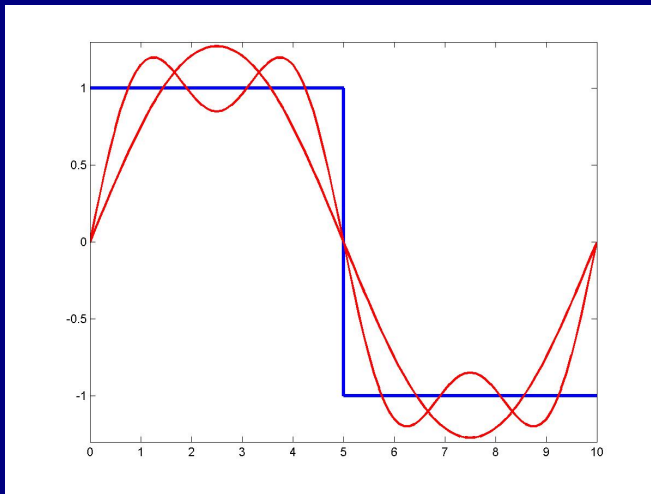
The square wave function ($l = 10$).





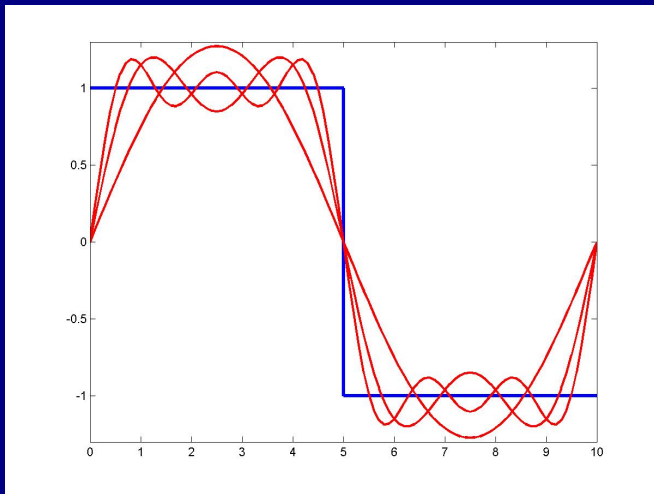
The square wave function. First coefficient.





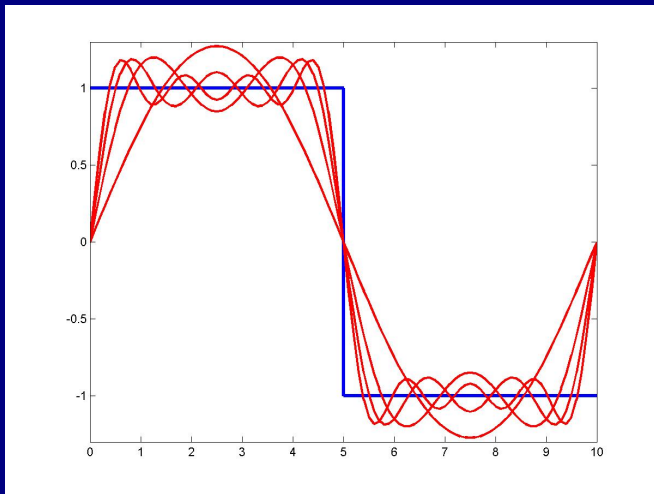
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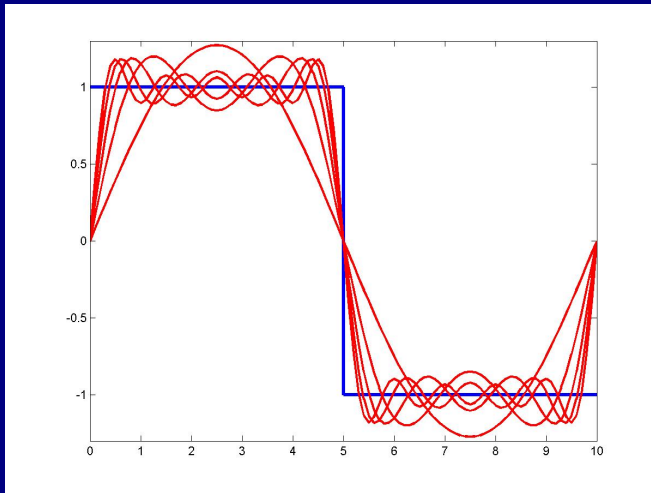
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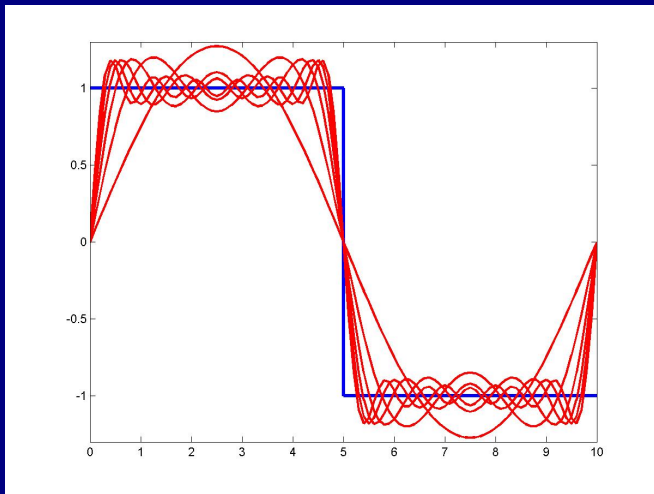
The square wave function. First four coefficients.





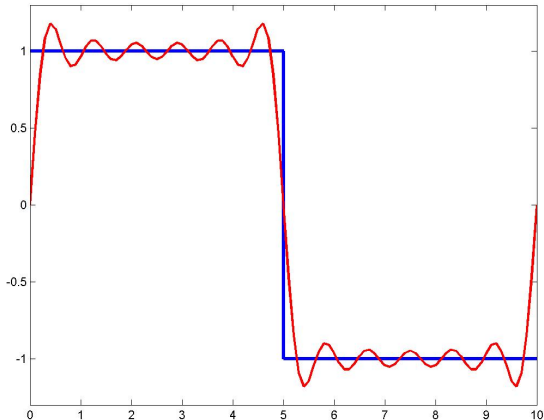
The square wave function. First five coefficients.





The square wave function. First six coefficients.

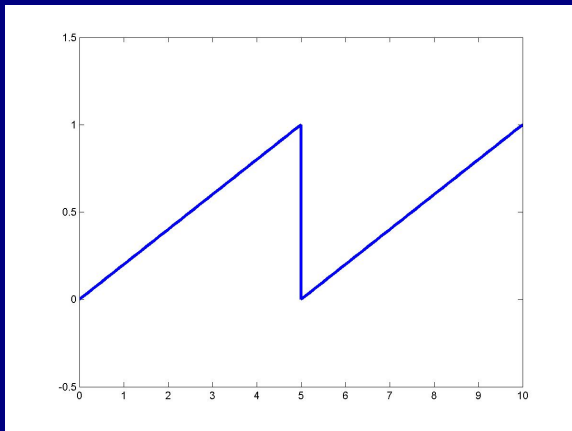




First six coefficients. Note Gibbs Phenomenon.



Exercise: Analysis of a Sawtooth Wave



Find the Fourier coefficients of the sawtooth function.



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At the initial time,

$$\Phi(x, t) = \Phi_0(x) = \sum_{n=1}^{\infty} A_n \Psi_n(x).$$

Also, because of the chosen form of solution,

$$\Phi_t(x, 0) = 0.$$



Again,

$$\Phi(x, t) = \Phi_0(x) = \sum_{n=1}^{\infty} A_n \Psi_n(x).$$



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This gives us the values of the coefficients:

$$A_n = \frac{2}{\ell} \int_0^{\ell} \Psi_n(x) \Phi_0(x) dx$$



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The problem is now completely solved:

$$\Phi(x, t) = \sum_{n=1}^{\infty} A_n \Psi_n(x) \cos \omega_n t,$$

with coefficients that are now known.



Again,

$$\Phi(x, t) = \Phi_0(x) = \sum_{n=1}^{\infty} A_n \Psi_n(x).$$

This gives us the values of the coefficients:

$$A_n = \frac{2}{\ell} \int_0^{\ell} \Psi_n(x) \Phi_0(x) dx$$

The problem is now completely solved:

$$\Phi(x, t) = \sum_{n=1}^{\infty} A_n \Psi_n(x) \cos \omega_n t,$$

with coefficients that are now known.

The eigenfunctions and eigenvalues are defined by

$$\Psi_n(x) \equiv \sin k_n x, \quad k_n = n\pi/\ell, \quad \omega_n = c k_n.$$



End of Part 1

