The material in this section is based largely on Lectures on Dynamical Meteorology by Roger Smith.

A *front* is the sloping interfacial region of air between two air masses, each of more or less uniform properties. [Fronts also occur also in the ocean, but we will not discuss them.]

A *front* is the sloping interfacial region of air between two air masses, each of more or less uniform properties. [Fronts also occur also in the ocean, but we will not discuss them.]

The primary example is the *polar front*, a zone of relatively large horizontal temperature gradient in the mid-latitudes that separates air masses of more uniform temperatures that lie polewards and equatorwards of the zone.

A *front* is the sloping interfacial region of air between two air masses, each of more or less uniform properties. [Fronts also occur also in the ocean, but we will not discuss them.]

The primary example is the *polar front*, a zone of relatively large horizontal temperature gradient in the mid-latitudes that separates air masses of more uniform temperatures that lie polewards and equatorwards of the zone.

This is associated with the midlatitude westerlies, having their maximum at the jetstream in the upper troposphere. This is the *Polar Front Jet*.

A *front* is the sloping interfacial region of air between two air masses, each of more or less uniform properties. [Fronts also occur also in the ocean, but we will not discuss them.]

The primary example is the *polar front*, a zone of relatively large horizontal temperature gradient in the mid-latitudes that separates air masses of more uniform temperatures that lie polewards and equatorwards of the zone.

This is associated with the midlatitude westerlies, having their maximum at the jetstream in the upper troposphere. This is the *Polar Front Jet*.

Regionally, where the polar front is particularly pronounced, we have *cold* and *warm fronts* associated with extra-tropical cyclones.

A *front* is the sloping interfacial region of air between two air masses, each of more or less uniform properties. [Fronts also occur also in the ocean, but we will not discuss them.]

The primary example is the *polar front*, a zone of relatively large horizontal temperature gradient in the mid-latitudes that separates air masses of more uniform temperatures that lie polewards and equatorwards of the zone.

This is associated with the midlatitude westerlies, having their maximum at the jetstream in the upper troposphere. This is the *Polar Front Jet*.

Regionally, where the polar front is particularly pronounced, we have *cold* and *warm fronts* associated with extra-tropical cyclones.

Sharp temperature differences can occur across a frontal surface: several degrees over a few kilometres.

The following Figure shows the passage of a cold front.





Composite meridional cross-section at 80°W of mean temperature and the zonal component of geostrophic wind computed from 12 individual cross-sections. The means were computed with respect to the position of the polar front in individual cases (from Palmén and Newton, 1948).



Analysed surface pressure, storm in October, 2000.

Margules' Model

Margules' Model



Max Margules (1856–1920)

The simplest model representing a *frontal discontinuity* is Margules' model.

In this model, the front is idealized as a sharp, plane, temperature discontinuity separating two inviscid, homogeneous, geostrophic flows.



Configuration of Margules' frontal model. Subscripts 1 and 2 refer to the warm and cold air masses.

Further, we assume:

1. The Boussinesq approximation: neglect variations in density except where they are coupled with gravity.

Further, we assume:

- 1. The Boussinesq approximation: neglect variations in density except where they are coupled with gravity.
- 2. The flow is everywhere parallel to the front and there are no along-front variations in it; i.e., $\partial v/\partial y = 0$.

Further, we assume:

- 1. The Boussinesq approximation: neglect variations in density except where they are coupled with gravity.
- 2. The flow is everywhere parallel to the front and there are no along-front variations in it; i.e., $\partial v/\partial y = 0$.
- 3. Diffusion effects are absent so that the frontal discontinuity remains sharp.

Further, we assume:

- 1. The Boussinesq approximation: neglect variations in density except where they are coupled with gravity.
- 2. The flow is everywhere parallel to the front and there are no along-front variations in it; i.e., $\partial v/\partial y = 0$.
- 3. Diffusion effects are absent so that the frontal discontinuity remains sharp.

We assume that the temperature difference between the air masses is small in the sense that $(T_1 - T_2)/\overline{T} \ll 1$, where $\overline{T} = (T_1+T_2)/2$ is the mean temperature of the two air masses, T_1 the temperature of the warm air and T_2 the temperature of the cold air.

The equations of motion are then:

The geostrophic equations:

$$u = 0, \qquad fv = \frac{1}{\bar{\rho}}\frac{\partial p}{\partial x}$$

The equations of motion are then: The geostrophic equations:

$$u = 0, \qquad fv = \frac{1}{\bar{\rho}}\frac{\partial p}{\partial x}$$

The hydrostatic equation:

$$\frac{1}{\bar{\rho}}\frac{\partial p}{\partial z} = -\frac{g(T-T')}{\bar{T}}$$

The equations of motion are then:

The geostrophic equations:

$$u = 0, \qquad fv = \frac{1}{\bar{\rho}}\frac{\partial p}{\partial x}$$

The hydrostatic equation:

$$\frac{1}{\bar{\rho}}\frac{\partial p}{\partial z} = -\frac{g(\bar{T} - T')}{\bar{T}}$$

The continuity equation:

$$\frac{\partial u}{\partial x} + \frac{\partial w}{\partial z} = 0$$

In Margules' model the vertical temperature gradient in each air mass is assumed to be zero.

In Margules' model the vertical temperature gradient in each air mass is assumed to be zero.

In fact, the temperature in each airmass is constant, varying neither in the horizontal nor in the vertical direction.

In Margules' model the vertical temperature gradient in each air mass is assumed to be zero.

In fact, the temperature in each airmass is constant, varying neither in the horizontal nor in the vertical direction.

We consider this to be the limiting case of the situation in which the temperature gradients are very small except across the frontal zone, where they are very large.



Frontal zone

Vertical cross-section through a (smeared-out) front. The coloured lines indicate isotherms.

In the frontal zone T = T(x, z). Otherwise, T is constant.

$$\delta T = 0 = \frac{\partial T}{\partial x} \delta x + \frac{\partial T}{\partial z} \delta z$$

$$\delta T = 0 = \frac{\partial T}{\partial x} \delta x + \frac{\partial T}{\partial z} \delta z$$

Therefore, the local slope $|\delta z/\delta x|$ of an isotherm in the frontal zone is given by

$$\tan \varepsilon = -\frac{\delta z}{\delta x} = \frac{\partial T/\partial x}{\partial T/\partial z}$$

$$\delta T = 0 = \frac{\partial T}{\partial x} \delta x + \frac{\partial T}{\partial z} \delta z$$

Therefore, the local slope $|\delta z/\delta x|$ of an isotherm in the frontal zone is given by

$$\tan \varepsilon = -\frac{\delta z}{\delta x} = \frac{\partial T/\partial x}{\partial T/\partial z}$$

Note that $\delta x > 0$ implies $\delta z < 0$ if, as assumed, $0 < \varepsilon < \pi/2$.

$$\delta T = 0 = \frac{\partial T}{\partial x} \delta x + \frac{\partial T}{\partial z} \delta z$$

Therefore, the local slope $|\delta z/\delta x|$ of an isotherm in the frontal zone is given by

$$\tan \varepsilon = -\frac{\delta z}{\delta x} = \frac{\partial T/\partial x}{\partial T/\partial z}$$

Note that $\delta x > 0$ implies $\delta z < 0$ if, as assumed, $0 < \varepsilon < \pi/2$.

Eliminating pressure from the monentum and hydrostatic equations by cross-differentiation gives

$$f\frac{\partial v}{\partial z} = \frac{1}{\bar{\rho}}\frac{\partial^2 p}{\partial x \partial z} = \frac{g}{\bar{T}}\frac{\partial T}{\partial x} = \frac{g}{\bar{T}}\frac{\partial T}{\partial z}\tan\varepsilon$$

$$\delta T = 0 = \frac{\partial T}{\partial x} \delta x + \frac{\partial T}{\partial z} \delta z$$

Therefore, the local slope $|\delta z/\delta x|$ of an isotherm in the frontal zone is given by

$$\tan \varepsilon = -\frac{\delta z}{\delta x} = \frac{\partial T/\partial x}{\partial T/\partial z}$$

Note that $\delta x > 0$ implies $\delta z < 0$ if, as assumed, $0 < \varepsilon < \pi/2$.

Eliminating pressure from the monentum and hydrostatic equations by cross-differentiation gives

$$f\frac{\partial v}{\partial z} = \frac{1}{\bar{\rho}}\frac{\partial^2 p}{\partial x \partial z} = \frac{g}{\bar{T}}\frac{\partial T}{\partial x} = \frac{g}{\bar{T}}\frac{\partial T}{\partial z}\tan\varepsilon$$

This is simply the *thermal wind equation* relating the vertical shear across the front to the horizontal temperature contrast across it.

 $f\frac{\partial v}{\partial z} = \frac{g}{\bar{T}}\frac{\partial T}{\partial z}\tan\varepsilon$

$$f\frac{\partial v}{\partial z} = \frac{g}{\bar{T}}\frac{\partial T}{\partial z}\tan\varepsilon$$

Solving for the slope, we get

$$\tan \varepsilon = \frac{f\bar{T}\partial v/\partial z}{g\partial T/\partial z}$$

$$f\frac{\partial v}{\partial z} = \frac{g}{\bar{T}}\frac{\partial T}{\partial z}\tan\varepsilon$$

Solving for the slope, we get

$$\tan \varepsilon = \frac{f\bar{T}\partial v/\partial z}{g\partial T/\partial z}$$

Integrating across the frontal zone, we get

$$\tan \varepsilon = \frac{f\bar{T}}{g}\frac{\delta v}{\delta T}$$

where δv and δT are the changes in along-front wind and temperature across the front.

$$f\frac{\partial v}{\partial z} = \frac{g}{\bar{T}}\frac{\partial T}{\partial z}\tan\varepsilon$$

Solving for the slope, we get

$$\tan \varepsilon = \frac{f\bar{T}\partial v/\partial z}{g\partial T/\partial z}$$

Integrating across the frontal zone, we get

$$\tan \varepsilon = \frac{f\bar{T}}{g}\frac{\delta v}{\delta T}$$

where δv and δT are the changes in along-front wind and temperature across the front.

This is *Margules' formula* and relates the slope of the frontal surface to the change in geostrophic wind speed across it and to the temperature difference across it.

Note that, with $0 < \varepsilon < \pi/2$, as drawn in the figures: 1. $\delta T = T_1 - T_2 > 0$, otherwise the flow is gravitationally unstable Note that, with $0 < \varepsilon < \pi/2$, as drawn in the figures:

- 1. $\delta T = T_1 T_2 > 0$, otherwise the flow is gravitationally unstable
- **2.** $\delta v > 0$ if f > 0 i.e., there is always a cyclonic change in v across the frontal surface.

Note that, with $0 < \varepsilon < \pi/2$, as drawn in the figures:

- 1. $\delta T = T_1 T_2 > 0$, otherwise the flow is gravitationally unstable
- **2.** $\delta v > 0$ if f > 0 i.e., there is always a cyclonic change in v across the frontal surface.

Note, however, that it is not necessary that $v_1 > 0$ and $v_2 < 0$ separately; only the change in v is important.

Note that, with $0 < \varepsilon < \pi/2$, as drawn in the figures:

- 1. $\delta T = T_1 T_2 > 0$, otherwise the flow is gravitationally unstable
- **2.** $\delta v > 0$ if f > 0 i.e., there is always a cyclonic change in v across the frontal surface.
- Note, however, that it is not necessary that $v_1 > 0$ and $v_2 < 0$ separately; only the change in v is important.
- There are three possible configurations as illustrated below.



Surface isobars in Margules' stationary front model hemisphere showing the three possible cases with the cold air to the left: (left) $v_1 > 0$, $v_2 < 0$; (centre) $0 < v_2 < v_1$; (right) $v_2 < v_1 < 0$; The surface pressure variation along the line AB is also shown.

Margules' formula is a *diagnostic one* for a stationary, or quasi-stationary front; it tells us nothing about the formation (frontogenesis) or decay (frontolysis) of fronts.

Margules' formula is a *diagnostic one* for a stationary, or quasi-stationary front; it tells us nothing about the formation (frontogenesis) or decay (frontolysis) of fronts.

It is of little practical use in forecasting, since active fronts, which are responsible for a good deal of the 'significant weather' in middle latitudes, are much more dynamic.

Margules' formula is a *diagnostic one* for a stationary, or quasi-stationary front; it tells us nothing about the formation (frontogenesis) or decay (frontolysis) of fronts.

It is of little practical use in forecasting, since active fronts, which are responsible for a good deal of the 'significant weather' in middle latitudes, are much more dynamic.

Real fronts are always associated with *rising vertical motion* and are normally accompanied by *precipitation*.

Margules' formula is a *diagnostic one* for a stationary, or quasi-stationary front; it tells us nothing about the formation (frontogenesis) or decay (frontolysis) of fronts.

It is of little practical use in forecasting, since active fronts, which are responsible for a good deal of the 'significant weather' in middle latitudes, are much more dynamic.

Real fronts are always associated with *rising vertical motion* and are normally accompanied by *precipitation*.

Moreover, real cold and warm fronts are generally not stationary, but may have speeds comparable to the horizontal wind itself. We illustrate fronts in motion in the following figure.



Schematic representation of (left) a translating cold front and (right) a translating warm front as they might be drawn on a mean sea level synoptic chart for the northern hemisphere. Note the sharp cyclonic change in wind direction and the discontinuous slope of the isobars.

However, there are technical difficulties in constructing a dynamical extension of Margules' model to fronts that translate with a uniform eostrophic flow.

However, there are technical difficulties in constructing a dynamical extension of Margules' model to fronts that translate with a uniform eostrophic flow.

Nevertheless, fronts analyzed on weather charts are drawn on the assumption that this is possible, and Margules' model is found to provide a valuable if highly simplified conceptual framework.

However, there are technical difficulties in constructing a dynamical extension of Margules' model to fronts that translate with a uniform eostrophic flow.

Nevertheless, fronts analyzed on weather charts are drawn on the assumption that this is possible, and Margules' model is found to provide a valuable if highly simplified conceptual framework.

It is interesting that Margules developed his model of an atmospheric discontinuity some fifteen years before the emergence of the frontal models of the Norwegian School. There were other precursors of frontal theory in Germany and in Britain. **Exercise:** Check the dimensional consistency of Margules' Formula.

Calculate the frontal slope using Margules' Formula, assuming that the mean temperature is $\overline{T} = 280 \text{ K}$, the Coriolis parameter $f = 10^{-4} \text{ s}^{-1}$, $g = 10 \text{ m s}^{-2}$, the difference in windspeed across the front is $\delta v = 12 \text{ m s}^{-1}$ and the difference in temperature is $\delta t = 4 \text{ K}$.

Exercise: Check the dimensional consistency of Margules' Formula.

Calculate the frontal slope using Margules' Formula, assuming that the mean temperature is $\overline{T} = 280 \text{ K}$, the Coriolis parameter $f = 10^{-4} \text{ s}^{-1}$, $g = 10 \text{ m s}^{-2}$, the difference in windspeed across the front is $\delta v = 12 \text{ m s}^{-1}$ and the difference in temperature is $\delta t = 4 \text{ K}$.

Solution: ...

Answer: $\epsilon \approx 1/120$.