§5.5: Ekman Pumping

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$$D = \frac{\pi}{\gamma} = \pi \sqrt{\frac{2K}{f}} = \pi \sqrt{\frac{2 \times 5}{10^{-4}}} = 993 \,\mathrm{m} \approx 1 \,\mathrm{km}$$

Thus, the effective depth of the Ekman boundary layer is about one kilometre.

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Note that D depends on the values of f and K so the particular value 1 km is more an indication of the scale that a sharp quantitative estimate.

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The modified Ekman Layer is discussed on Holton (§5.3.6). We will not discuss it here.

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We will now calculate the vertical velocity at the top of the Ekman layer.

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The result is thus

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Here we have used the fact that

$$e^{-\pi} \approx 0.0432 \ll 1$$

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For example, it can be done by means of integration by parts (twice), or by expressing the sin-function in terms of complex exponentials.

Next, integrate the continuity equation through the PBL:

$$\int_0^D \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z}\right) dz = \int_0^D \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}\right) dz + [w(D) - w(0)] = 0$$

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$$\left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}\right) = \frac{\partial v}{\partial y} = \left(\frac{\partial u_{\rm g}}{\partial y}\right) e^{-\pi z/D} \sin(\pi z/D)$$

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Substituting this into the equation for w(D) gives

$$w(D) = -\frac{\partial u_{\rm g}}{\partial y} \int_0^D e^{-\pi z/D} \sin(\pi z/D) \, dz = -\frac{1}{2} \left(\frac{D}{\pi}\right) \frac{\partial u_{\rm g}}{\partial y}$$

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We now note that the geostrophic vorticity is given by

$$\zeta_{\rm g} = \left(\frac{\partial v_{\rm g}}{\partial x} - \frac{\partial u_{\rm g}}{\partial y}\right) = -\frac{\partial u_{\rm g}}{\partial y}$$

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$$\begin{bmatrix} \mathbf{Cyclonic} \\ \mathbf{Flow} \end{bmatrix} \Longleftrightarrow \begin{bmatrix} \mathbf{Positive} \\ \mathbf{Vorticity} \end{bmatrix} \Longleftrightarrow \begin{bmatrix} \mathbf{Upward} \\ \mathbf{Velocity} \end{bmatrix}$$

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Note on **Dines Mechanism** to be added later.

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$$w(D) = \left(\frac{1}{2\pi}\right) \times 10^3 \times (5 \times 10^{-5}) = \frac{5 \times 10^{-2}}{2\pi} \approx 8 \,\mathrm{mm\,s^{-1}} \sim 1 \,\mathrm{cm\,s^{-1}}$$

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If it is sufficient to lift air to its LCL, then latent heat release allows stronger updrafts within the convective clouds.

Storms in Teacups



Standing waves in a tea cup, induced by the propeller rotation of an airoplane.

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The centrifugal force is given, as usual, by

$$\frac{V^2}{r} = \omega^2 r$$

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This can be integrated immediately to give

$$p = p_0 + \frac{1}{2}\rho_0\omega^2 r^2$$

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As a result, there is radial inflow near the bottom. By continuity of mass, this must result in upward motion near the centre.

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Furthermore, outflow must occur in the fluid above the boundary layer.



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This secondary circulation is completed by downward flow near the edges of the container.

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The geostrophic vorticity is given, in cylindrical coordinates, by

$$\zeta_{\rm g} = \mathbf{K} \cdot \nabla \times \mathbf{V}_{\rm g} = \frac{1}{r} \left[\frac{\partial (rV_{\rm g})}{\partial r} - \frac{\partial U_{\rm g}}{\partial \theta} \right] = \frac{1}{r} \frac{\partial (\omega r^2)}{\partial r} = 2\omega$$

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We may compare this to the azimuthal velocity. At r = 5 cm we have

$$V_{\rm g} = \omega r = 2\pi \,\mathrm{s}^{-1} \times 5 \,\mathrm{cm} \approx 30 \,\mathrm{cm} \,\mathrm{s}^{-1}$$

Thus, the secondary circulation is relatively weak compared to the primary (solid rotation) circulation. But it is dynamically important.

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Exercise: Create a storm in a teacup: Stir your tea (no milk) and observe the leaves.

Exercise:

- Calculate the mass influx through the sides of a cyclone.
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- Deduce an expression for the vertical velocity.

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For solid body rotation, $u_g = \omega R$ and the geostrophic vorticity is $\zeta_g = 2\omega$, so

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Assuming w(H) = 0 and substituting the Ekman pumping for W(D) we get

$$\frac{d\zeta_{\rm g}}{dt} = -\frac{f}{(H-D)} \left(\frac{1}{2\pi}\right) D\zeta_{\rm g}$$

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 $\zeta_{\rm g} = \zeta_{\rm g}(0) \exp(-t/\tau_{\rm Ekman})$

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The size of τ_{Ekman} may be estimated for typical values:

$$\tau_{\rm Ekman} = \frac{2\pi H}{fD} = \frac{2\pi \times 10^4}{10^{-4} \times 10^3} \approx 6 \times 10^5 \,\mathrm{s}$$

which is about seven days.

$$\frac{\partial u}{\partial t} = K \frac{\partial^2 u}{\partial z^2}$$

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For the values already assumed $(K = 5 \text{ m}^2 \text{ s}^{-1} \text{ and } H = 10 \text{ km})$ we get

$$\tau_{\rm Diff} = \frac{H^2}{K} \approx \frac{10^8}{5} = 2 \times 10^7 \, {\rm s}$$

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However, cumulonumbus convection can produce rapid transport of heat and momentum through the entire troposphere.

Ekman Layer in the Ocean



Ekman spiral in the ocean.



Typical La Niña Pattern



Mean sea surface temperature, eastern Pacific Ocean 5 September to 5 October, 1998.

End of §5.5